Ab initio theory and simulation of semiconductor nanostructures

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Introduction on nanoscience

Ab initio simulations: Density Functional Theory (DFT)

Quantum confinement in Si nanowires (Si NWs)

- Electronic structure
- Transport properties
- Doping effects

Crystal phase engineering in Si NWs

- Crystal phase effects on bulk
- Crystal phase effects on NWs
- STEM-EELS experiments



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Nanoscience





Nanoscience



The study of structures, dynamics, and properties of systems in which one or more of the spatial dimensions is nanoscopic (1-100 nm) Dynamics and properties that are distincly different (often in an extraordinary way) from both small-molecule systems and systems macroscopic in all dimensions U.S. National Nanotechnology Initiative, nano.gov







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The importance of size



The importance of size



Gold nanoparticles



As the size of the material is reduced, and the nanoscale regime is reached, it is possible that the same material will display totally different properties



Nano building blocks



Nano building blocks

What are the nano building blocks that would play an analogous role to macro building blocks ?



Nano building blocks

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Clusters and molecular nanostructures





Nanotubes and related systems

Quantum wells, wires, films and dots





2D materials

From nanoscience to ...



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From nanoscience to ...



... nanotechnology



... nanotechnology



Biomedicine



Nanoelectronics



Biomaterials



Quantum information



Optoelectronics







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The application of new extraordinary experimental tools has created an urgent **need for a quantitative understanding** of matter at nanoscale



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- The **absence of quantitative models** that describe newly observed phenomena increasingly limits progress in the field



- The application of new extraordinary experimental tools has created an urgent **need for a quantitative understanding** of matter at nanoscale
- The **absence of quantitative models** that describe newly observed phenomena increasingly limits progress in the field
- The absence of such tools would also **seriously inhibit widespread applications** ranging from molecular electronics to biomolecular materials





 $J = -k \nabla T$



 $J = -k \nabla T$

Is the Fourier's law applicable to nanostructures?

How size effects can affect thermal conductivity?

Can present theories interpret satisfactorily experiments?



 $J = -k \nabla T$

Is the Fourier's law applicable to nanostructures?

How size effects can affect thermal conductivity?

Can present theories interpret satisfactorily experiments?



Measured thermal conductivity lower than bulk and diameter dependent

Quantum effects not negligible

Breakdown of Fourier's law

Where do we need a theoretical effort?



Where do we need a theoretical effort?

Transport in nanostructures



Coherence and decoherence tunneling

Electronic and optical properties



Soft/hard matter interfaces



Spintronics





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Length scales of materials modelling







Ab initio computational modelling



Ab initio computational modelling

The developement and use of mathematical models for describing and predicting certain properties of materials at a **quantitative level**

'Ab initio' or 'from first principles' refers to a bottom-up modelling strategy in which we **do not use any empirical parameters**

Such kind of calculations are completely based on **quantum mechanics**, that can be considered as an engineering tool

The complexity of such problems requires the use of **supercomputers**

This is a discipline at the boundary between **materials science, physics and chemistry** on the one side, and **applied mathematics and software engineering**



Precision of ab initio modelling




Density Functional Theory



Density Functional Theory

DFT is a very effective technique for studying molecules, nanostructures, solids, surfaces and interfaces by directly solving approximate versions of the Schrödinger equation



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M. Dumaz et al. Scientometrics 126, 6681 (2021)

The many body electronic structure



The many body electronic structure



N_e electrons N_n nuclei

Schrödinger equation for interacting particles $\widehat{H}\Psi(\{R\}\,,\{r\})=E\Psi(\{R\}\,,\{r\})$

 $\begin{aligned} \widehat{H} &= \widehat{T}_n(\{R\}) + \widehat{V}_{nn}(\{R\}) + \widehat{T}_e(\{r\}) + \\ &+ \widehat{V}_{ee}(\{r\}) + \widehat{U}_{en}(\{R\}, \{r\}) \end{aligned}$



The many body hamiltonian



The many body hamiltonian

$$\widehat{T}_n = \sum_{I=1}^{N_n} -\frac{\nabla_I^2}{2M_I}$$
 Kinetic energy of nuclei

7 7

$$\widehat{T}_e = \sum_{i=1}^{N_e} -\frac{\nabla_i^2}{2m_e}$$

Kinetic energy of electrons

$$\widehat{V}_{nn} = \frac{1}{2} \sum_{I,J,I \neq J}^{N_n} \frac{Z_I Z_J}{|R_I - R_J|}$$

Ion-ion interactions

$$\widehat{V}_{ee} = \frac{1}{2} \sum_{i,j,i \neq j}^{N_e} \frac{1}{|r_i - r_j|}$$

Electron-electron interactions

$$\widehat{U}_{en} = -\sum_{j,J}^{N_e,N_n} \frac{Z_J}{|R_J - r_j|}$$

Electron-ion interactions

Exactly solvable for two particles (analytically) and very few particles (numerically)

How to deal with N ~ 10^{23} particles?



Density Functional Theory (DFT)



Density Functional Theory (DFT)

It can map, exactly, the interacting problem to a non-interacting one

Hohenberg and Kohn, Phys. Rev. B 136, 864 (1964)



W. Kohn Nobel Prize in Chemistry in 1998



interacting particles in a real external potential a set of non-interacting electrons (with the same density as the interacting system) in some effective potential



Popularity of DFT



Popularity of DFT

Transferability: same technique for describing different classes of materials

Simplicity: based on simply and intuitive equations

Realiability: possibility of making direct and quantitative comparison with experiments

Software sharing: online platforms and adoption of opensource software model

Reasonable starting point: even when it fails in describing correctly a property, it represents an accurate starting point for more accurate theory (more computationally demanding)



Some specific comments

Some specific comments

In principle DFT is exact and pure predictive, but in practice it needs some approximations (exchange-correlation functional)

Good scaling of computational cost with system size

Calculations on large and complex systems (surfaces, interfaces, nanostructures, defects)

Accuracy of properties prediction is well known (i.e. 3% for bond length, 10% for bulk modulus and phonon frequencies)

Good compromise between accuracy and computational time



Ab initio properties of materials



Ab initio properties of materials

STRUCTURAL PROPERTIES

Density Functional Theory, Total energies, Forces, Stresses, Atomic and cell geometries

SURFACE PROPERTIES

Nudge Elastic Bands Activation barrier, Surface and adsorption energies

ELECTRONIC STRUCTURE Density Functional Theory,

Band structure, Density of states, Wave function, OPTICAL RESPONSE Many Body Perturbation Theory, Optical Absorption, Excitons, GW approach, EELS

VIBRATIONAL PROPERTIES

Density Functional Perturbation Theory (DFPT), Raman spectra, phonons, phase diagrams

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ELASTIC PROPERTIES

Density Functional Theory, Elastic constants and properties of solids

Ab initio properties of materials

STRUCTURAL PROPERTIES

Density Functional Theory, Total energies, Forces, Stresses, Atomic and cell geometries

ELECTRONIC STRUCTURE

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OPTICAL RESPONSE Many Body Perturbation Theory, Optical Absorption, Excitons, GW approach, EELS

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ELASTIC PROPERTIES

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Theoretical method and codes



Theoretical method and codes

siesta

SIESTA code http://icmab.es/siesta

Density Functional Theory (DFT)

- Exchange-correlation functional: LDA
- Optimized double-z polarized basis-set
- Troullier-Martins pseudopotentials



VASP code http://www.vasp.at

Hybrid - DFT

- Exchange-correlation functional: PBE-HSE06
- Plane-wave basis set
- PAW approach pseudopotentials



YAMBO code <u>www.yambo-code.org</u>

GW and optical properties calculations

- GW one short perturbative approach
- Plane-wave basis set
- RPA and BSE approaches



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Space for electrons in materials



Space for electrons in materials



T. Tsurumi et al. Nanoscale physics for materials science, Taylor and Francis (2010)

- In macroscopic materials the Fermi wave length is of the same order as the unit cells
- Thus the inner space is too large for electrons to feel boundaries
- The possible energy levels for electrons E(k)=(ħ²/2m*)k² are dense with very small energy spacing

Quantum confinement effect (QCE)



Quantum confinement effect (QCE)



T. Tsurumi et al. Nanoscale physics for materials science, Taylor and Francis (2010)

Quantum confinement effect (QCE) is defined as: a reduction in the degrees of freedom of the carrier particles, implying a reduction in the allowed phase space

E.G. Barbagiovanni et al. Appl. Phys. Rev. 1, 011302 (2014)



Consequences of QCE



Consequences of QCE

Electrons traveling along the direction that has been
reduced in size reach the boundaries and are confined

- This confinement causes the quantization of the electron wavelength and of the energy spectrum
- The contribution of edge-localized surface states can become relevant

Size effects dominate the physics of nanostructures



The case of Si nanowires (Si NWs)



The case of Si nanowires (Si NWs)

One-dimensional nanostructures with precise composition, morphology, interfaces and electrical properties (d=2-3 nm)

Vapor-Liquid-Solid (VLS) growth



O. Hayden et al., Nanotoday 3 5-6 (2008)

Lithography and etching processes



N. Singh et al., IEEE Trans. Electron Devices 55, 11 (2008)





Real system





Quantum Mechanics (2000)



The motion of electrons is restricted to be in the direction of confinement. Their kinetic energy increases and the eigenstates energies are given by:

$$E_n = \frac{\hbar^2 n^2 \pi^2}{2m^* d^2}$$

m*=effective mass d=width of the well

Not only the energy levels but even the spacing between them increase when the diameter is reduced (QCE increases)

QCE has a dramatic effect on semiconductor NWs (like Si NWs) because it affects the **energy band gap**





Ma, D. D. D. et al., Science 299, 1874 (2003)

STS measurements



Ma, D. D. D. et al., Science 299, 1874 (2003)

$$E_{gap} = E_{bulk, gap} + C(1/d)^{\alpha}$$

with $\alpha = 1.1 - 1.4$

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M. Bruno et al., Phys. Rev. Lett. 98, 036807 (2007)

QCE in Si NWs: DFT scaling of the gap



QCE in Si NWs: DFT scaling of the gap





 $\alpha_{hex} = 1.16$ $\alpha_{cub} = 1.34$

M. Amato et al., Nano Lett. 16, 5694 (2016) 26
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QCE in NWs: energy quantization



QCE in NWs: energy quantization



The total energy of electrons is given by the formula:





QCE in NWs: Landauer approach



QCE in NWs: Landauer approach



Electrons move with a group velocity v(k) and transmission probability T(k)



QCE in NWs: quantized conductance



QCE in NWs: quantized conductance



T. Tsurumi et al., Taylor and Francis (2010)

Current

$$I = \frac{2e}{h} \int_0^\infty \left\{ f(E,\mu) - f(E,\mu-eV) \right\} T(E) dE$$

Conductance (if the bias is small)

$$G = \frac{2e^2}{h} \int_0^\infty \left\{ -\frac{\partial f(E,\mu)}{\partial E} \right\} T(E) dE \equiv \frac{2e^2}{h} T_\mu$$



QCE in NWs: quantized conductance

- In 1D systems the conductance of electrons is quantized in units of 2e²/h
- Electrons states along the wire are associated with quantized states in the plane
- Each of the quantum states in the plane has equal unit of conductance $(2e^2/h)T_u$ along the wire axis (where T_u is dimensionless)
- The conductance of NWs is hence the product of the number of quantum states and their quantized conductance $(2e^2/h)T_{\mu}$



Experimental demonstration on SiGe core-shell NWs

QCE in Si NWs: DFT+Landauer approach



QCE in Si NWs: DFT+Landauer approach

Conductance is calculated in terms of transmission probability T(E) through the available transmitting channels

> How many? For an infinitely long NW: as many as electron states at that energy





For a pristine Si NW: T(E) = 1

For a defected NW: T(E) < 1

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- Ionization energy ~ few hundredths of eV
- Ionized impurity at room temperature
- High doping efficiency

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- Ionization energy ~ few hundredths of eV
- Ionized impurity at room temperature
- High doping efficiency

- Ionization energy increase
- Donor level deep into the band gap
- Doping deactivation with respect to bulk

QCE in Si NWs: P deactivation in DFT



QCE in Si NWs: P deactivation in DFT



M. Amato et al., Nano Lett. 11, 594–598 (2011)

P-impurity state deep into the band gap Activation energy is too high (~ 100 meV)



Some comments on QCE in Si NWs



Some comments on QCE in Si NWs

The diameter is a powerful tool to modulate physical properties

- Electronic Structure. Reducing the size can induce an opening of the bulk band gap that can cause a modulation of optical properties
- **Transport.** The quantization due to the confinement leads to a quantization of the electrical conductance. Si NWs are hence quantum confined ballistic conductors
- **Doping.** Intentional addition of impurities in Si NWs is not efficient as for bulk systems due to the impurity deactivation



Some comments on QCE in Si NWs

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- **Doping.** Intentional addition of impurities in Si NWs is not efficient as for bulk systems due to the impurity deactivation

Is it possible to further modify the physical properties of Si NWs?



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All these experiments confirmed that, as can be intuitively expected, Si NWs crystallize in the diamond-like cubic structure as in their bulk counterpart (3C phase)





NWs growth











M. Amato et al., Chem. Rev. 114, 1371 (2014)







L. Tizei and M. Amato, Eur. Phys. J. B 93, 16 (2020) 36
Hexagonal-diamond (2H) NWs



Hexagonal-diamond (2H) NWs

Hexagonal Si nanoshells on top of hexagonal GaP



H. T. Hauge et al., Nano Lett. 15, 5855 (2015)



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Hexagonal SiGe nanoshells on top of hexagonal GaP



H. T. Hauge et al., Nano Lett. 17, 85-90 (2017)



Hexagonal-cubic junctions 2H/3C NWs



Hexagonal-cubic junctions 2H/3C NWs

Hexagonal-cubic diamond Ge NWs homojunctions



L. Vincent et al., Nano Lett. 14, 4828 (2014)



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Hexagonal-cubic diamond Si NWs homojunctions



J. Tang et al., Nanoscale 9, 8113 (2017)



A not obvious result for bulk



A not obvious result for bulk

Bulk hexagonal diamond Si metastable and obtained only at high pressures



Pirouz et al., Acta Metall. Mater. 38, 313 (1990)



A not obvious result for bulk

Bulk hexagonal diamond Si metastable and obtained only at high pressures



Pirouz et al., Acta Metall. Mater. 38, 313 (1990)



J. Z. Hu and I. L. Spain, Sol. State Comm. 51, 263 (1984)



Light emission from 2H-SiGe

Light emission from 2H-SiGe



Light emission from 2H-SiGe



tow (nm) W.H.J. Peeters et al. Nat. Commun. 15, 5252 (2024)

15

20

25

30

10

5

0.40

0.35

35

quantum wells

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Band structure of 2H and 3C bulk



Band structure of 2H and 3C bulk







DFT – HSE06

Element	Crystal structure	$E_g \text{ at } \Gamma \text{ (eV)}$	$E_g (eV)$	E_g (SOC) (eV)	Band-edge	Notes
	Cubic-diamond (3C)	3.29	1.14		$\Gamma - X$	This work (HSE06)
		3.20	1.14		$\Gamma - X$	Cal. Ref. ¹⁶
Si		3.33	1.15		$\Gamma - X$	Cal. Ref. ³³
		3.40	1.17		$\Gamma - X$	Exp. Ref. ⁶¹
	Hexagonal-diamond (2H)	1.70	0.95		$\Gamma - M$	This work (HSE06)
		1.63	0.95		$\Gamma - M$	Cal. Ref. ¹⁶
	Cubic-diamond (3C)	0.84	0.72	0.61	$\Gamma - L$	This work (HSE06)
		0.94	0.85	0.71	$\Gamma - L$	Cal. Ref. ³³
\mathbf{Ge}		0.81	0.66		$\Gamma - L$	Exp. Ref. ⁵⁴
	Hexagonal-diamond (2H)	0.32		0.23		This work (HSE06)
	Hexagonal-diamond (2H)	0.23				Cal. Ref. ⁵⁶
		0.31				Cal. Ref. ¹⁷

Table 2: The calculated band gaps of Si and Ge.



DFT – HSE06

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DFT – HSE06

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GW – BSE







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For both Si and Ge bulk, the **2H phase** has a **smaller band gap** with the respect to the 3C one



For both Si and Ge bulk, the **2H phase** has a **smaller band gap** with the respect to the 3C one

2H Ge bulk phase is characterized by a **direct band gap** whose optical transition has **a low oscillator strength** [PRM 3, 034602 (2015)]



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2H Ge bulk phase is characterized by a **direct band gap** whose optical transition has **a low oscillator strength** [PRM 3, 034602 (2015)]

2H Si bulk absorbs more in the **visible (above 2.5 eV)** and its spectrum presents a **big overlap** near the **maximum** of the solar spectrum



For both Si and Ge bulk, the **2H phase** has a **smaller band gap** with the respect to the 3C one

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What happens when we reduce the size? Should we expect the same electronic and optical behaviour?



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Quantum confinement in 3C Si NWs



Quantum confinement in 3C Si NWs



Ma, D. D. D. et al., Science 299, 1874 (2003)

STS measurements



Ma, D. D. D. et al., Science 299, 1874 (2003)

$$E_{gap} = E_{bulk, gap} + C(1/d)^{\alpha}$$

with $\alpha = 1.1 - 1.4$

M. Bruno et al., Phys. Rev. Lett. 98, 036807 (2007)

Band gap of 3C and 2H Si NWs



Band gap of 3C and 2H Si NWs





Band gap of 3C and 2H Si NWs





 $\alpha_{cub} = 1.34$
Bands and optical absorption

DFT – LDA



M. Amato et al., Nano Lett. 16, 5694 (2016)



Bands and optical absorption

DFT – LDA

IP – RPA



M. Amato et al., Nano Lett. 16, 5694 (2016)



Bands and optical absorption

DFT – LDA

IP – RPA





Introduction on nanoscience

Ab initio simulations: Density Functional Theory (DFT)

Quantum confinement in Si nanowires (Si NWs)

- Electronic structure
- Transport properties
- Doping effects

- Crystal phase effects on bulk
- Crystal phase effects on NWs
- STEM-EELS experiments



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Dielectric and loss functions of 2H–Si and 2H–Ge nanowires can be computed from first principles and compared to experiments



Synthesized high-quality 2H–Si and 2H–Ge nanowires via vapor–liquid–solid growth on wurtzite GaAs stems





L. Tizei et al., Nano Lett. 25, 8604-8611 (2025)



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Nanowires are defect-free, stress-free, and contamination-free, ideal for optical studies





















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For 2H–Ge NWs, the intersecting-specimen EELS shows increased absorption near 0.9 eV, consistent with ab initio calculations (no absorption near the band gap)









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An intense peak at 2 eV is observed in 2H-Ge, attributed to a thin 3C–Ge shell formed during growth

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This coupling helps interpret complex EELS signals and guides the design of nanoscale optoelectronic devices

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Bibliography



FELICIANO GIUSTINO

Materials modelling using Density Functional Theory F. Giustino, Oxford Univeristy Press 2014

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