

Characterizing the electronic and optical properties of semiconductors

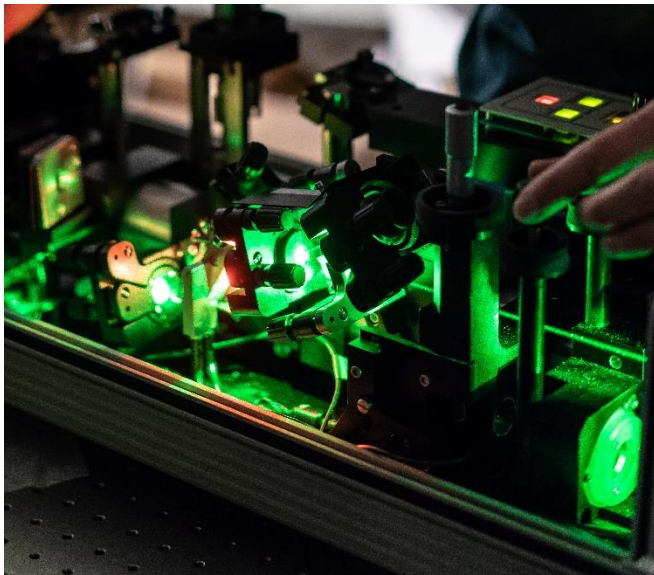
Delphine LAGARDE
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MATEPI Summer School
Porquerolles - June 2025

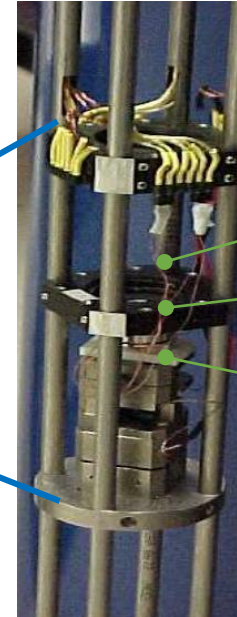
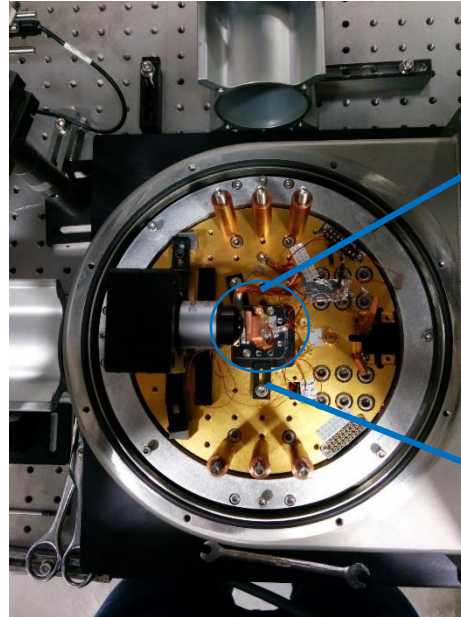


LPCNO lab – OPTO team

Laser

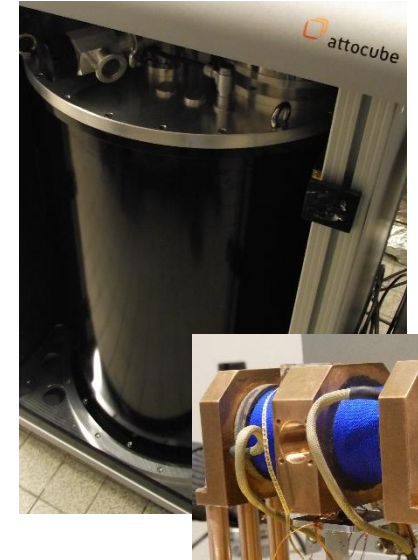


Cryostat 4K

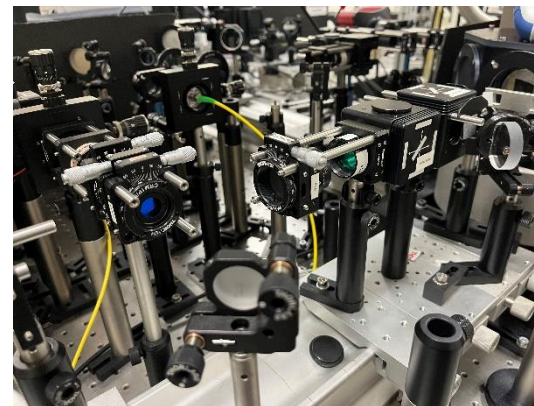


Objective
Sample
Nano-
positioners

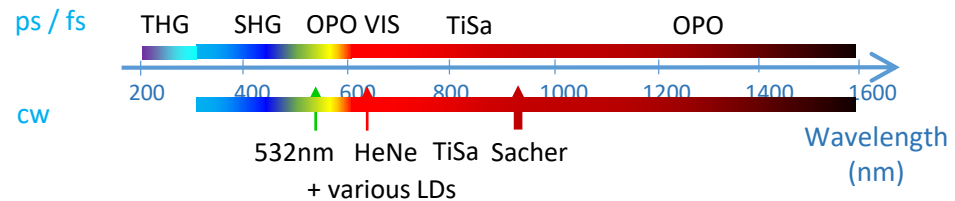
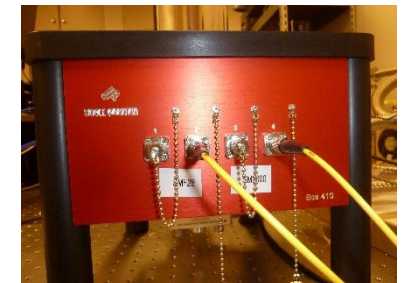
Magnetic field



Optics

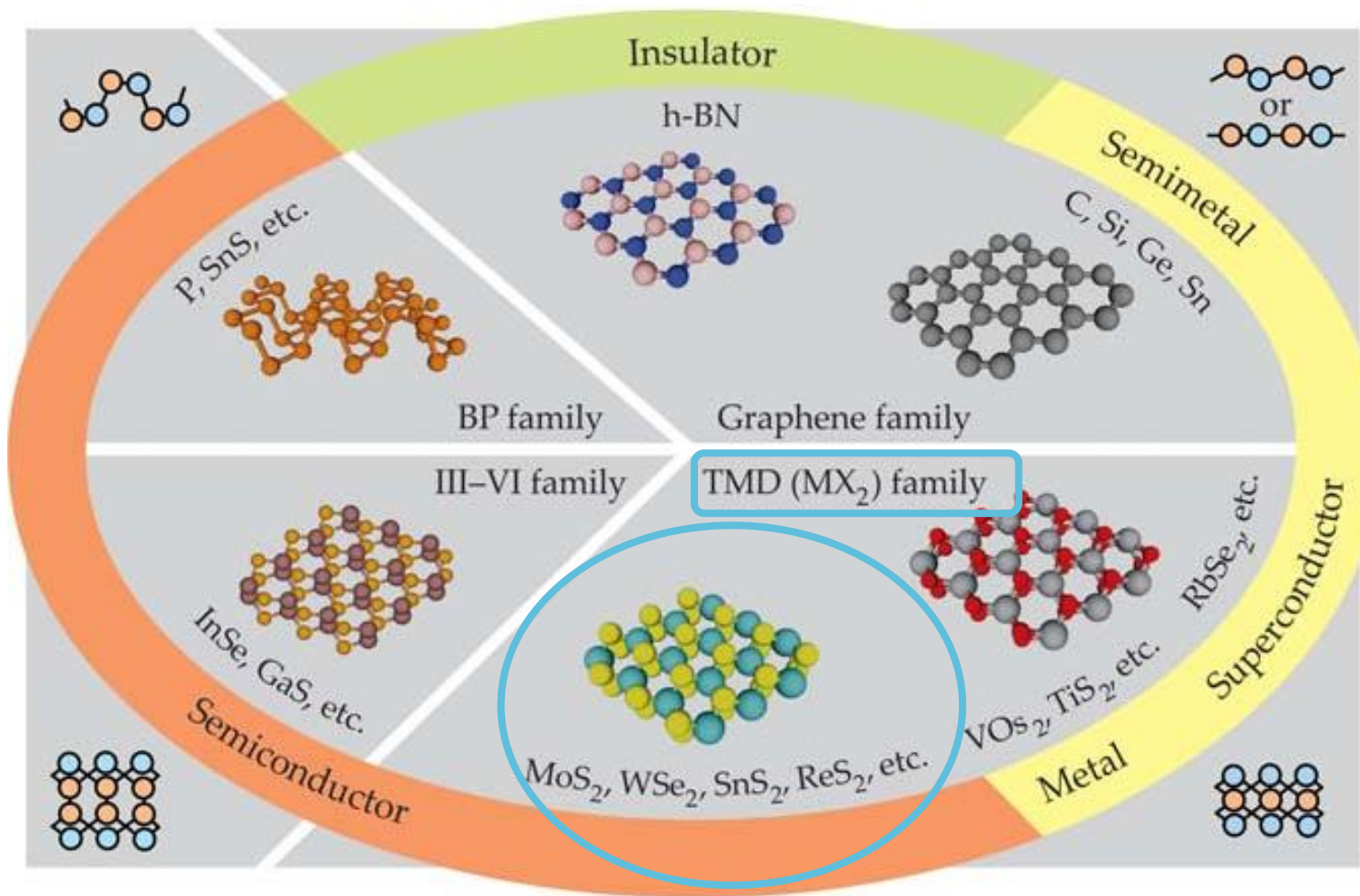


Spectrometer /
Detector



Introduction

- Large variety of methods that exploit light-matter interactions across the entire electromagnetic spectrum
- This presentation aims to give an overview of some available techniques and their potential applications for electronic and optical characterization of semiconductor epitaxial layers or micro/nanostructures



Outline

Basics of SC physics

- * SC band structure
- Density of states
- Doping
- * Interaction with light
- * Dimensionality

Optical characterization of SC structures

- * Absorption
- * Photoemission
- * Reflectivity
- * Photoluminescence (PL)
 - PL Eexcitation (PLE)
- * Time-Resolved PL
 - TR-Cathodoluminescence

Electric/Optical characterization of SC structures

- * Electroluminescence (EL)
- * Photocurrent (PC)

Raman spectroscopy

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Basics of SC physics

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Density of states
Doping
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Optical characterization of SC structures

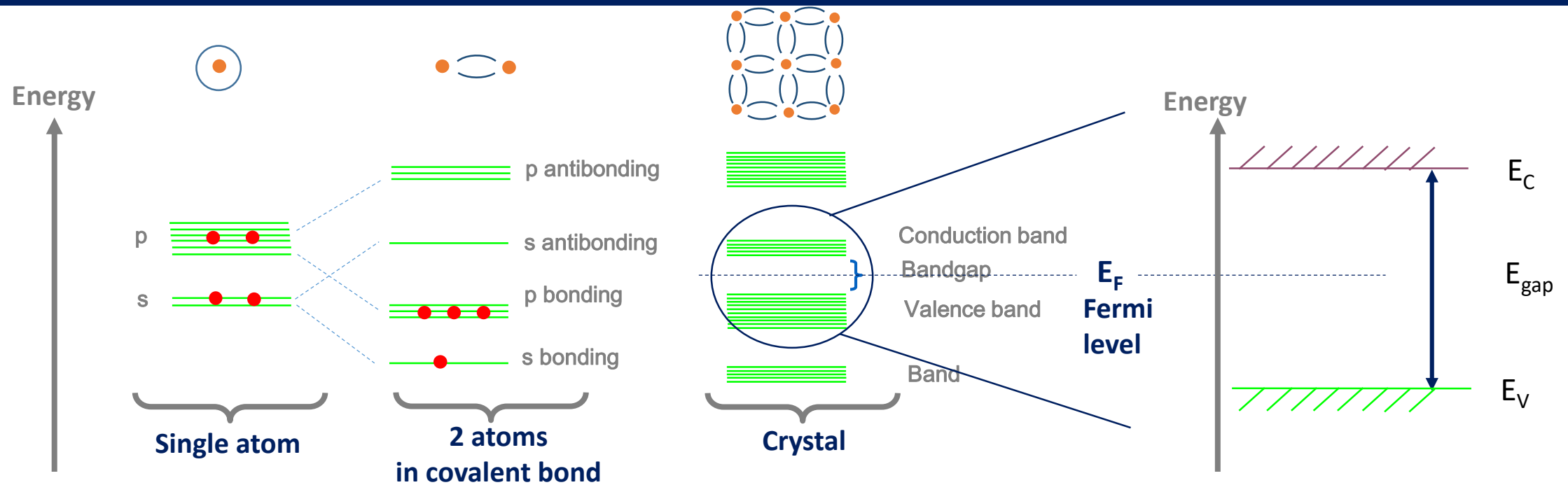
- * Absorption
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Electric/Optical characterization of SC structures

- * Electroluminescence (EL)
- * Photocurrent (PC)

Raman spectroscopy

Band structure



Electrons in atom

- Electrons can be described by Schrödinger equation
- Ground vs excited state
- Pauli exclusion principle

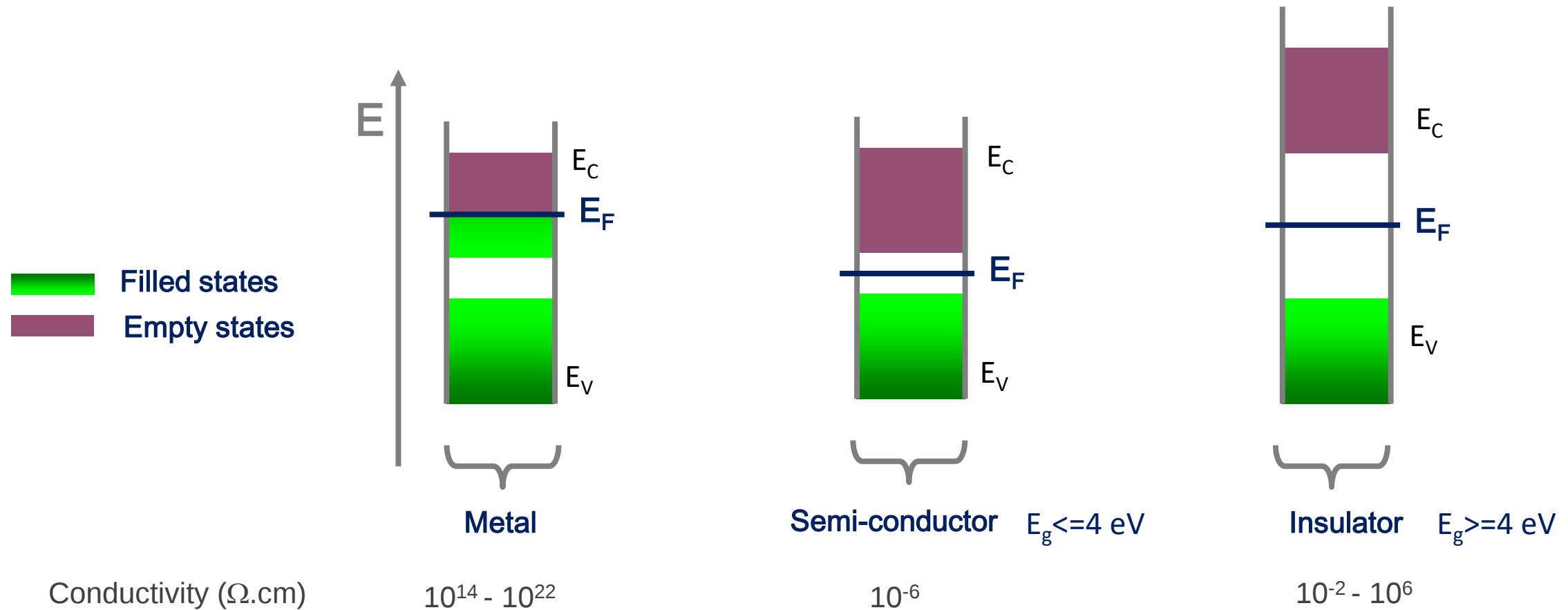
Electrons in interacting atoms

- Outer orbitals of the atoms overlap and interact strongly

Electrons in crystal

- Highest nearly-filled band : **valence band**
- Lowest nearly-empty band : **conduction band**

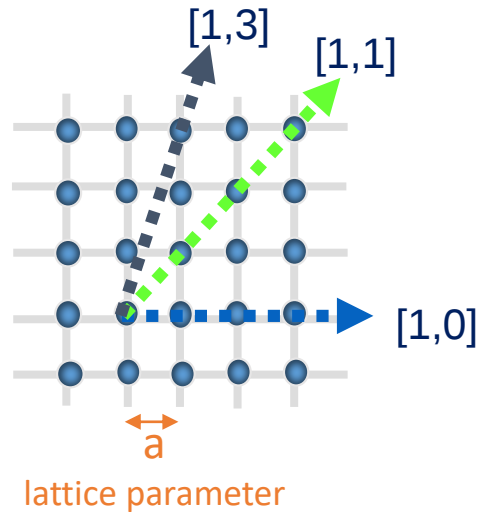
Band structure



- metal: weakly bound electrons → free mobile electrons
- insulator: every electrons bound to the lattice
- semiconductor: some electrons can gain enough E to get into the conduction band

Band structure

Crystal symmetry



32 classes of crystal according to their point group symmetry (invariance by translations, rotations...)

Crystal symmetry will affect electronic wavefunctions and optical and transport properties

Crystal periodicity

Periodic arrangement of atoms in a crystal



periodic potential $V(\vec{r} + \vec{R}) = V(\vec{r})$ with $\vec{R}$ reciprocal lattice vector

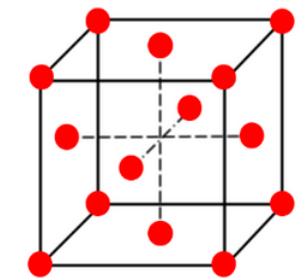
Wavefunctions of an electron in a crystal

Block's theorem

$$\psi_{n,\vec{k}}(\vec{r}) = u_{n,k} e^{i\vec{k} \cdot \vec{r}}$$

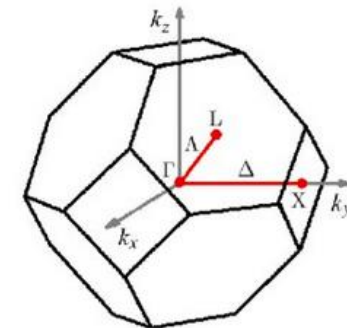
$\vec{k}$: wave vector in reciprocal space (linked to crystal translation invariance)
 $u_{n,k}$ has the periodicity of the crystal lattice $u_{n,k}(\vec{r} + \vec{R}) = u_{n,k}(\vec{r})$

Real space



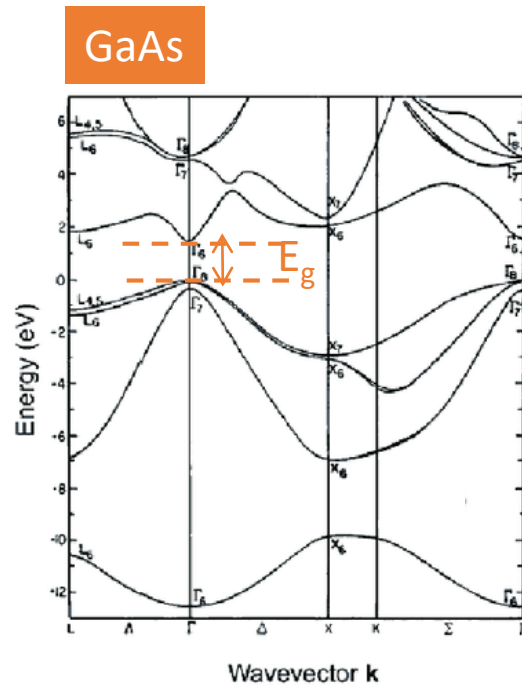
Cubic face centered

Reciprocal space



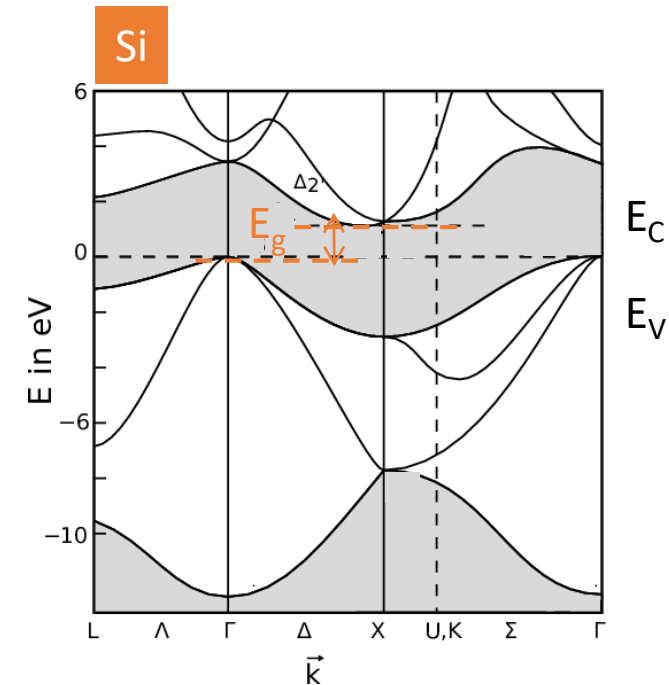
Band structure

Optical properties : consider only the band structure in a region close to band extrema, near $\vec{k}=0$, center of Brillouin zone



GaAs: $E_g = 1.43$ eV (300K)

Direct Bandgap

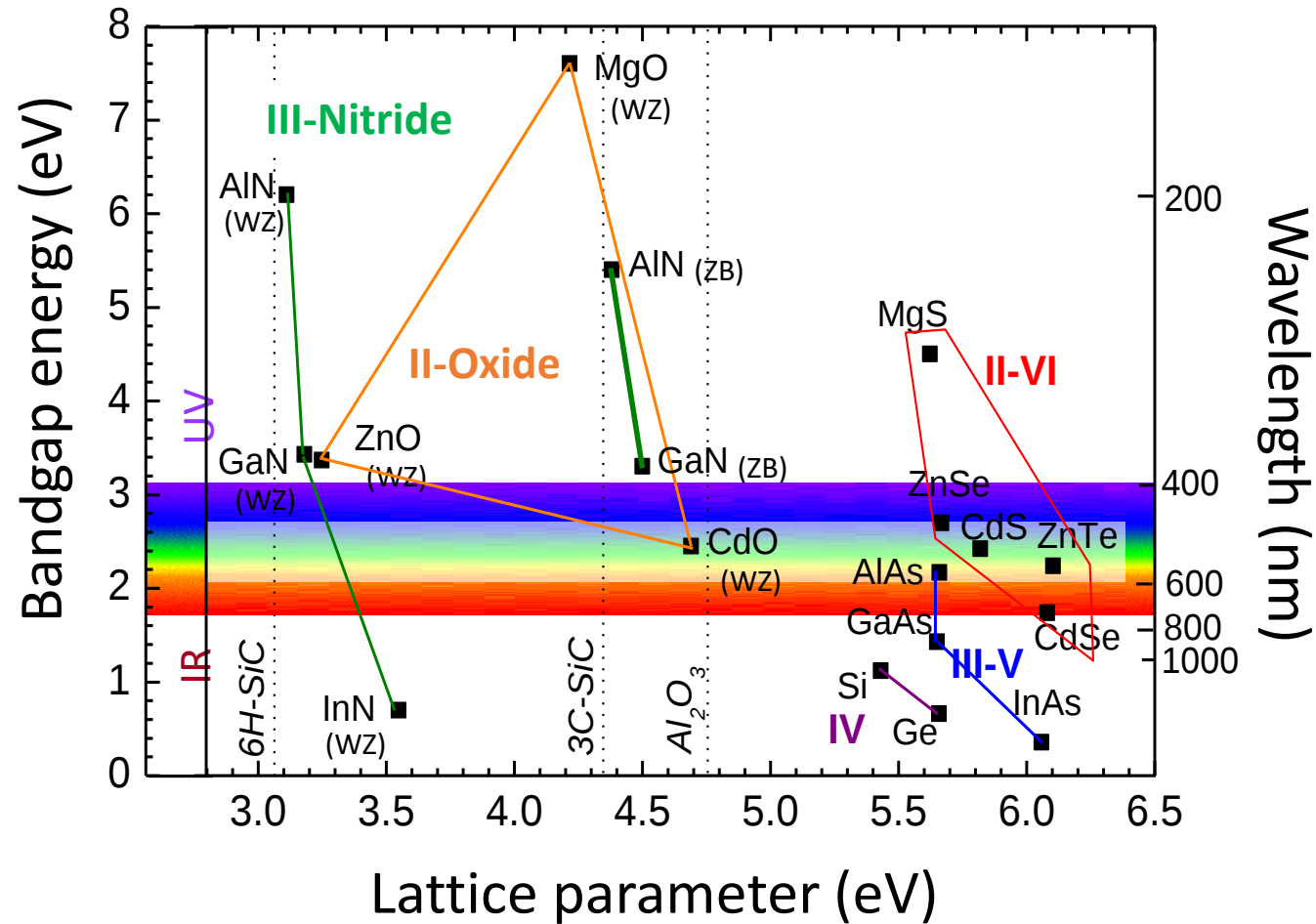


Silicon: $E_g = 1.1$ eV

Indirect Bandgap

Max of VB and min of CB occur at the same point of the first Brillouin zone
 → **Strong interaction with light**

Bandgap engineering



Density of states

for 3D crystal

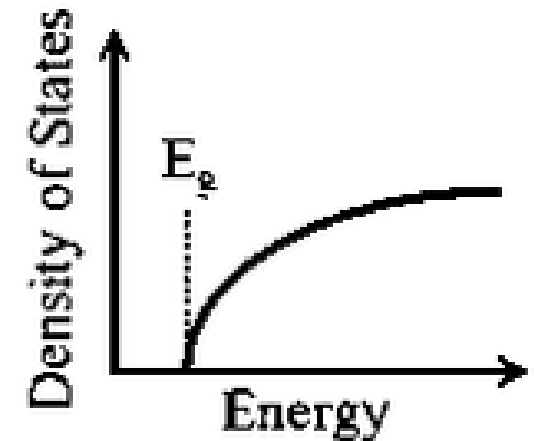
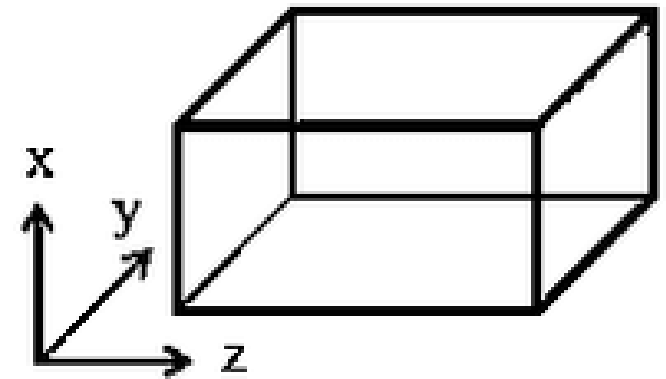
CB

$$\rho(E)_c = \frac{1}{2\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} \sqrt{E - E_c}$$

VB

$$\rho(E)_v = \frac{1}{2\pi^2} \left(\frac{2m}{\hbar^2} \right)^{3/2} \sqrt{E_v - E}$$

Bulk (3-D)



The density of states $\rho(E)$ is a completely different function of carrier energy for 3D bulk, 2D quantum wells or 0D quantum dots

Carrier density

Carrier density linked to probability of occupation of each electronic states at thermal equilibrium

$$n_0 = \int_{E_c}^{\infty} \rho_c(E) f(E) dE$$

$$p_0 = \int_{-\infty}^{E_v} \rho_v(E) [1 - f(E)] dE$$



Concentration of electrons/holes per unit volume:

$$n_0 = N_c e^{(E_F - E_c)/kT}$$

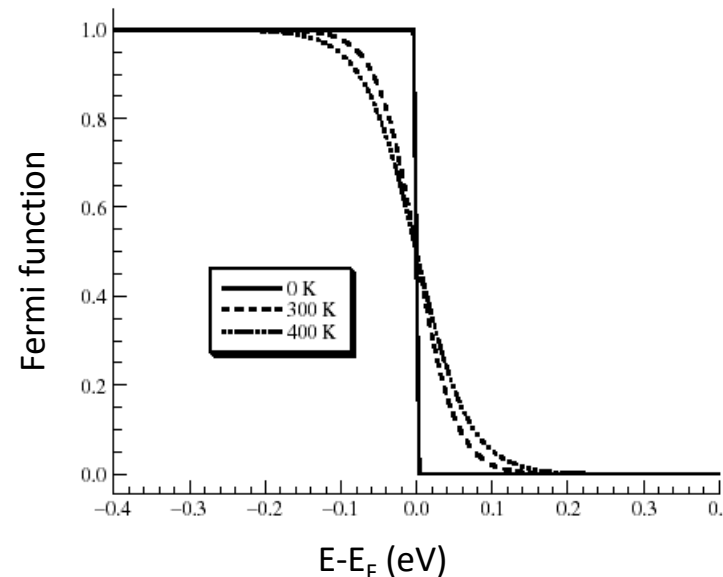
$$p_0 = N_v e^{(E_v - E_F)/kT}$$

Fermi-Dirac statistics

$$f(E) = \frac{1}{1 + e^{(E - E_F)/kT}}$$

takes into account
Pauli exclusion principle
(one fermion per state)

E_F : Fermi Energy



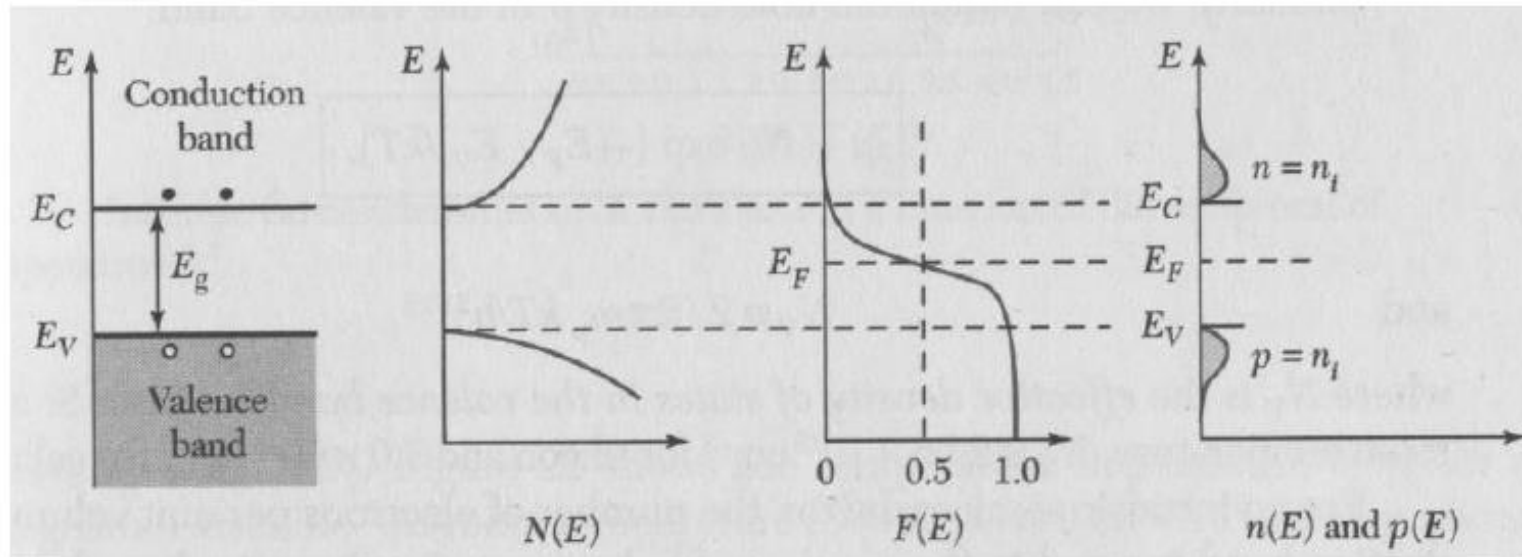
with

$$\left[\begin{aligned} N_c &= 2 \left(\frac{2\pi m_n^* kT}{h^2} \right)^{3/2} \\ N_v &= 2 \left(\frac{2\pi m_p^* kT}{h^2} \right)^{3/2} \\ E_F &= E_c + kT \ln \left(\frac{n_0}{N_c} \right) \end{aligned} \right.$$

Summary

Semiconductor band diagram

Fermi distribution function



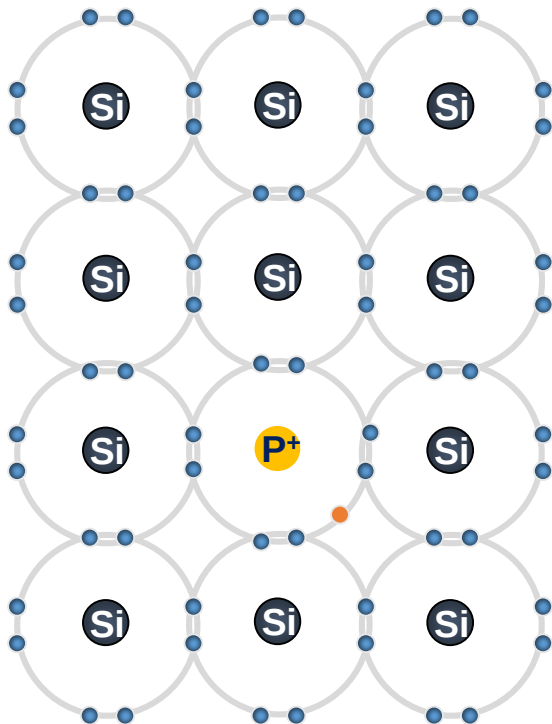
Density of states
Density of allowed states
per energy range

Carrier concentration

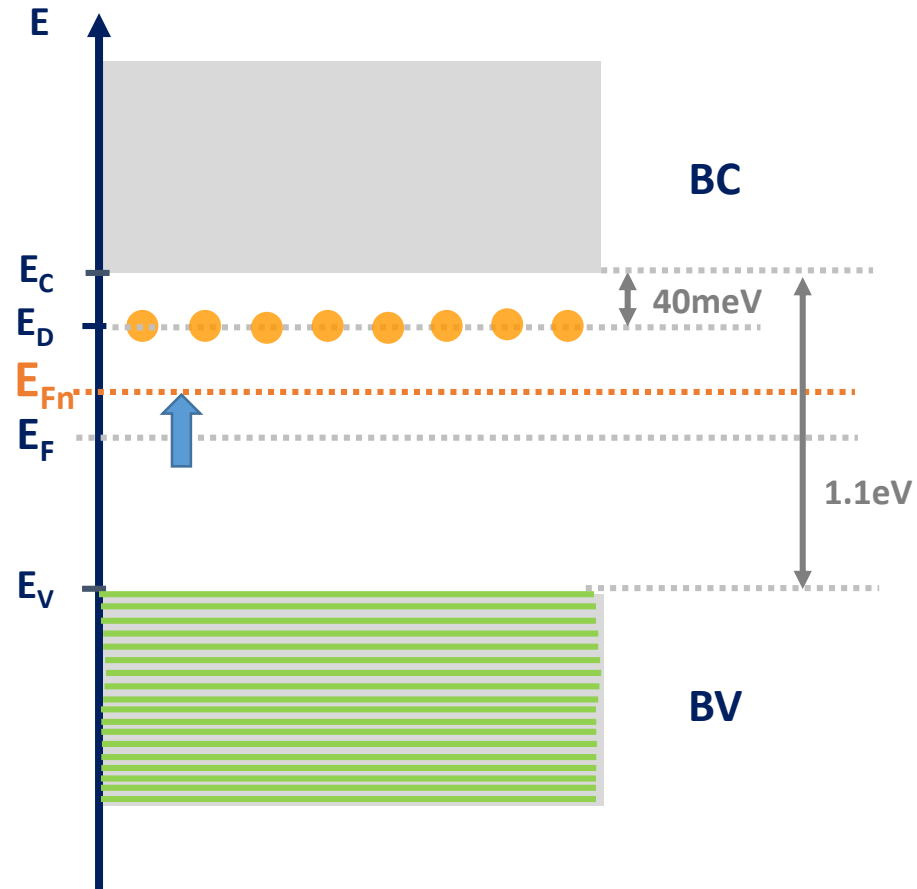
$$n = \int_{E_C}^{\infty} N(E) F(E) dE$$

Defects and Doping in SC

n-doping

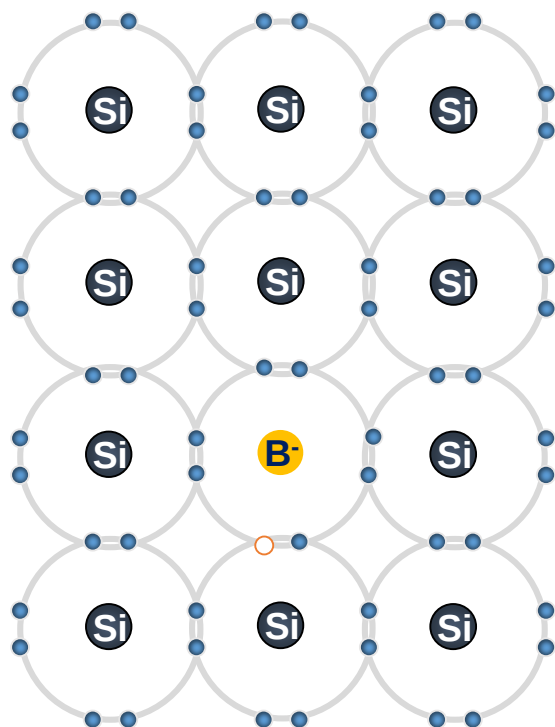


Concentration $n=N_D$ density of dopants



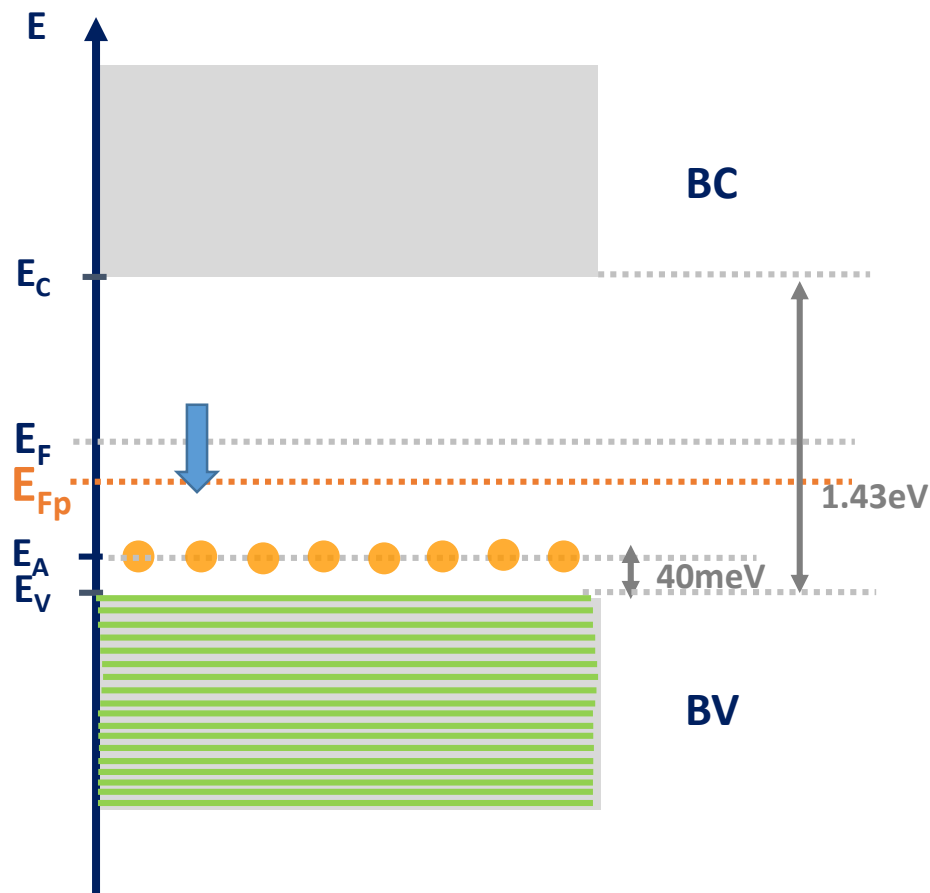
Defects and Doping in SC

p-doping



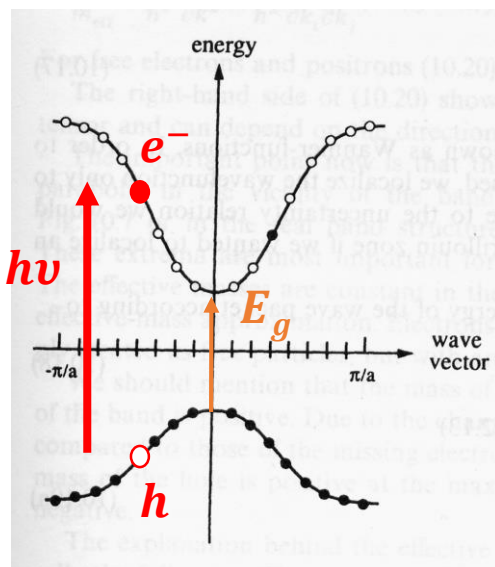
$B \text{ atom} \rightarrow B_{\text{fixed}}^- + \text{hole}$

Concentration $p=N_A$ density of dopants



Interaction of light with matter

Absorption



Incident photon

$$E = h\nu$$

$$K = \frac{2\pi}{\lambda}$$

* Energy conservation in the transition

$$E_c(k_f) = E_v(k_i) + h\nu$$

$$\Rightarrow h\nu = E_g + \frac{\hbar^2 k^2}{2} \left(\frac{1}{m_c} + \frac{1}{m_v} \right)$$

Reduced effective mass $\frac{1}{\mu} = \frac{1}{m_c} + \frac{1}{m_v}$

Ex: GaAs: $m_c \sim 0.066m_0$; $m_v \sim 0.4m_0 \rightarrow \mu \sim 0.057m_0 \sim m_c$

* Momentum conservation in the transition

$$\vec{k}_f = \vec{k}_i + \vec{K}$$

Order of magnitude:

- Photons (UV - visible - near IR): $K = 2\pi/\lambda \sim 10^4$ to 10^6 cm^{-1}
- Electrons: in the Brillouin zone, $k = 2\pi/a \sim 10^8 \text{ cm}^{-1}$

Ex: GaAs: $a \sim 2.8 \text{ \AA}$

K wave vector of photon negligible $\rightarrow$ Vertical optical transitions



Only applies to direct bandgap SC

Absorption

Optical transition (Absorption) Rate for a photon of energy $h\nu = E_c - E_v$ per time unit [s^{-1}]

$$R_{abs}(h\nu) = B_{vc} \cdot \rho_j(h\nu) \cdot f_v(E_v) \cdot (1 - f_c(E_c))$$

Einstein
coefficient

Density of
electronic
states at
energy $h\nu$

Probability that
valence state
is occupied

Probability that
conduction state
is empty

Einstein coefficient

- * Calculated from perturbation theory: 2 discrete states 'perturbed' by a light wave
- * **Perturbation:** Interaction electron $\leftrightarrow$ e.m. wave $W(t) = -q\vec{E}(\vec{r}, t)\hat{r}$ $\hat{r}$ position operator
 q electron charge
Electric dipole Hamiltonian, see Cohen Tannoudji Ch. A_{xiii}

- * **Fermi's golden rule:**

$$B_{vc} = \frac{\rho q^2}{2\hbar} \cdot \left| \langle Y_c(r) | E \cdot \hat{r} | Y_v(r) \rangle \right|^2 \times d(E_c - E_v - \hbar\omega)$$

$\swarrow \quad \searrow$
 Coupling of Ψ_v to Ψ_c
 through $W(t)$

energy conservation

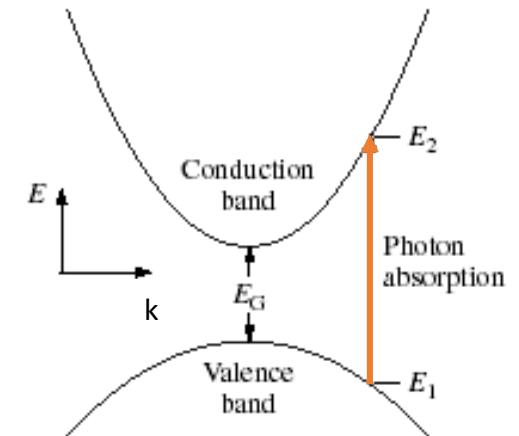
- * **Dipole matrix element:** gives strength of light-matter interaction

Absorption

Direct Absorption

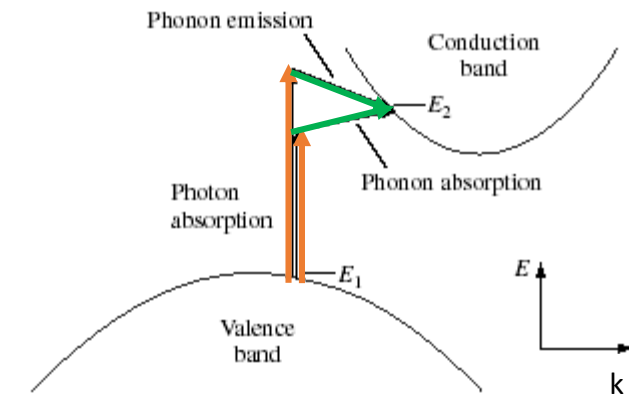
$$\alpha(h\nu) \approx A^* \cdot (h\nu - E_g)^{1/2}$$

$$h\nu = E_g + \frac{\hbar^2 k^2}{2} \left(\frac{1}{m_c} + \frac{1}{m_v} \right)$$

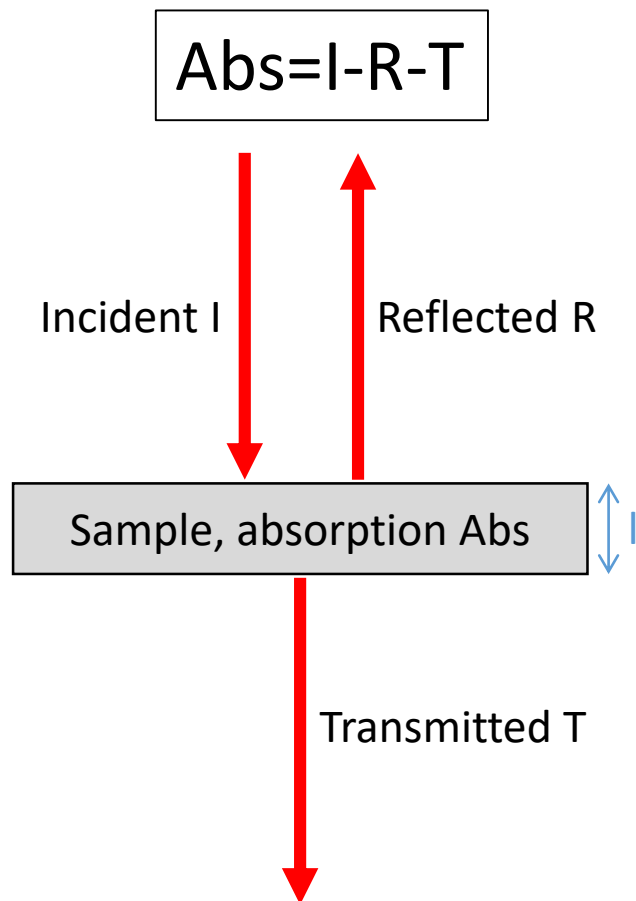


Indirect Absorption (phonons)

$$\alpha(h\nu) \approx \frac{A(h\nu - E_g + E_{\text{phonon}})^2}{e^{E_{\text{phonon}}/kT} - 1} + \frac{A(h\nu - E_g - E_{\text{phonon}})^2}{e^{-E_{\text{phonon}}/kT} - 1}$$



Absorption measurement



Absorption coeff measured by transmission of a thin platelet sample of thickness l

Beer's law $I(z) = I_0 e^{-\alpha z}$

Transmissivity of absorbing medium of thickness l

$$T = (1 - R)^2 e^{-\alpha l}$$

and

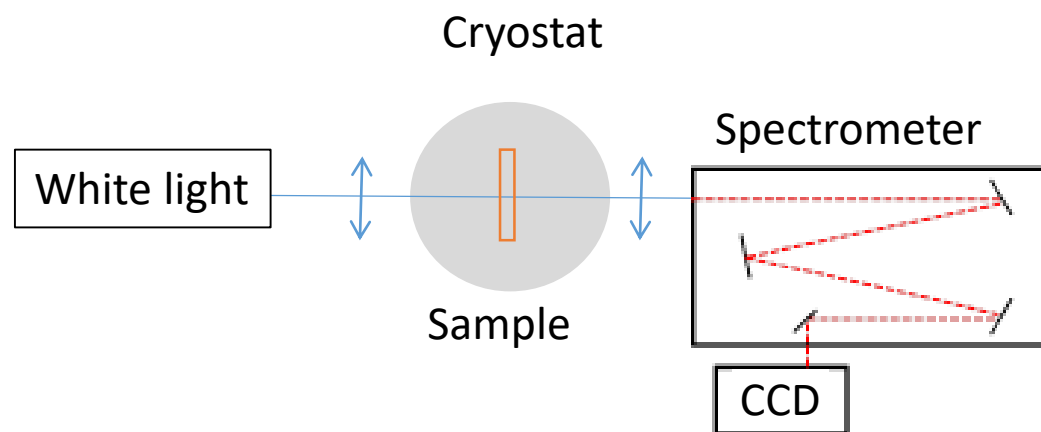
$$T_{\text{measured}} = \frac{I_{w/\text{Sample}}}{I_{w/o\text{Sample}}}$$

$$\alpha = \frac{1}{L} \ln \left(\frac{I_{w/o\text{Sample}}}{I_{w/\text{Sample}}} \right)$$

Requires transparent substrate
Requires two collection paths
Requires a reference

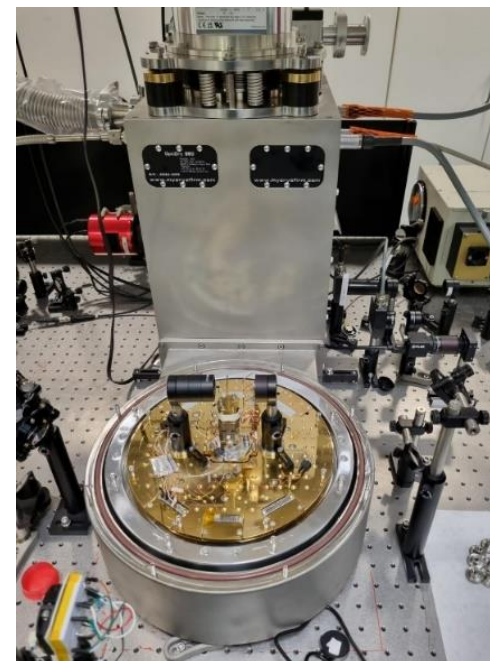
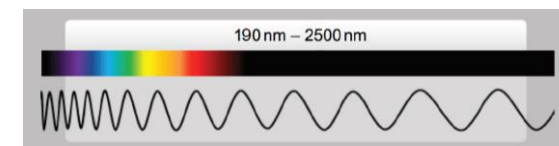
Absorption measurement

Experimental setup



- Low excitation power
- Semiconductor choice:
 → imposes choice of white lamp, mirrors/lenses, grating and detector
- Temperature control: need for a cryostat
- Measure a reference

Cryostat

Broad wavelength range
Light source

Spectrometer

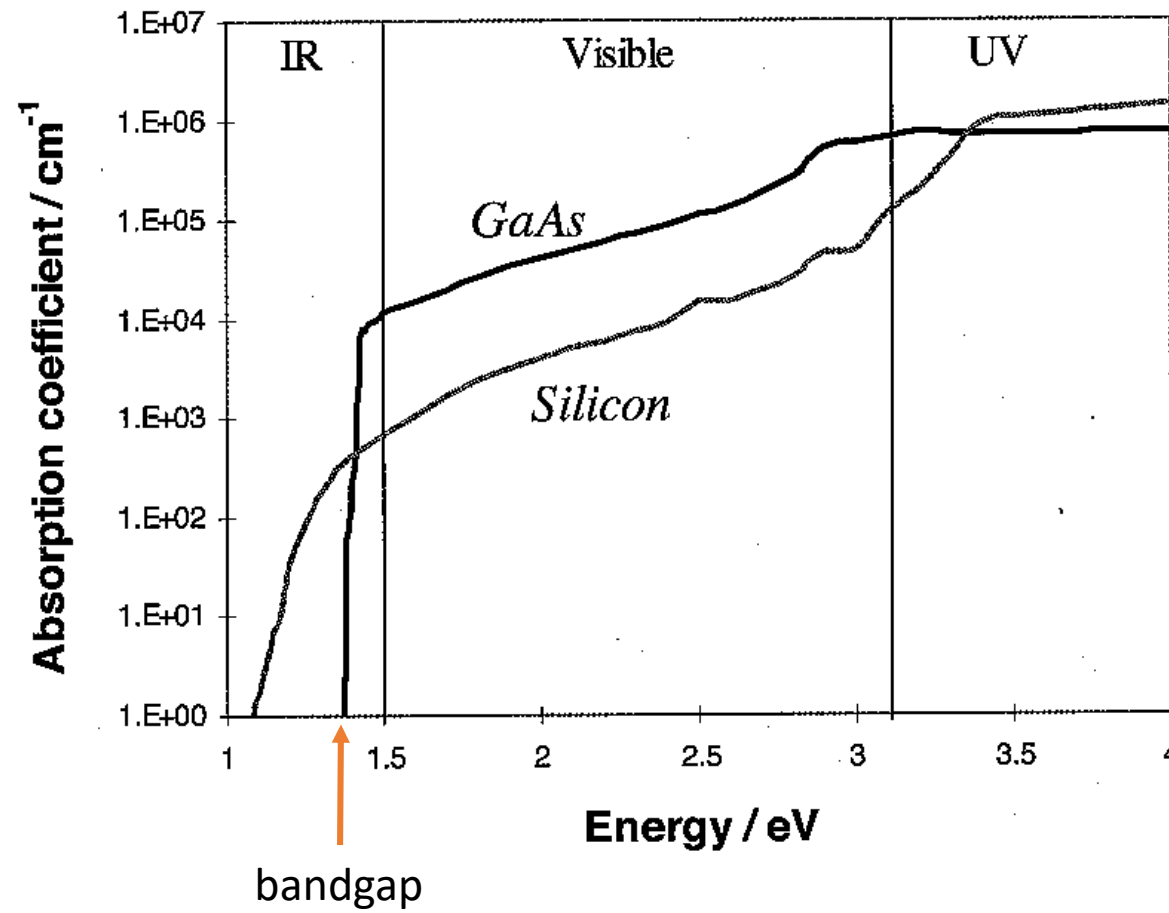


UV: PM

VIS: Si CCD or PD

NIR: InGaAs PD

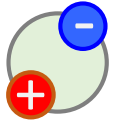
Absorption measurement



In SC

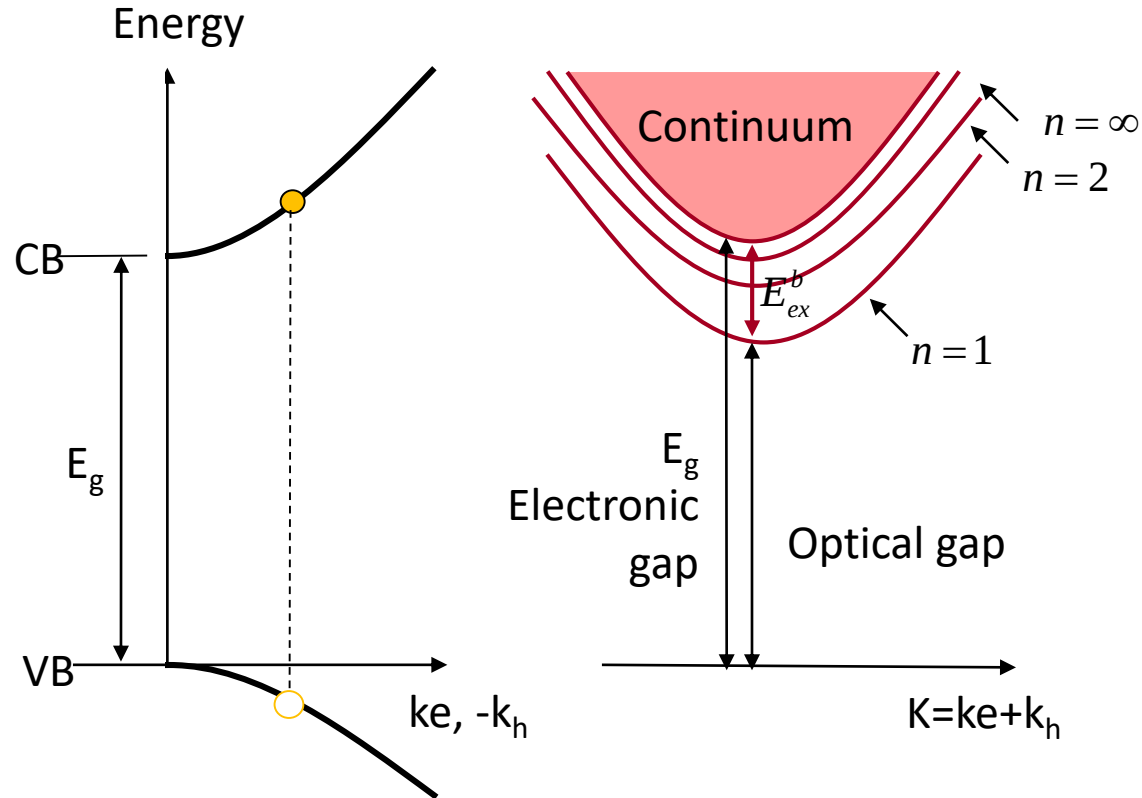
$$\alpha_{typ} \approx 10^{-4} \text{ cm}^{-1}$$

Excitons



Photogenerated electron AND hole are linked by coulombic interaction

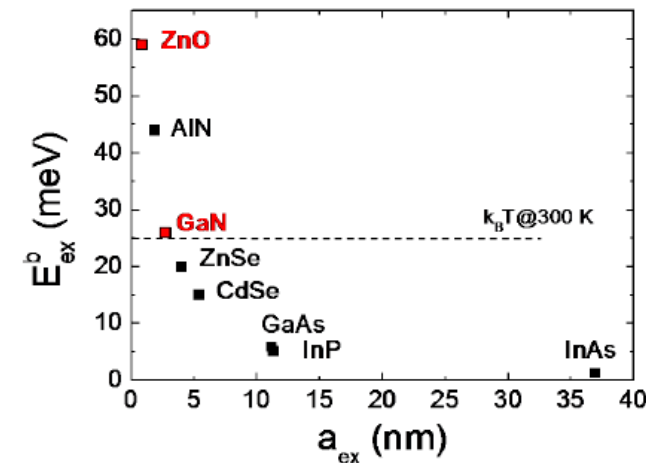
➡ **exciton**



Discrete exciton energy: $E_{X,n}(k=0) = E_g - \frac{R_y^2}{n^2}$

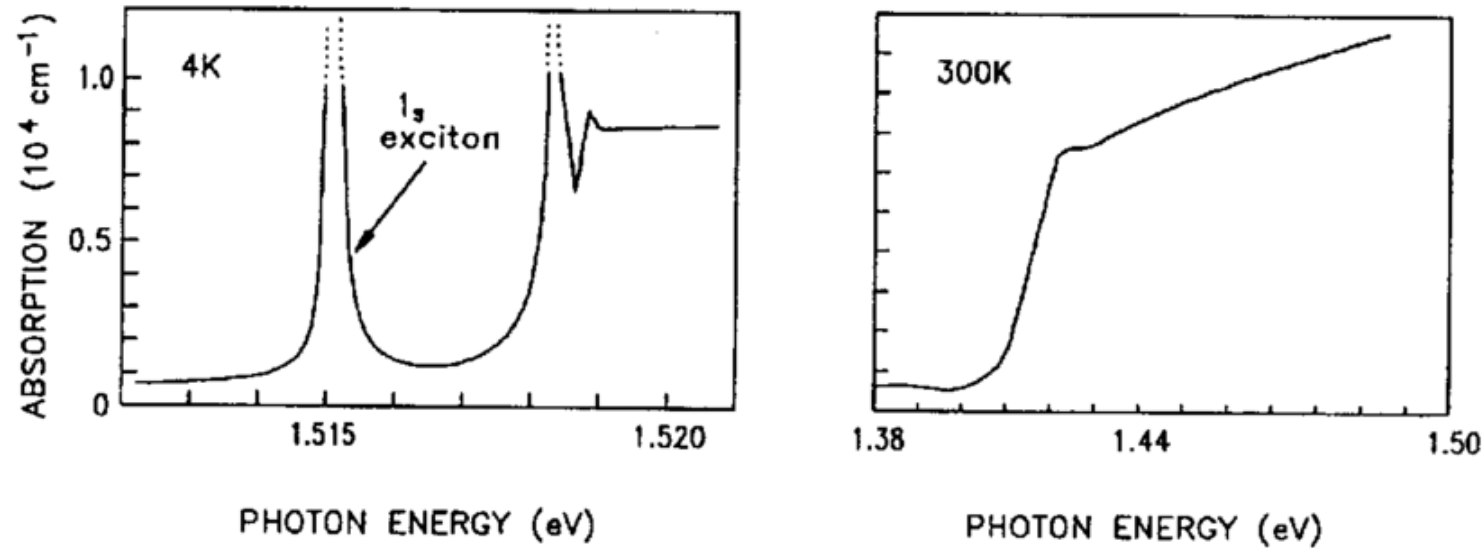
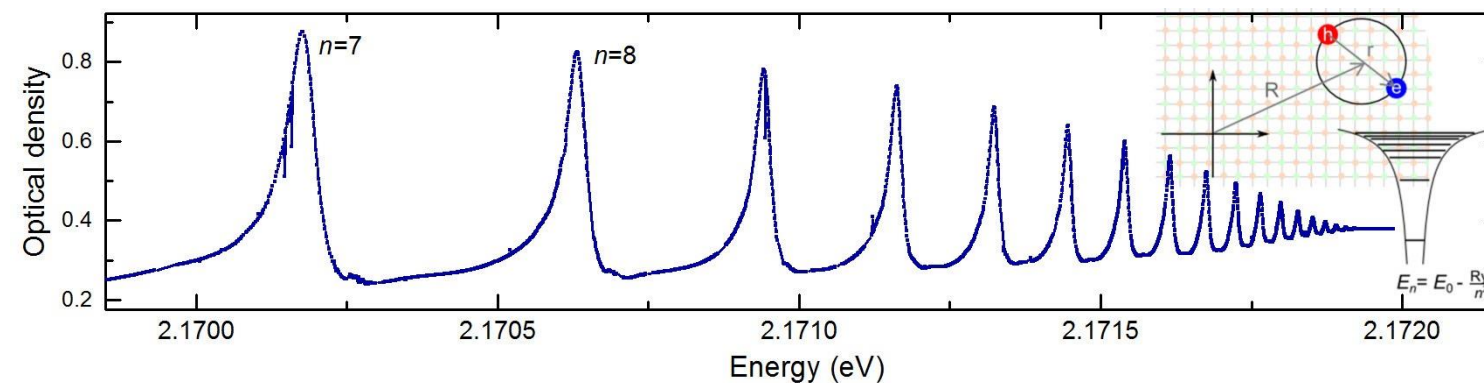
Exciton binding energy

$$E_b^x = \frac{e^2}{2\epsilon_0 a_B} \quad \text{Exciton Bohr radius}$$



Excitons

GaAs

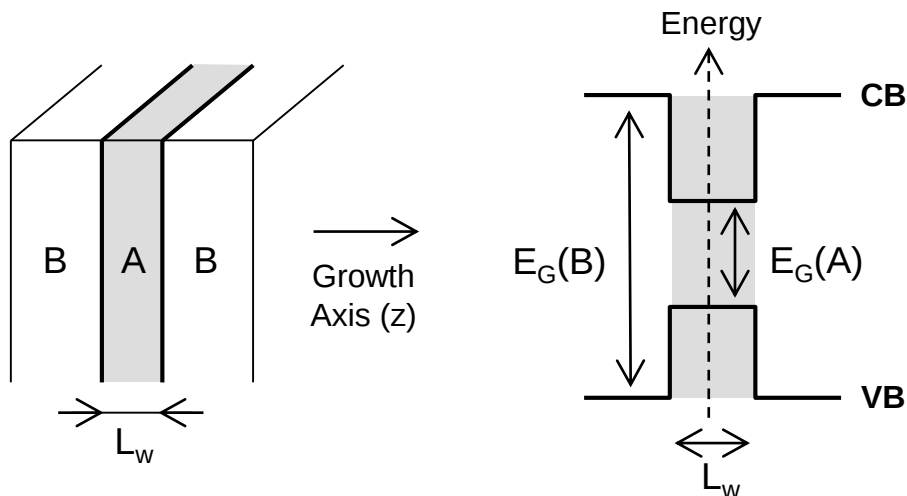
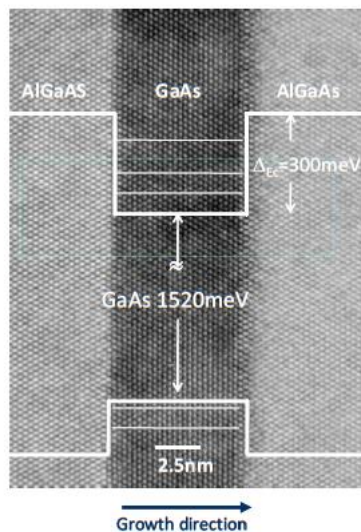
Cu₂O

Kazimierczuk, *et al.* Nature 514.7522, 343 (2014)

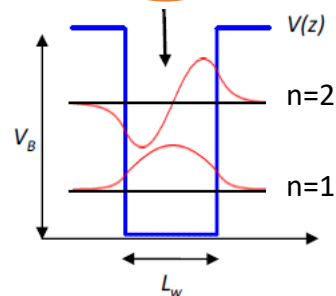
Dimensionality

Quantum wells / 2D

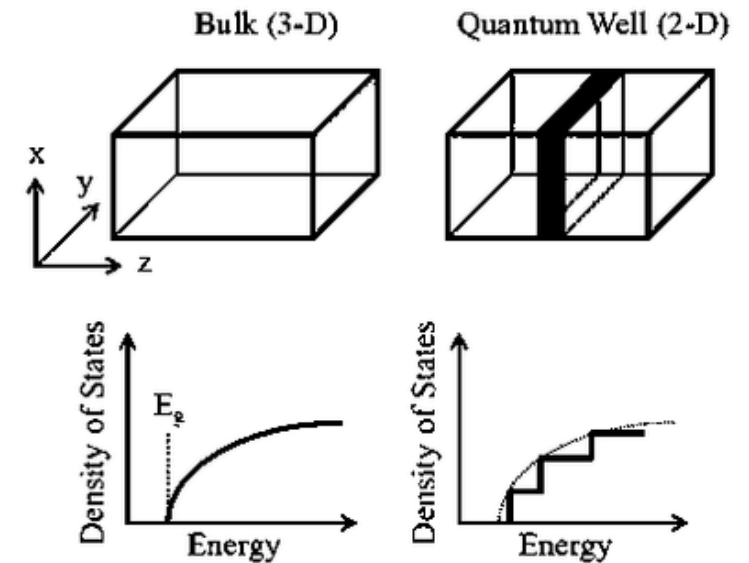
Typ. $L_w < \text{few } 1\text{-}10 \text{ nm}$



$$\Psi_i(\vec{r}) = \phi_c(z) e^{-i(k_x x + k_y y)}$$



ϕ_z z confined
 $e^{-ik_x x}$ x free
 $e^{-ik_y y}$ y free



$$E_{n,k} = E_n + \frac{\hbar^2 k_{\parallel}^2}{2m^*}$$

Optical transitions in 2D

Same approach as in 3D: time dependent perturbation theory

calculate absorption rate $R_{v,j \rightarrow c,i}^{abs}(h\nu)$ of photons of energy $h\nu$ between heavy holes of band j in the VB $\rightarrow$ electrons of band i in the CB using Fermi's golden rule:

$$R_{v,j \rightarrow c,i}^{abs}(h\nu) \propto [\text{dipole element}] \times [\text{joint density of state}] \times [\text{occupation probability}]$$

$$R_{v,j \rightarrow c,i}^{abs}(h\nu) = \frac{\pi}{2\hbar} \left| \langle \Psi_c^i | q \vec{E}_{em} \hat{r} | \Psi_v^j \rangle \right|^2 \cdot \frac{\mu}{\pi \hbar^2} \cdot f_v^j(E_v) \cdot [1 - f_c^i(E_c)] \quad [\text{s}^{-1} \cdot \text{cm}^{-2}]$$

for $k=0$ $Y_c^i(k_{//} = 0) = F_c^i(z)u_c(r)$

envelope function $\left| \langle Y_c^i | q \vec{E}_{em} \hat{r} | Y_v^j \rangle \right|^2 \mu \left| \langle u_c | \vec{E} \times \hat{r} | u_v \rangle \right|^2 \cdot \left| \langle F_c^i | F_v^j \rangle \right|^2$



Selection rule n° 2

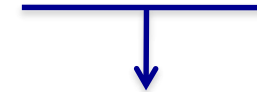


Selection rule n° 1

Optical transitions in 2D

envelope function

$$\left| \left\langle Y_c^i \left| q \vec{E}_{em} \hat{r} \right| Y_v^j \right\rangle \right|^2 \propto \left| \left\langle u_c \left| \vec{E} \times \hat{r} \right| u_v \right\rangle \right|^2 \cdot \left| \left\langle F_c^i \left| F_v^j \right\rangle \right|^2$$



Selection rule n° 1

Selection rule n° 2:

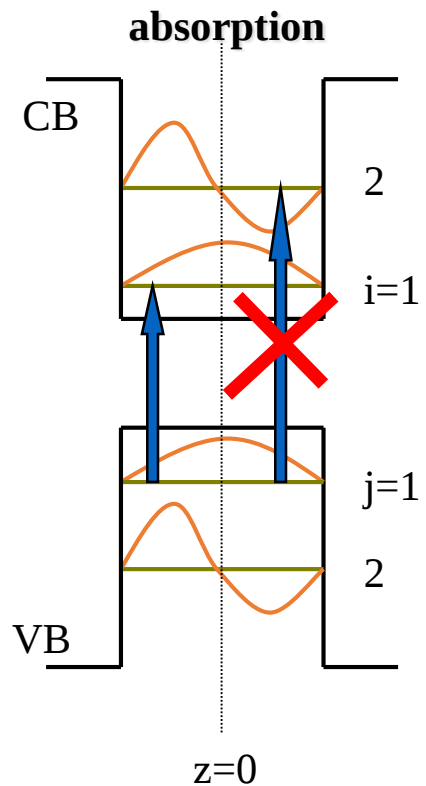
overlap integral of envelope functions

Polarization of emission

if i (or j) is even, the envelope function is even,
if i (or j) is odd, the envelope function is odd

Hence, if $(i+j)$ is odd $\Rightarrow \left\langle F_c^i \left| F_v^j \right\rangle = 0$

Transition is forbidden

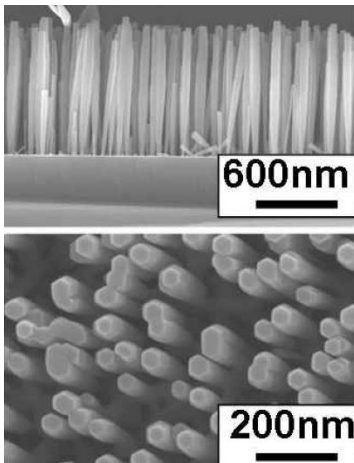
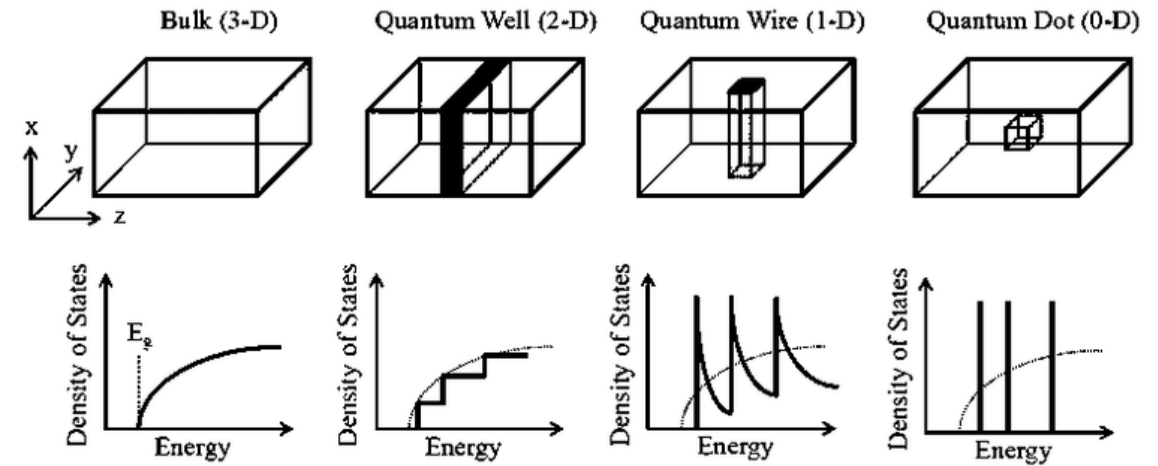


Dimensionality

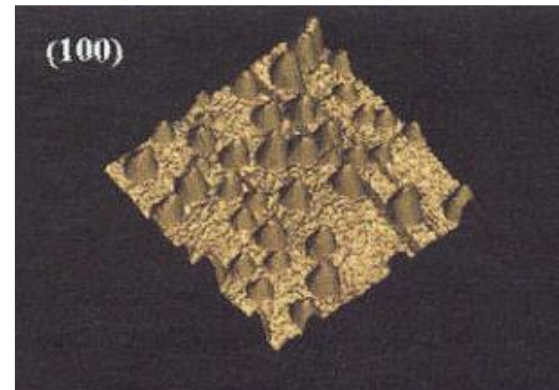
Quantum wells / 2D $\Psi_i(\vec{r}) = \phi_c(z)e^{-i(k_x x + k_y y)}$

1D $\Psi_i(\vec{r}) = \phi_c(y, z)e^{-ik_x x}$

Quantum dots $\Psi_i(\vec{r}) = \phi_c(x, y, z)$



GaN nanowires



InAs Quantum dots

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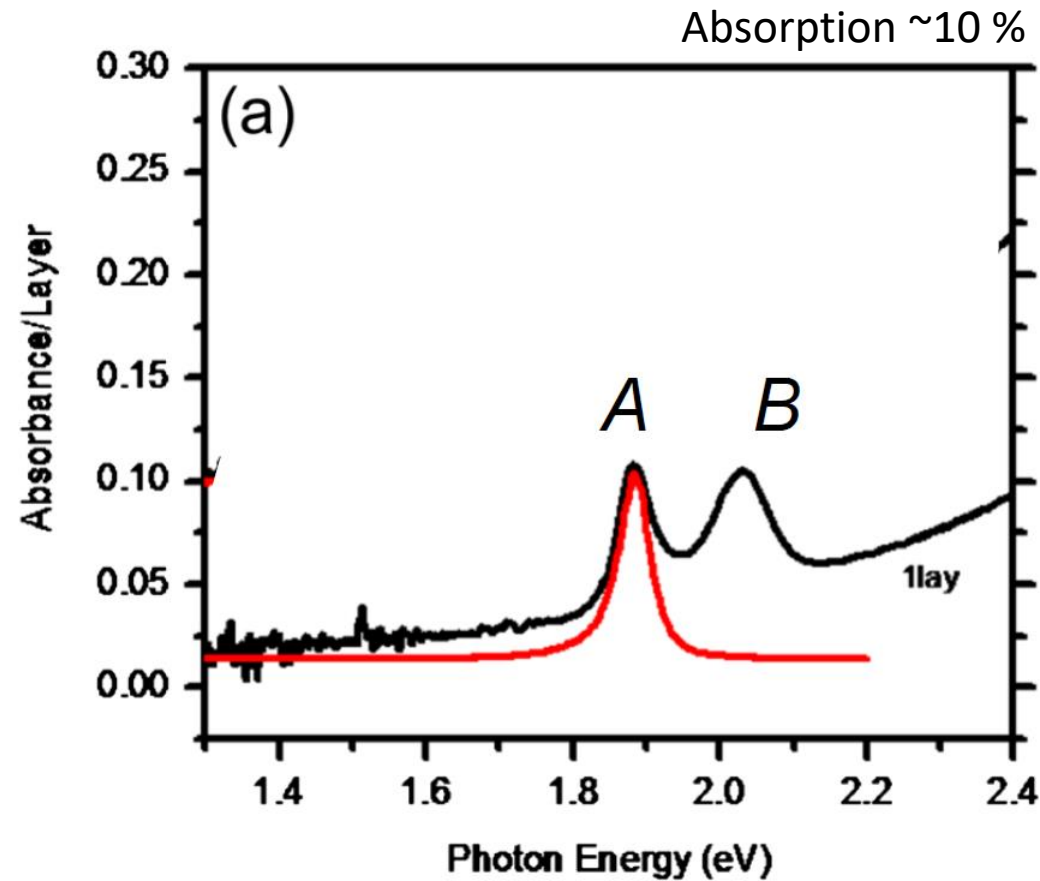
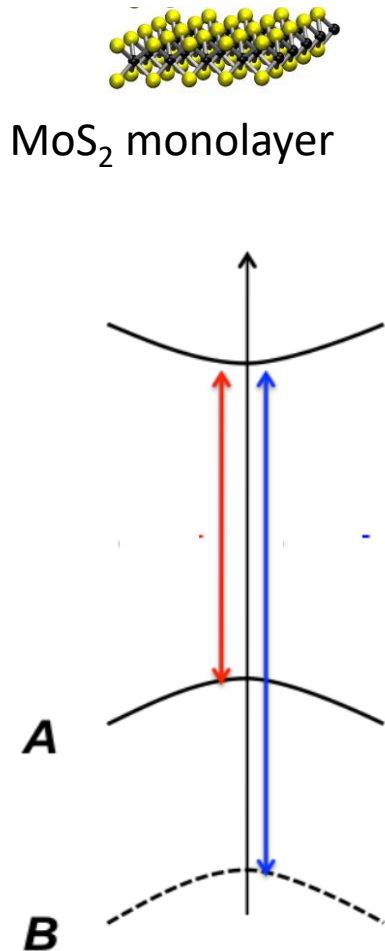
Electric/Optical characterization of SC structures

- * Electroluminescence (EL)
- * Photocurrent (PC)

Raman spectroscopy

Absorption

Ex:TMDC



Phys. Rev. Lett. 105, 136805 (2010)

Photoemission Spectroscopy (UPS & ARPES)

Photoemission

* based on the **photoelectric effect**:

Photons are absorbed and electrons are ejected above vacuum level $h\nu > E_F + \phi$

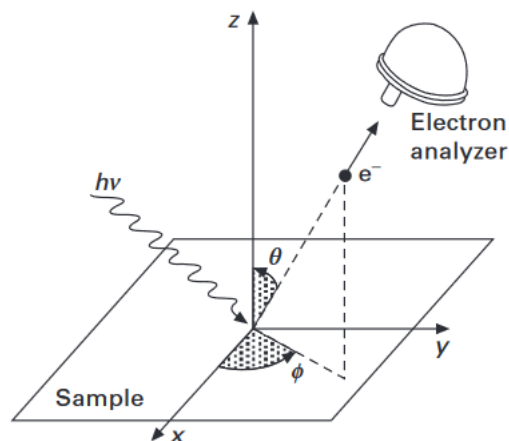
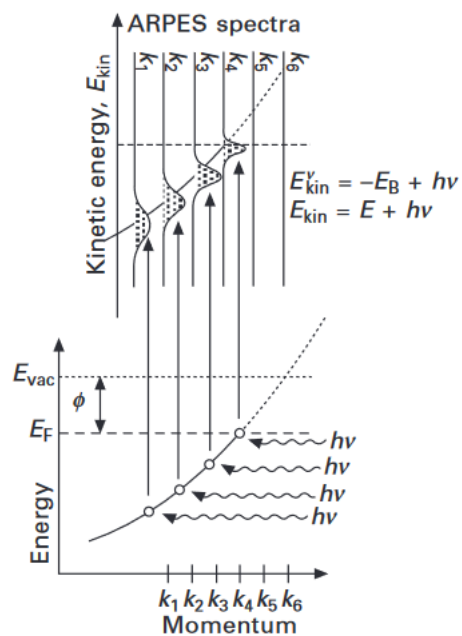
* need for highly energetic photon source (typ. UV rays or He, Ar, Xe gas discharge lamp)

* Only for surface characterization: short escape depth of the photoelectrons to the surface (<10nm)

* Direct way to measure the electronic density of states near E_F (for SC: electrons in top VB)

UPS : UV Photoemission Spectroscopy

ARPES: Angle resolved Photoemission Spectroscopy



Energy conservation $\rightarrow$ measurement of electron energy in the VB band

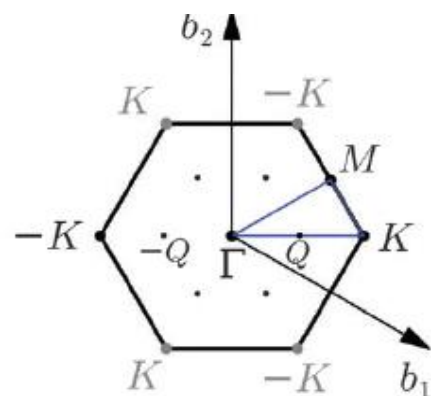
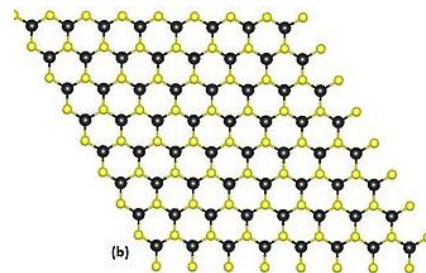
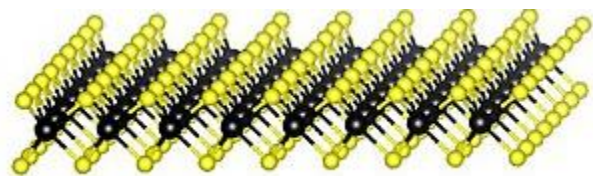
Momentum conservation $\rightarrow$ measurement of quasimomentum of the electron in the initial state given by the measurement of the photoelectron momentum

In-situ Charac. of thin film Growth 55 (2011)
Nature Reviews 2, 54 (2022)

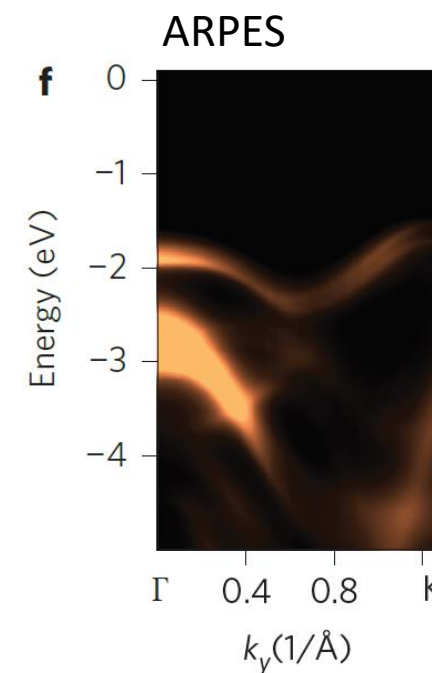
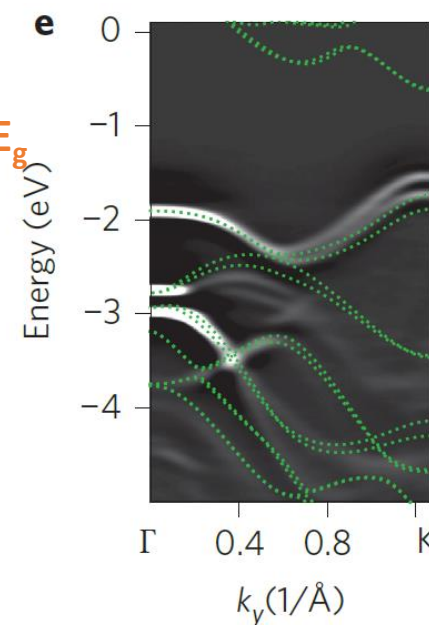
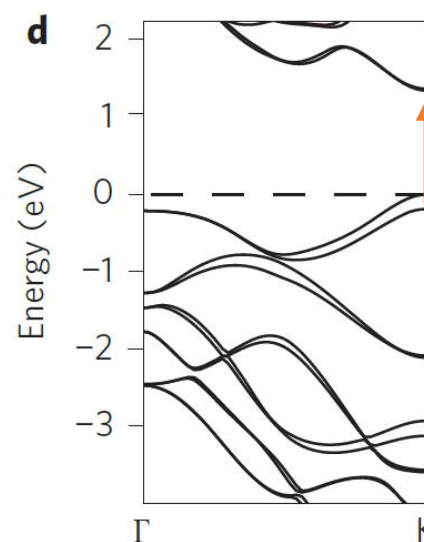
Photoemission Spectroscopy (ARPES)

Band structure of TMDC

MoSe₂ - ML



First Brillouin zone

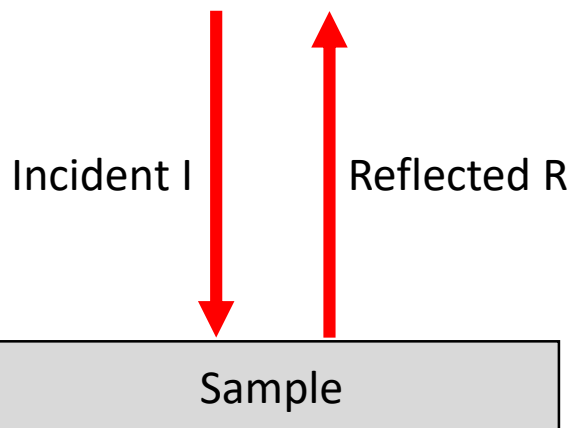


Tight-binding / k.p / DFT-GW calculations

PRB **79**, 115409 (2009), [Nancy](#)
 PRB **85**, 205302 (2012), [Cleveland](#)
 PRB **86**, 115409 (2012), [Massachusetts](#)
 2D Materials **2**, 22001 (2015), [Konstanz](#),
 ...

Nature Nanotechnology **9**, 111 (2013)

Reflectivity



Probes vertical transitions in k-space

Probes the dielectric function

$$\varepsilon(\omega)^2 = (\varepsilon_1(\omega) + i\varepsilon_2(\omega))^2 = n(\omega) + i\kappa$$

Intensity ruled by selection rules and oscillator strength of the transition

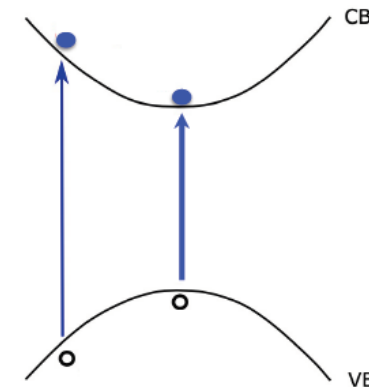
$$\varepsilon(\omega) = 1 + \sum_k \frac{f_k}{\omega_{0,k}^2 - \omega^2 - i\omega\gamma_k}$$

ω : incident light angular frequency

ω_0 : angular frequency of the oscillator

f_k : oscillator strength of oscillator and electric field of light

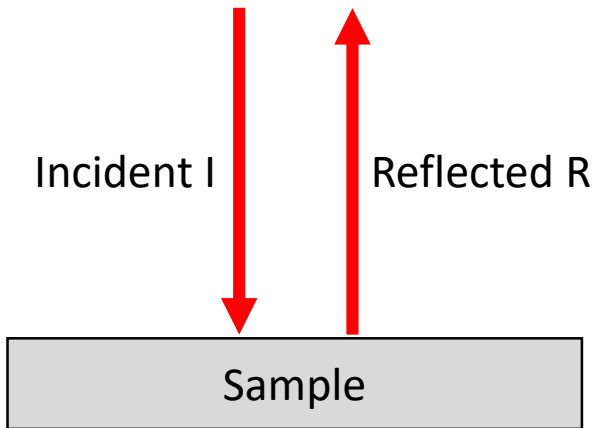
γ_k : damping factor



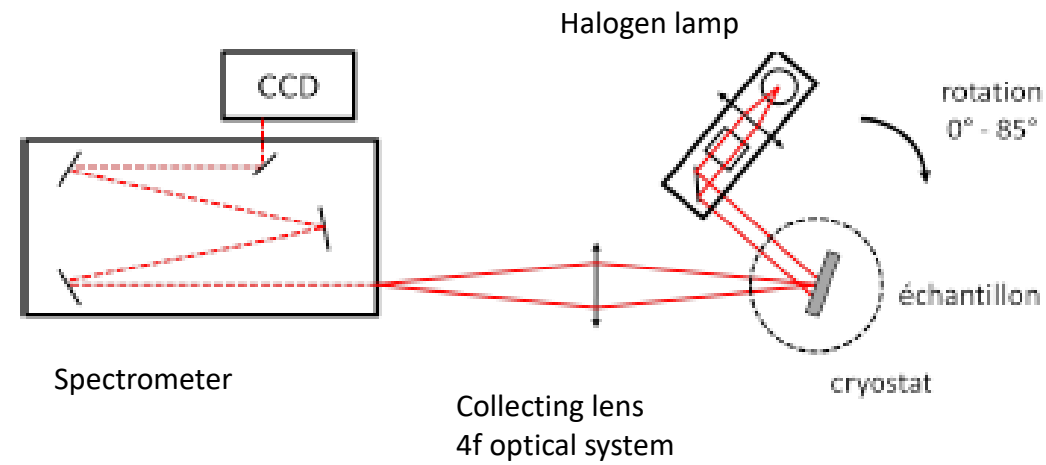
- ➡ Requires a model (transfer matrix for instance) or a transformation (Kramers Kronig) to extract quantitative parameters (energy of the transition, oscillator strength, linewidth,...)
- ➡ Requires the precise measurement of the optical index of the sample (measured by ellipsometry for ex)

Reflectivity measurement

Experimental setup



Angle-resolved reflectivity setup



Source: white lamp (Halogen, Xe,...)

Reflected beam:

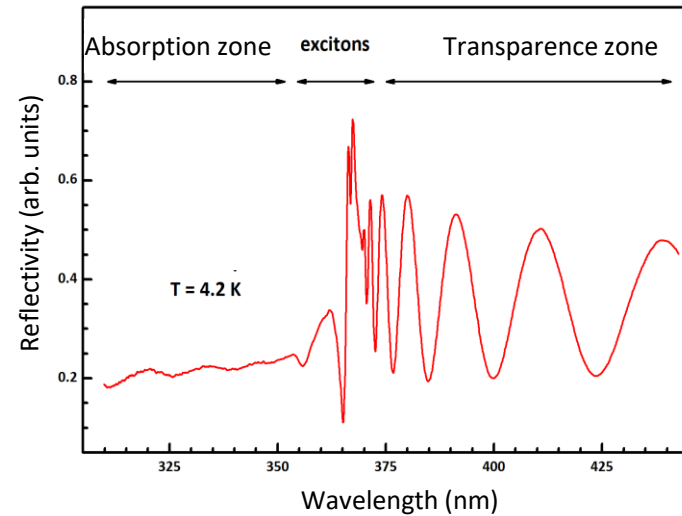
spectrally resolved using spectrometer

intensity recorded using CCD ou sCMOS camera

$$\text{Reflectivity } R = \frac{I_{w/sample}}{I_{white lamp}}$$

Reflectivity measurement

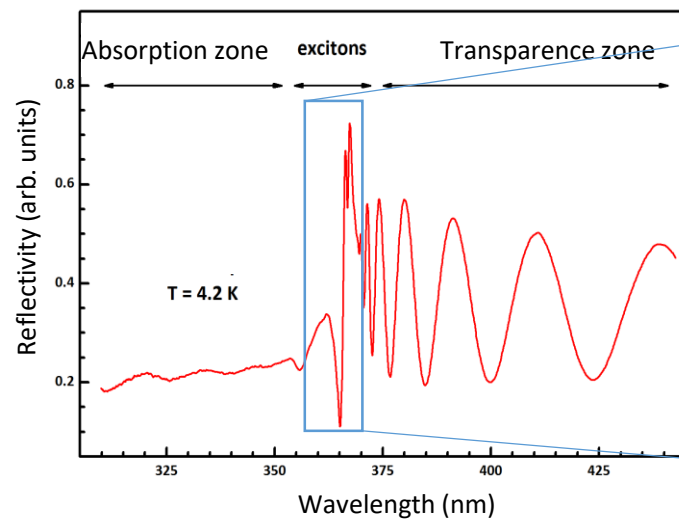
Ex: ZnO epilayer



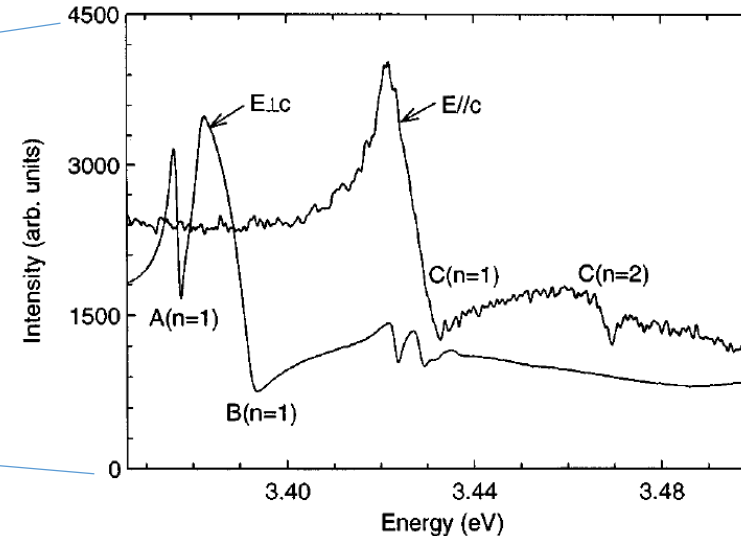
François Médard PhD thesis (2010)

Reflectivity measurement

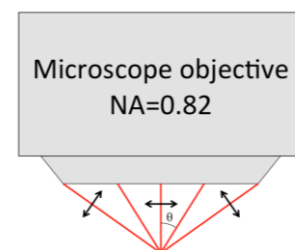
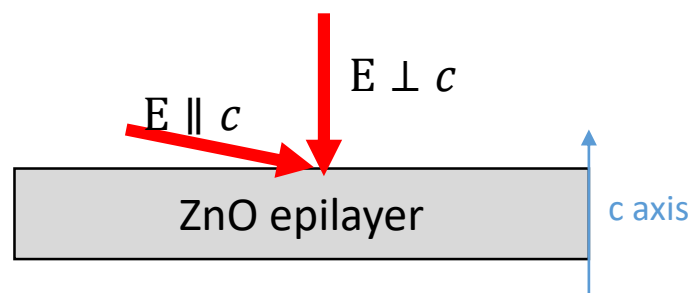
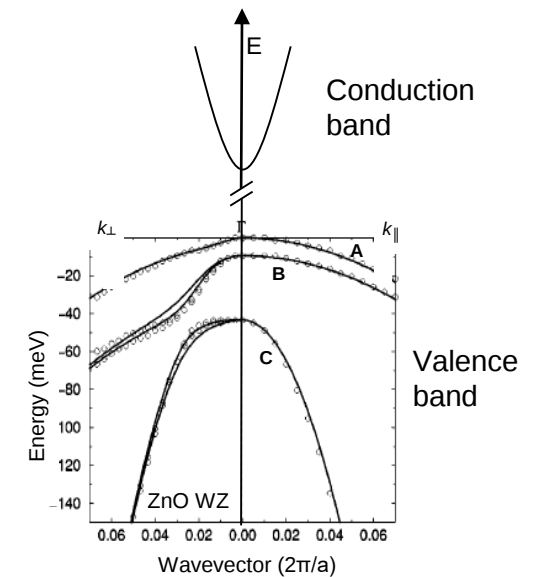
Ex: ZnO epilayer



François Médard PhD thesis (2010)



Reynolds, PRB 60, 2340 (1999)

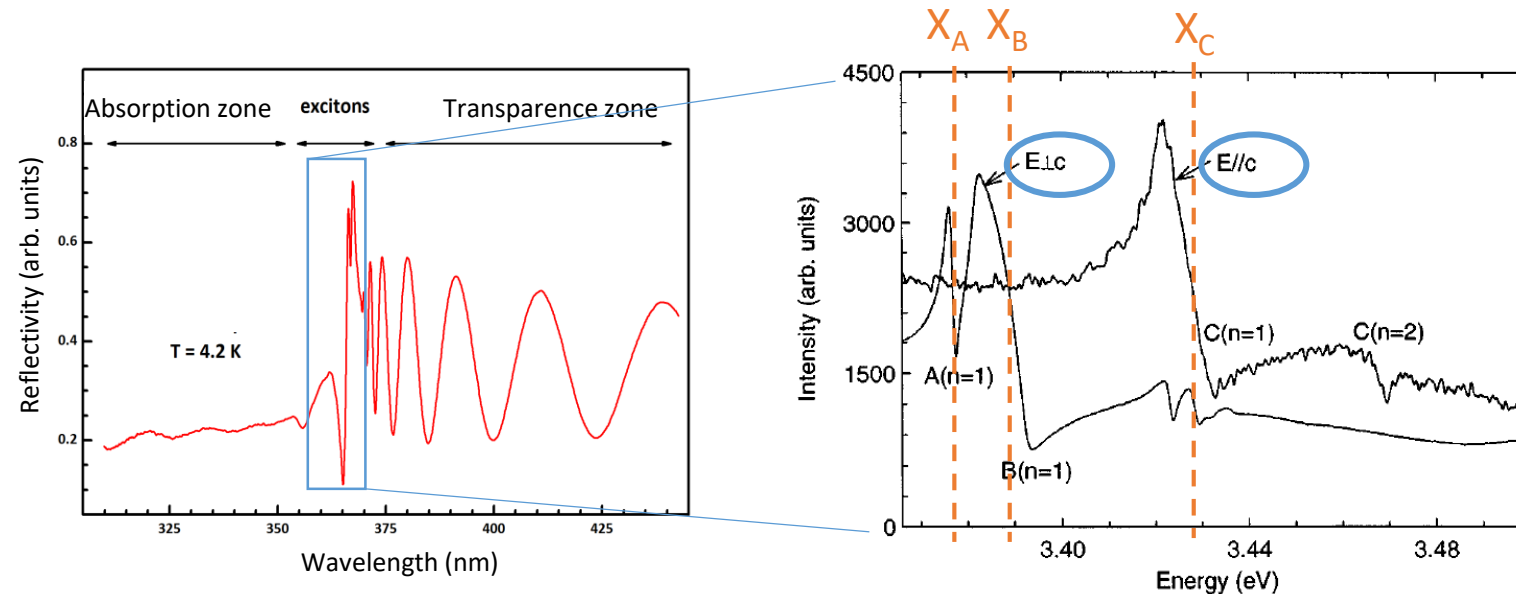


Selection rules:

A & B excitons mainly coupled to $E \perp c$
C excitons mainly coupled to $E \parallel c$

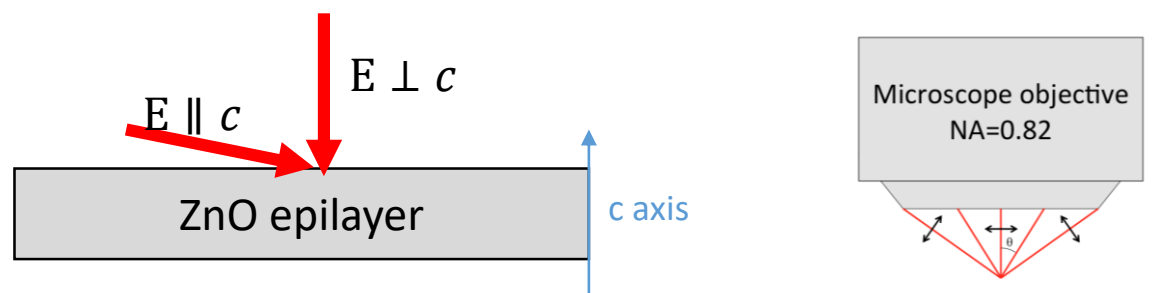
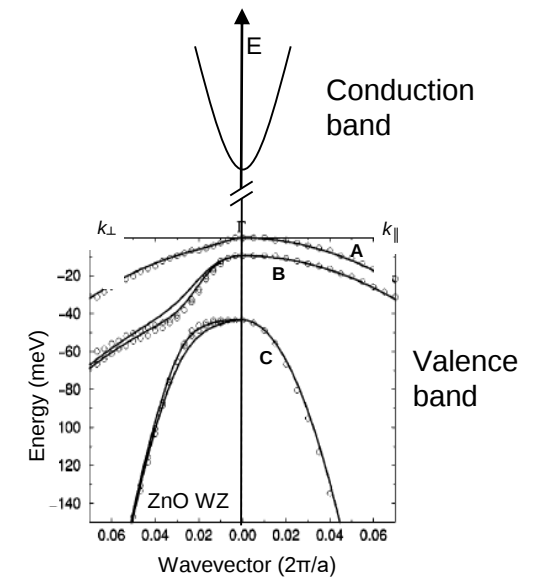
Reflectivity measurement

Ex: ZnO epilayer



François Médard PhD thesis (2010)

Reynolds, PRB 60, 2340 (1999)

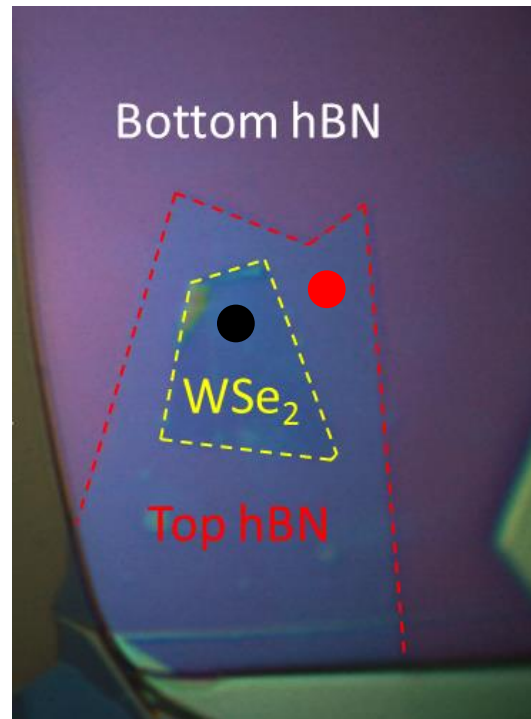
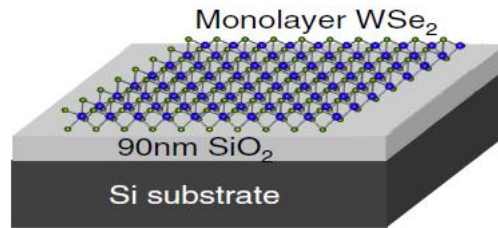


Selection rules:

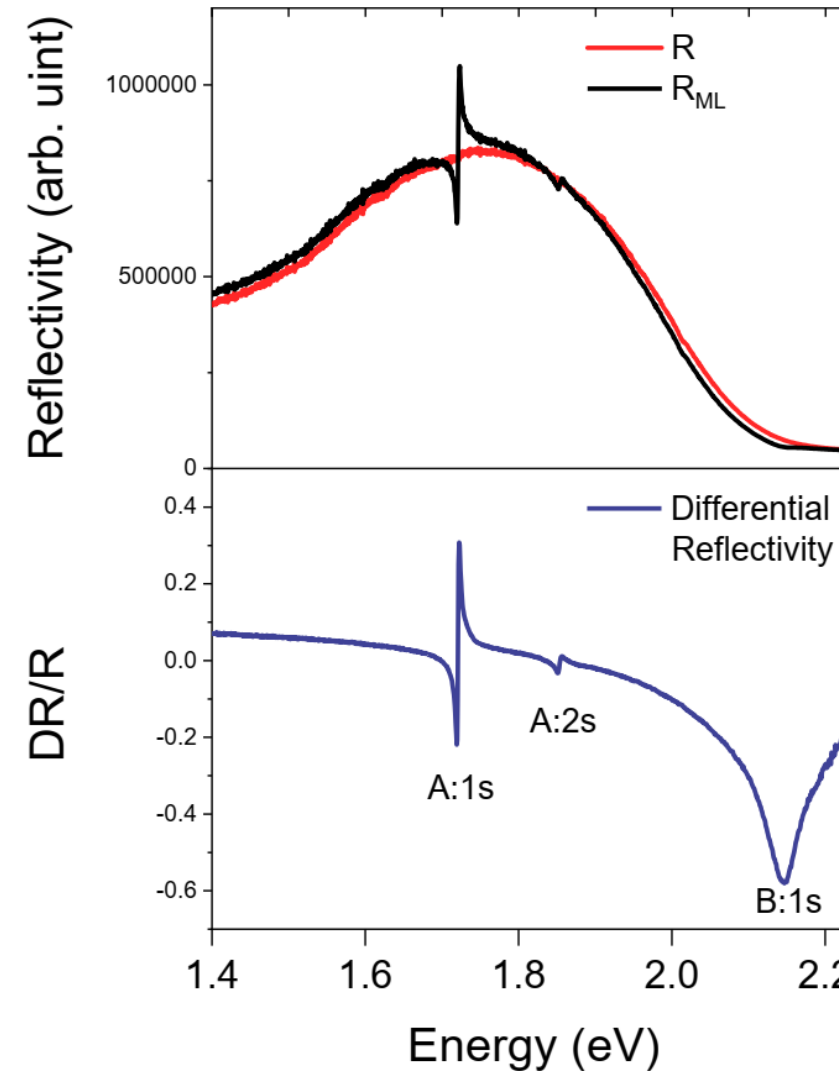
A & B excitons mainly coupled to $E \perp c$
C excitons mainly coupled to $E \parallel c$

Reflectivity

Ex: WSe₂ monolayer

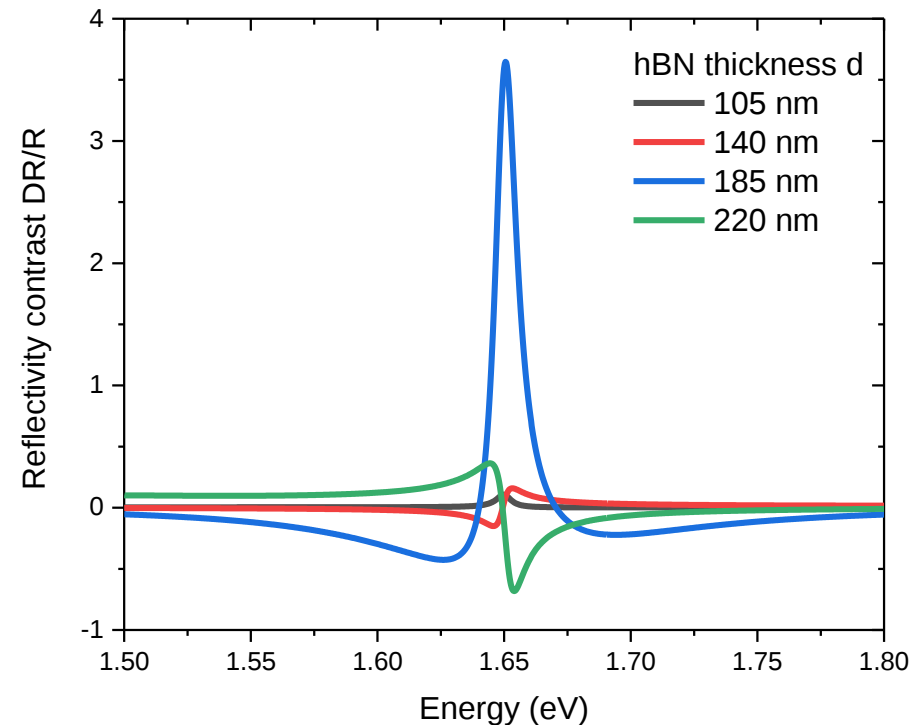
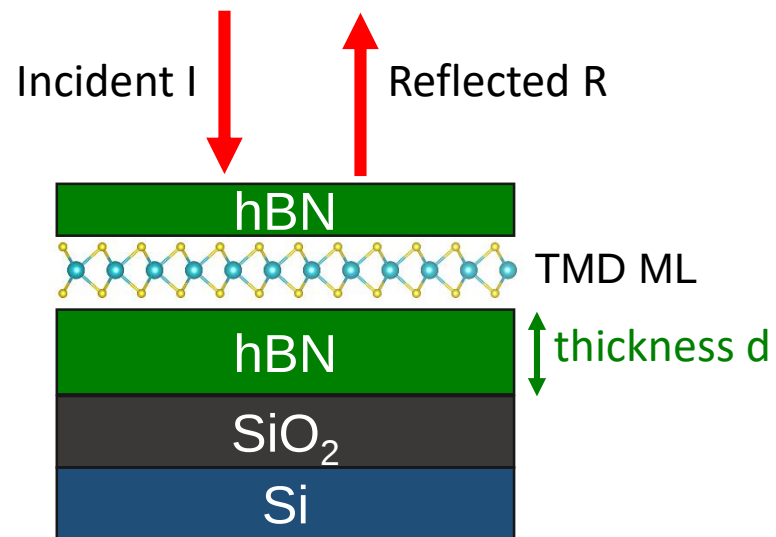


Reflectivity Contrast: $DR/R = (R_{ML} - R)/R$

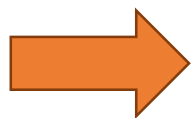


Reflectivity

Ex: WSe₂ monolayer

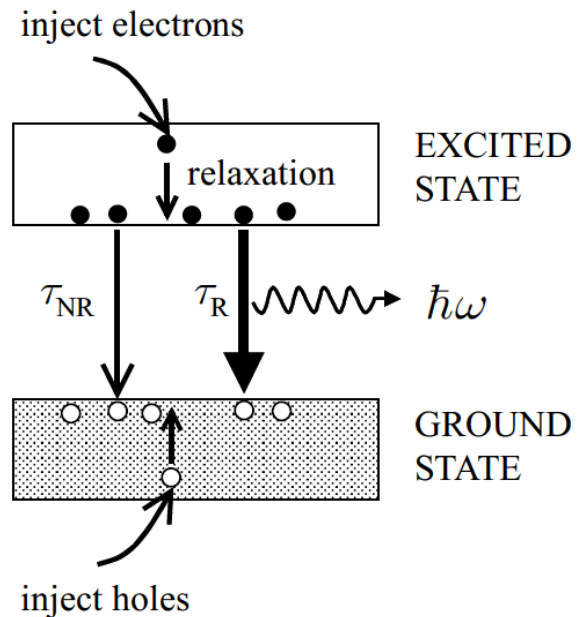


Amplitude and shape of reflectivity contrast depends on thicknesses of dielectric layers



Use of transfer matrix model needed to extract quantitative parameters (energy of the transition, oscillator strength, linewidth)

Luminescence



Probes how **excess energy is released by emission of light**

Photoluminescence: the emission of light after absorbing a photon of higher energy

Electroluminescence: the emission of light caused by running an electrical current through the material

... And others (**Cathodo-**, thermo-, chemi- etc)

$$\Phi(E) = |M|^2 \cdot g(E). \quad \text{level occupancy factor}$$

Matrix of light-mater interaction

Luminescence

Radiative recombination :

At equilibrium (no excitation – thermal emission)

$$n^0 = \int_{E_g}^{\infty} g_c(E) f_e(E) dE = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$

$$p^0 = \int_{E_g}^{\infty} g_v(E) f_h(E) dE = N_v \exp\left(-\frac{E_F - E_v}{kT}\right)$$

Fermi level

$$E_F = E_c + kT \ln\left(\frac{n}{N_c}\right) \quad n_i^2 = n^0 p^0$$

Out of equilibrium (under excitation) - excess carrier $\Delta n = \Delta p$

$$n = n^0 + \Delta n = N_c \exp\left(\frac{E_{Fn} - E_c}{kT}\right)$$

$$p = p^0 + \Delta p = N_v \exp\left(-\frac{E_{Fp} - E_v}{kT}\right)$$

Quasi-Fermi levels

$$\Delta\mu = qV = E_{Fn} - E_{Fp}$$

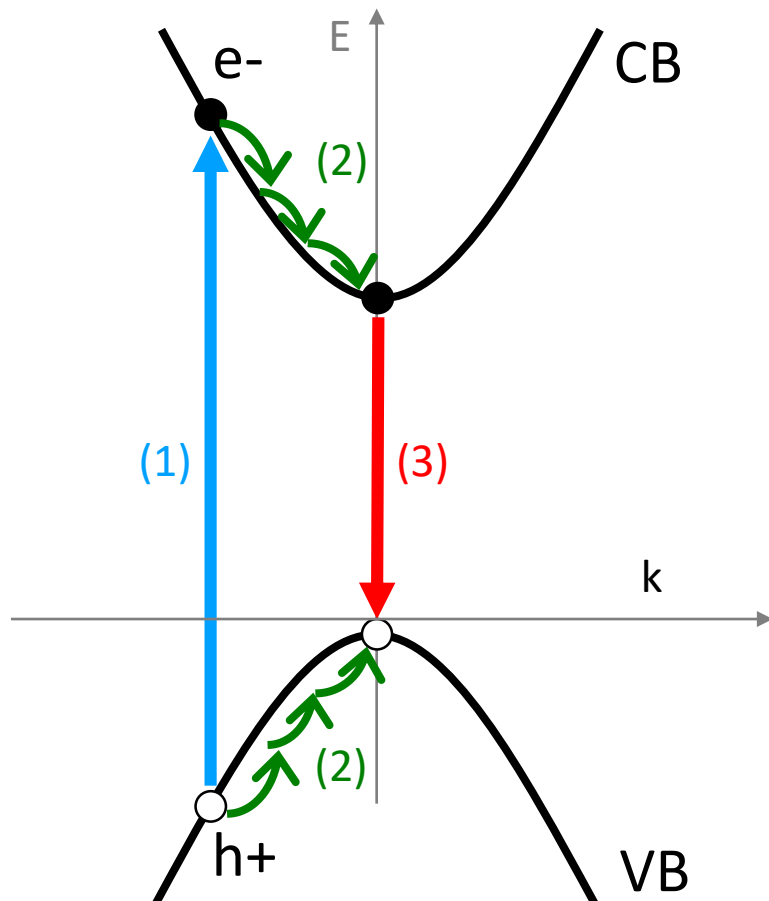
Simplified expression for recombination (out of equilibrium)

$$R = Bnp = R^0 \exp\left(\frac{E_{Fn} - E_{Fp}}{kT}\right)$$

Bimolecular recombination

Photoluminescence (PL)

a powerful tool to understand electronic structure of semiconductors



(1) **absorption**

How efficient is absorption?
Depend on which states is excited

(2) **relaxation**

Many possible interactions: excitons-phonons,
excitons-excitons, exciton-electrons

(3) **recombination**

Depend on oscillator strength
Depend on population of the given state
(But not simply proportional to the density of states)
Selection rules
Radiative vs non radiative

Requires complementary experiments and tuning knobs

Photoluminescence (PL)

Experimental setup

• Advantages

- *No sample preparation/no contacts/no transparent substrate
- *Background free : very weak signals can be detected – to the single-photon ultimate limit
- *Fixed wavelength laser excitation is enough (Choose $E_{\text{laser}} > E_g$)

• Drawbacks

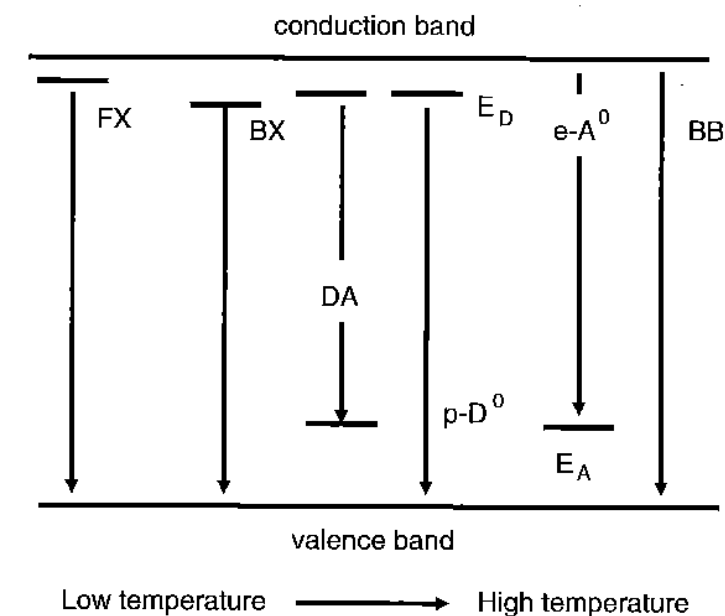
- *Requires low temperature
- *Probes **intrinsic** and **extrinsic** states
(linked to defects, energy levels in the bandgap...)



PL is not simply proportional to density of states
PL energy gives information on lowest states only

- *PL Intensity is affected if transitions are not radiative due to:

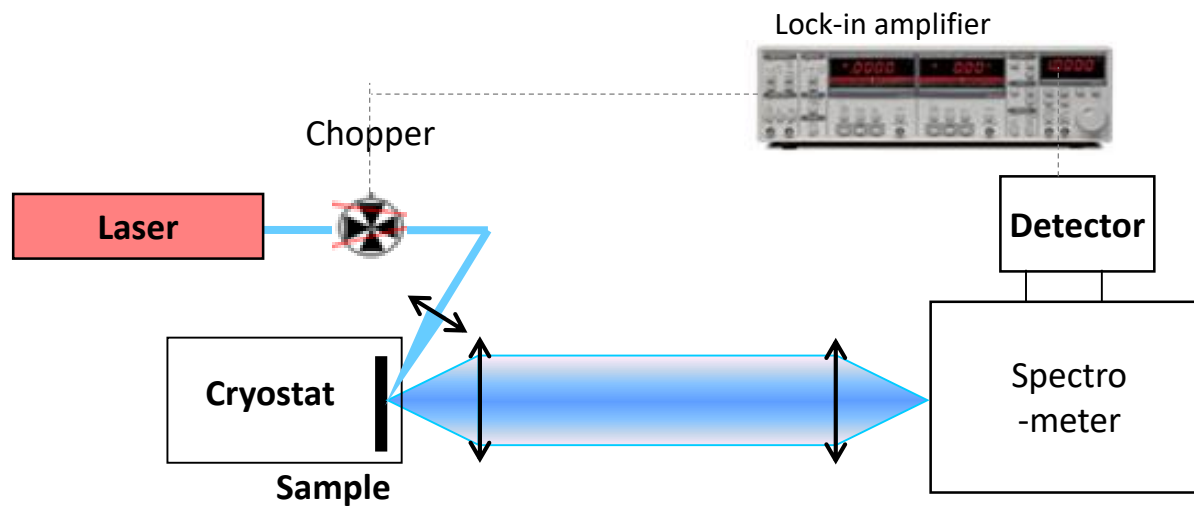
- Recombination on traps or impurities (vacancies, dislocations, joint grains...)
- Surface recombination



Photoluminescence (PL)

Experimental setup

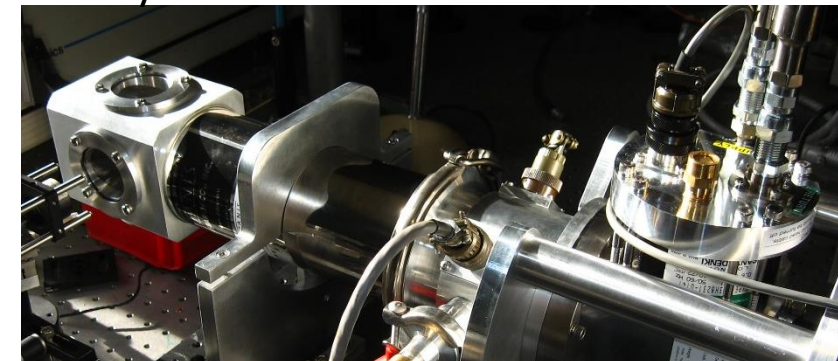
- *Choose $E_{\text{laser}} > E_g$ or Tunable laser
- *Laser spot size = spatial resolution $\sim 100\mu\text{m}$
- *Spectral resolution given by grating and spectrometer



Supercontinuum laser

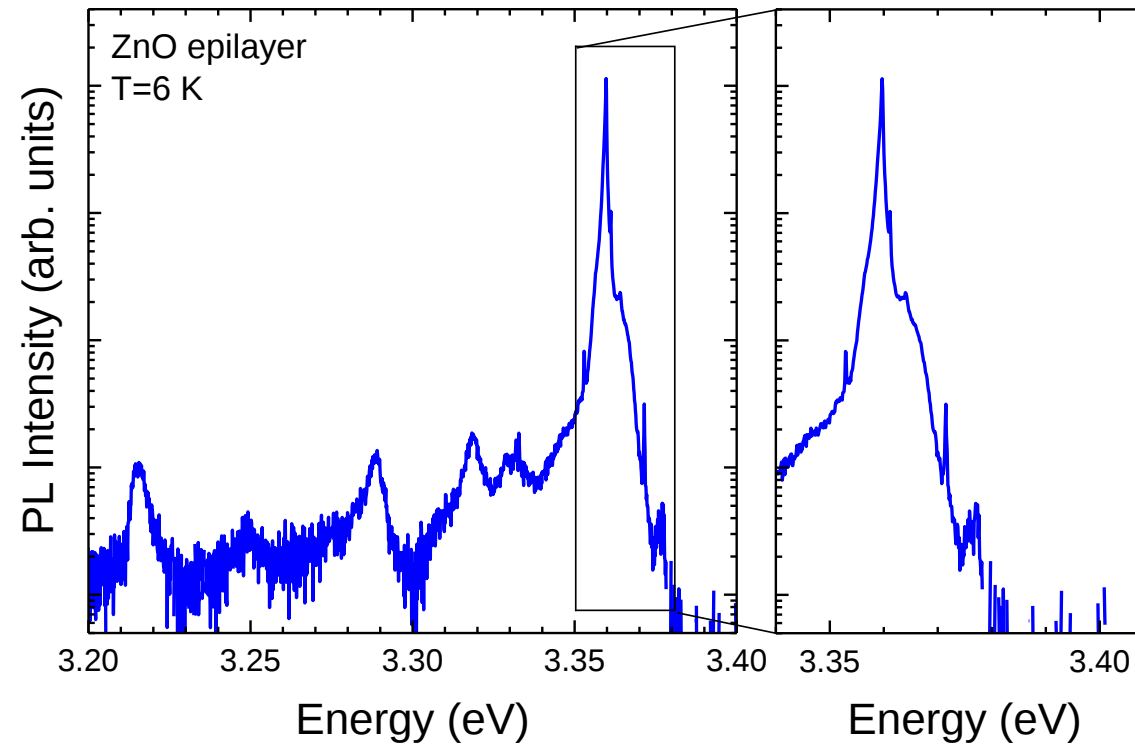


4K cryostat



Photoluminescence

Ex: ZnO epilayer



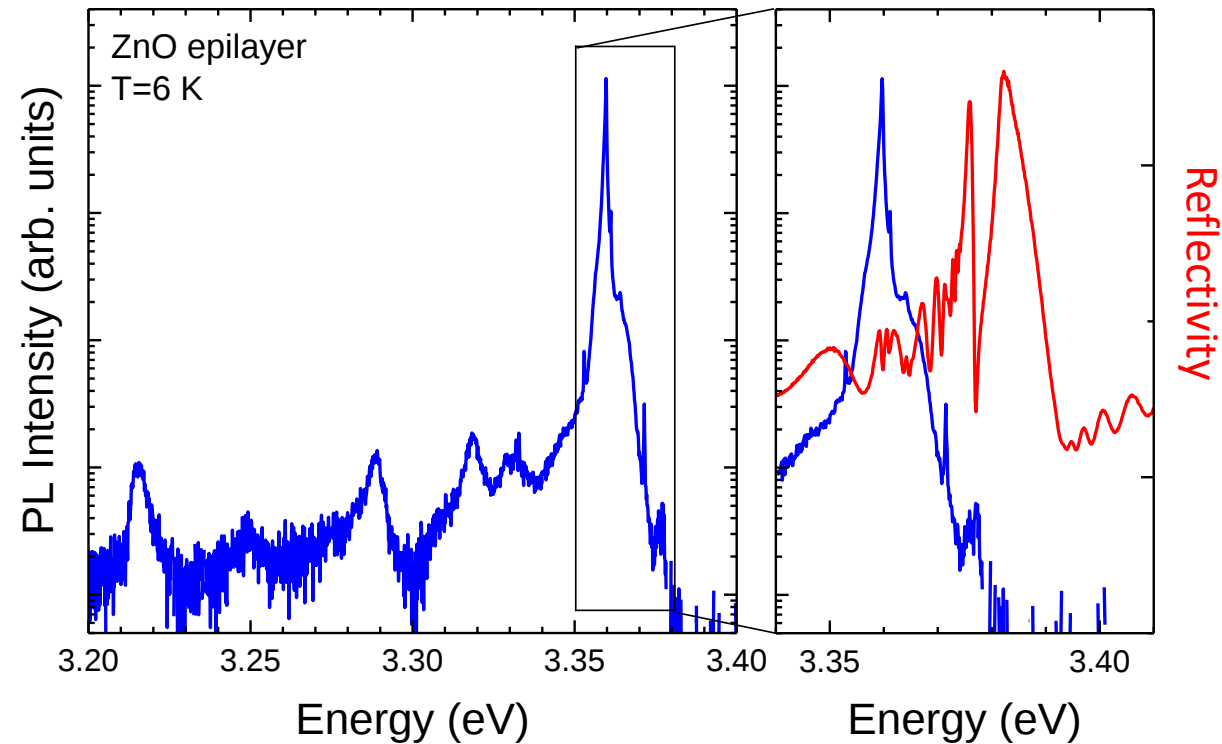
PL spectrum presents many lines !!

➡ Need for complementary experiments

Photoluminescence

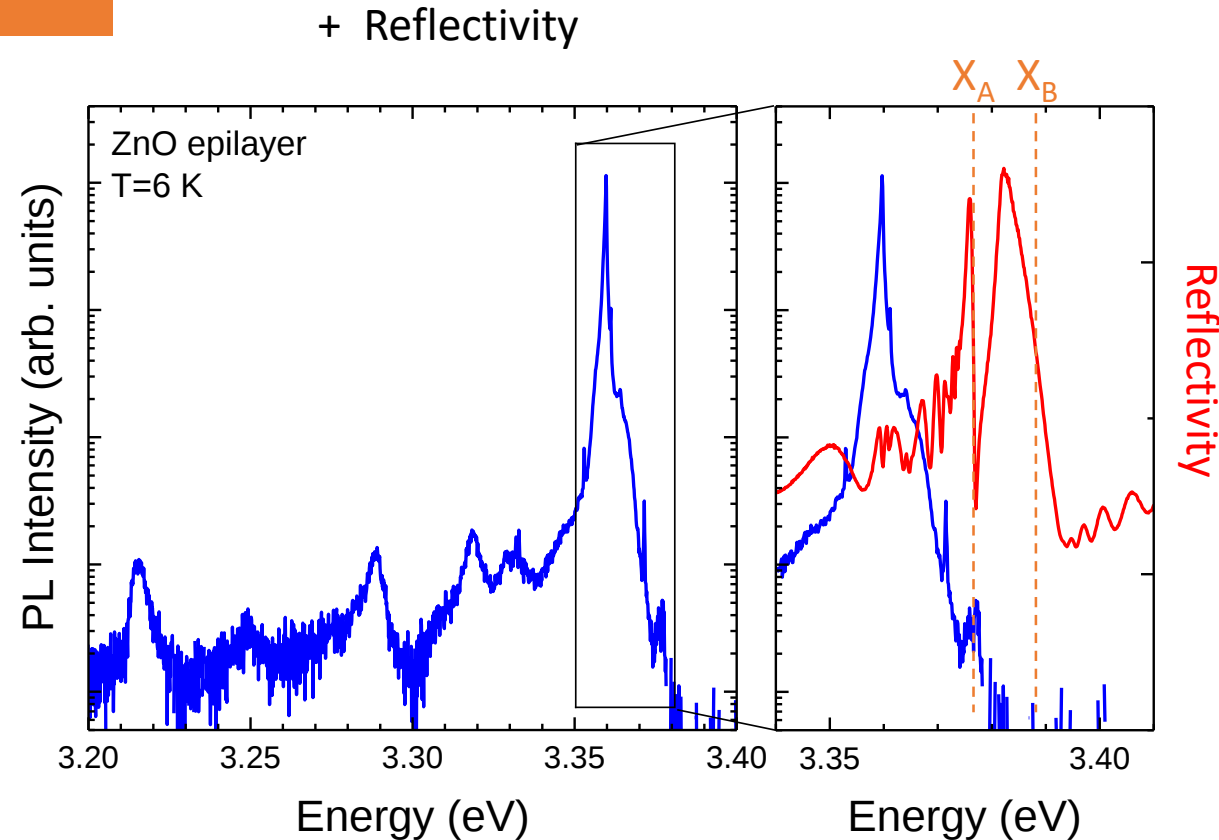
Ex: ZnO epilayer

+ Reflectivity



Photoluminescence

Ex: ZnO epilayer



- * Reflectivity reveals the free exciton lines
- * PL spectrum is dominated by other lines even if oscillator strengths of free exciton transitions are strong

↳ Excitonic complexes bound to donor states

Photoluminescence

Ex: InGaN quantum wells

PL energy gives information on lowest states only

PL widths gives information on inhomogeneity

tuning knob=temperature

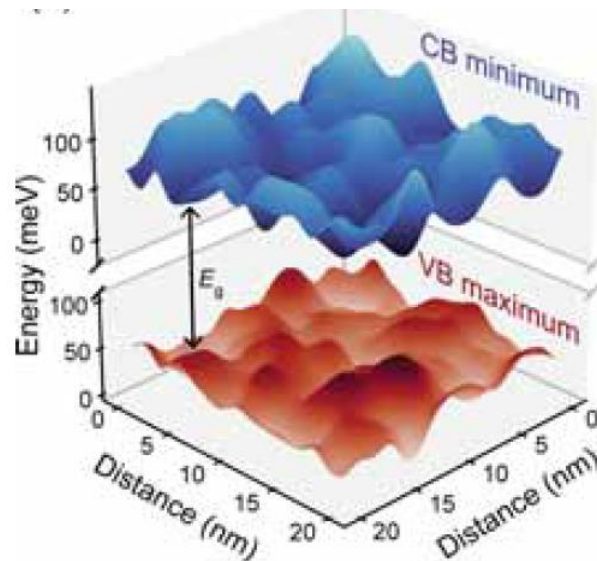
Random alloy disorder in InGaN QW



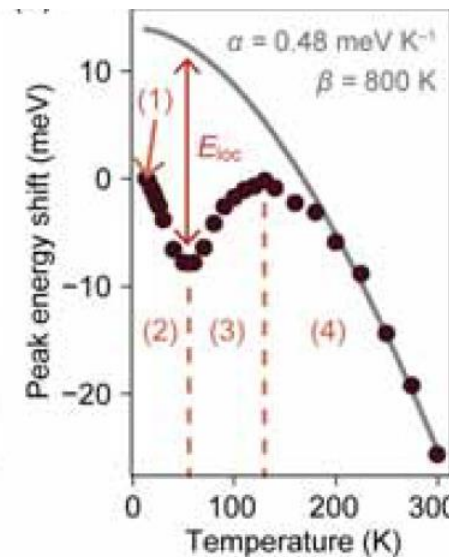
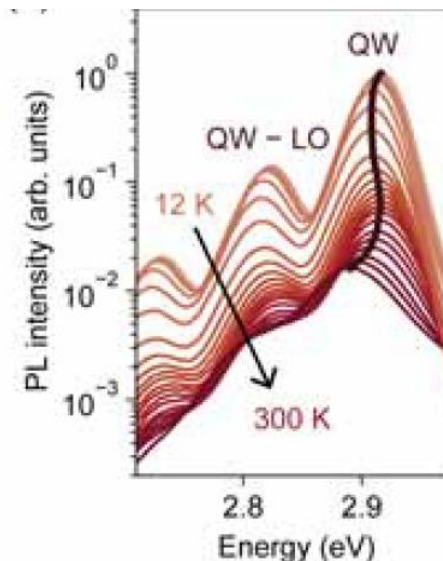
Inhomogeneous broadening of QW PL peak

Localisation: S-shape of $E=f(T)$

Non radiative channels: low signal at 300 K



EFFL Thomas Weitherley PhD thesis



Varshni's law

$$E_g(T) = E_g(0) - \frac{\alpha T}{T + \beta}$$

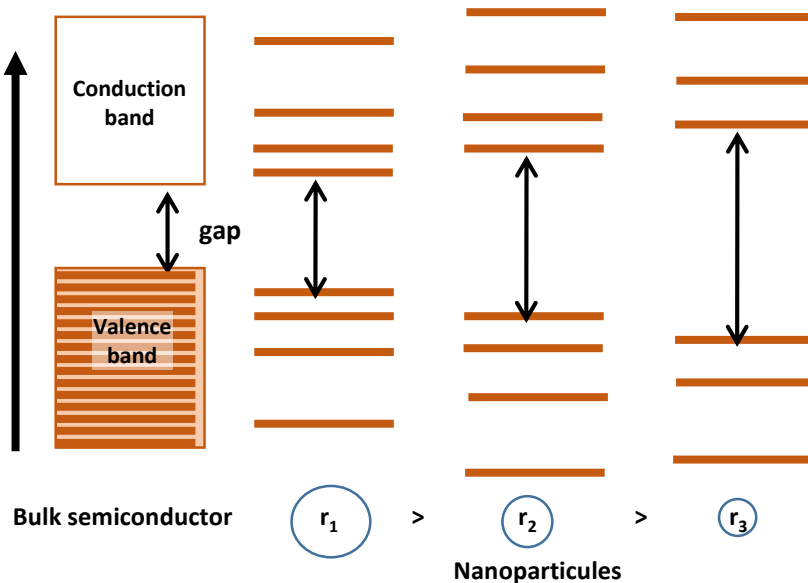
Photoluminescence

Ex: Quantum dots / Nanoparticles

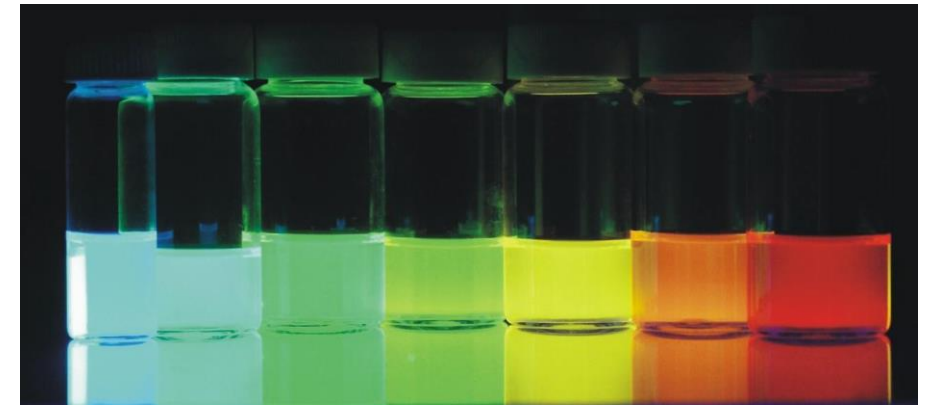
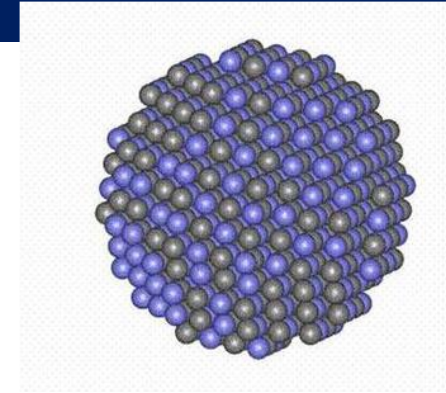


Nobel prize 2023

Confinement in nanoparticles: effect on the optical properties



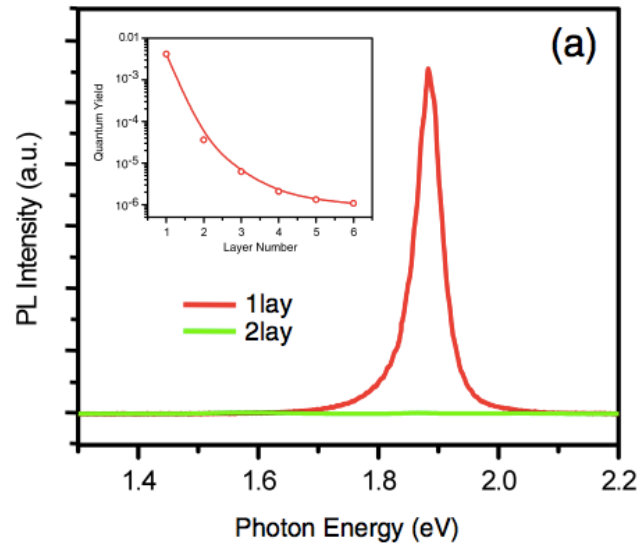
Reduction of size of NPs → increase of the energy bandgap



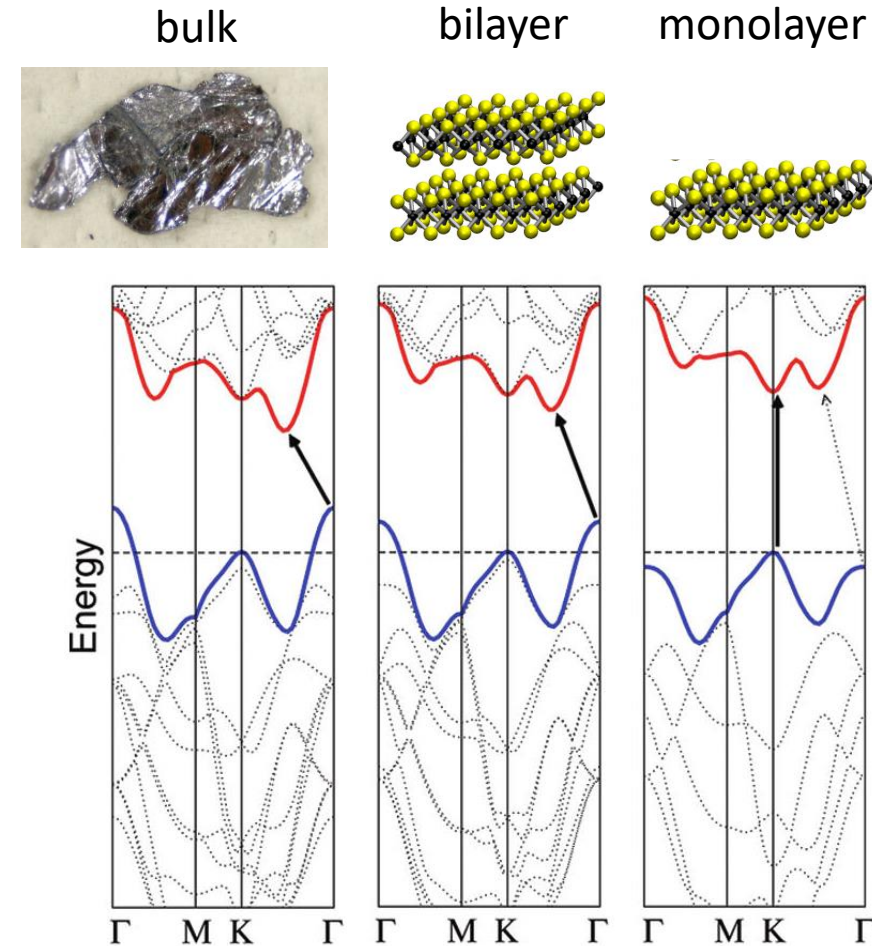
Size (nanometers)

Photoluminescence

Ex: MoS₂



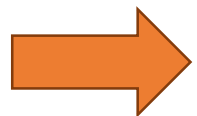
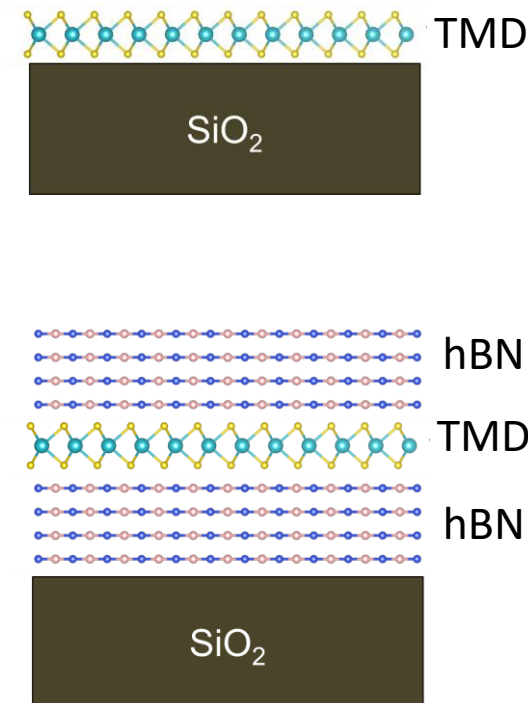
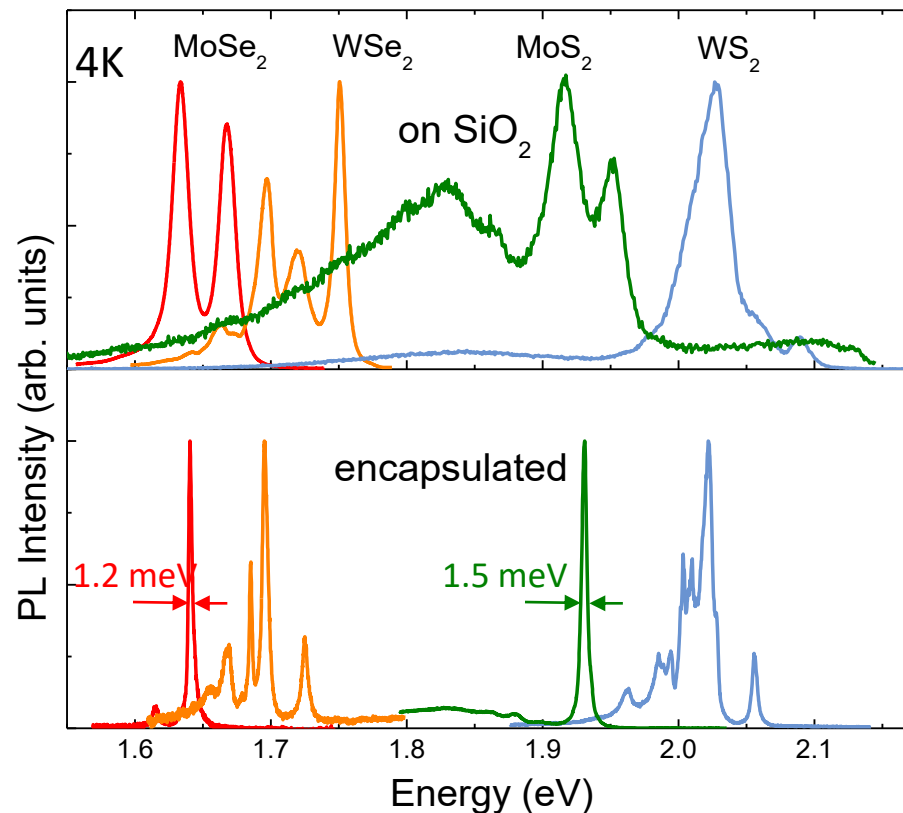
From multilayer to monolayer  Indirect to direct bandgap



- A. Splendiani et al. Nano Lett. 10, 1271 (2010)
- B. Mak et al. Phys. Rev. Lett. 105, 136805 (2010)

Photoluminescence

Ex: TMDC



The optical quality is strongly influenced by the environment

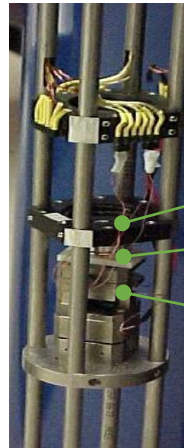
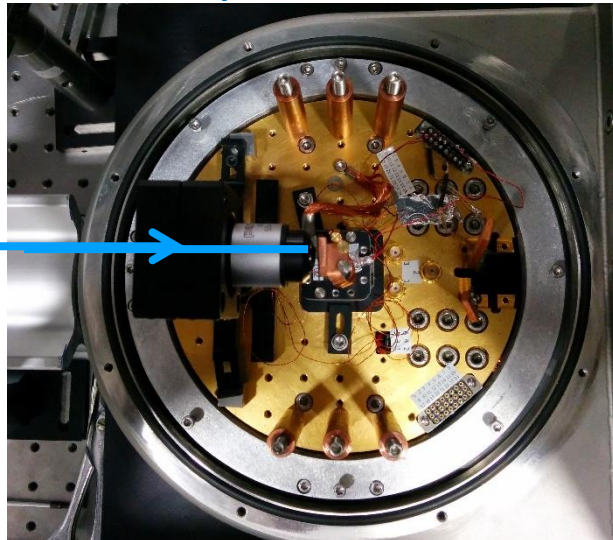
hBN encapsulated samples → narrow transitions → limited by homogeneous linewidth

μ -Photoluminescence

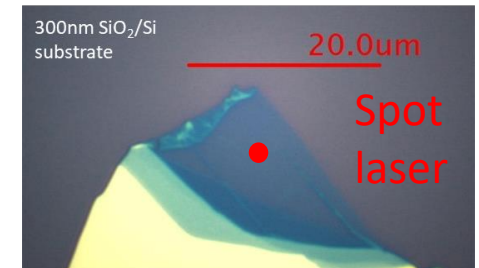
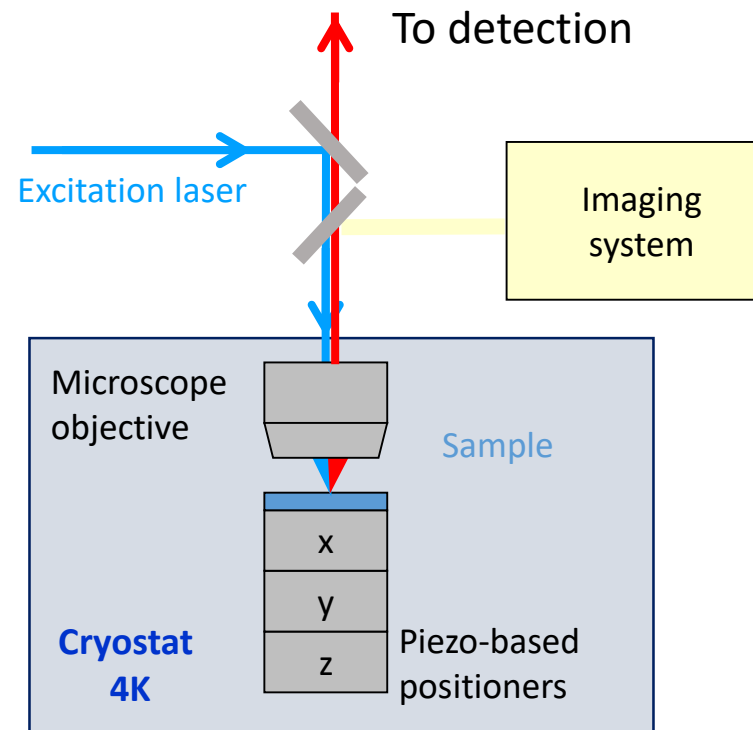
Experimental setup

- Size of monolayer $\sim 10\ \mu\text{m}$
 - Low temperature required
 - Wavelength-tunable LASER
 - Sensitive detector
- Need for laser spot $\sim 1\ \mu\text{m}$
 → Precise sample positioning
 → Optical cryostat with no vibration & adjustable temperature 4 K – 300 K

4K Cryostat

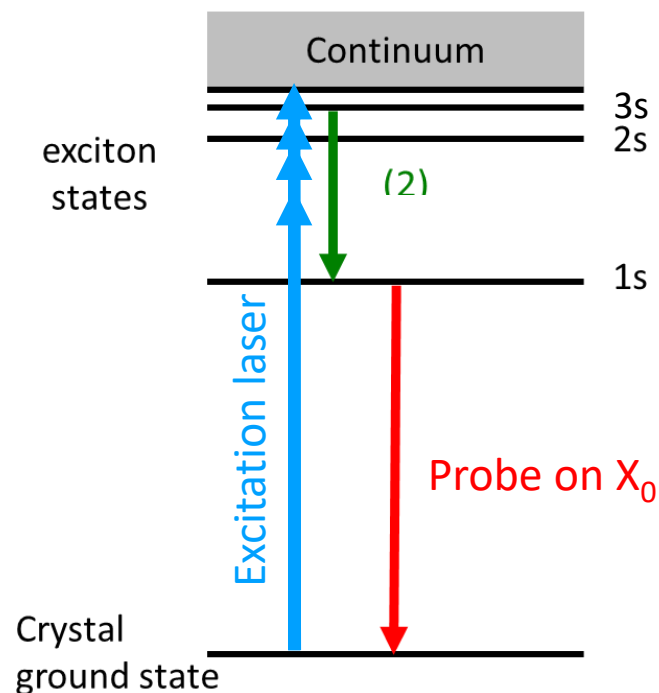


Objective
Sample
Nano-positioners



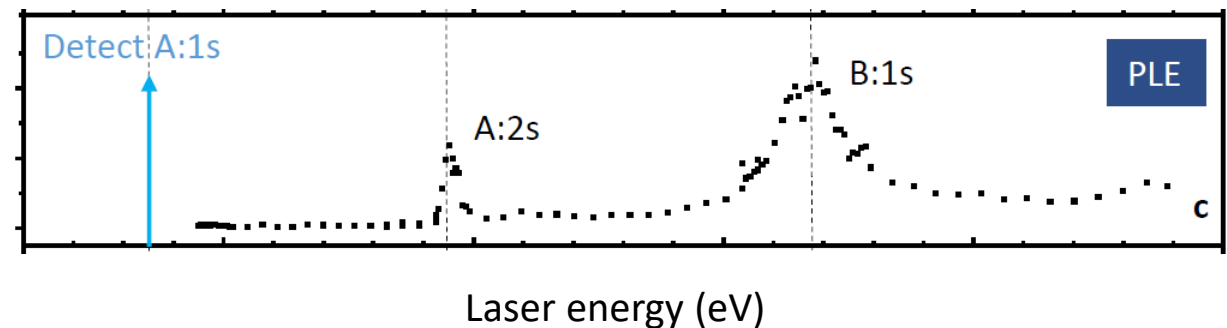
Photoluminescence excitation (PLE)

Ex: TMDC



tuning knob=laser energy

X^0 PL intensity

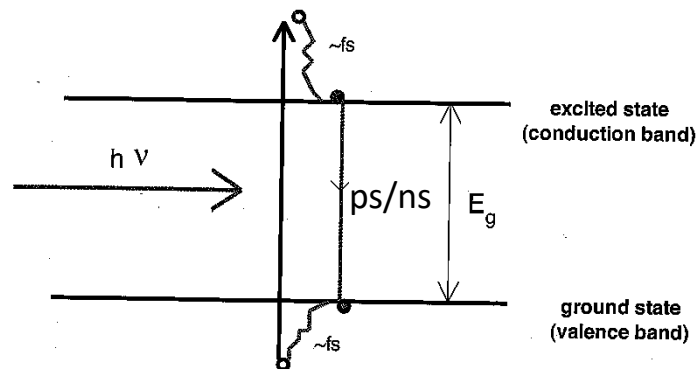


Phys. Rev. Lett. 114, 97403 (2015)

- *Probes the density of states but !! Sensitive to relaxation mechanisms
→ different from absorption measurements
- *Need for wavelength tunable laser source

Time-resolved spectroscopy

tuning knob=time



➤ Coherent Regime (< 1 ps)

- momentum scattering
- carrier-carrier scattering
- carrier-optical phonons

➤ Non-Thermal Regime (few ps)

- e-h scattering
- intersubband scattering
- capture

➤ Hot-Thermal Regime (< 100 ps)

- phonon scattering (acoustic, optical)

➤ Equilibrium with lattice (ns)

- recombination (radiative and non-radiative)

Four-Wave Mixing

- Pump-probe
- TR-Photolum.

Time-resolved PL

- *Information on the dynamics of the carrier population
- *Information on relaxation times
- *Information on radiative or non-radiative recombination channels

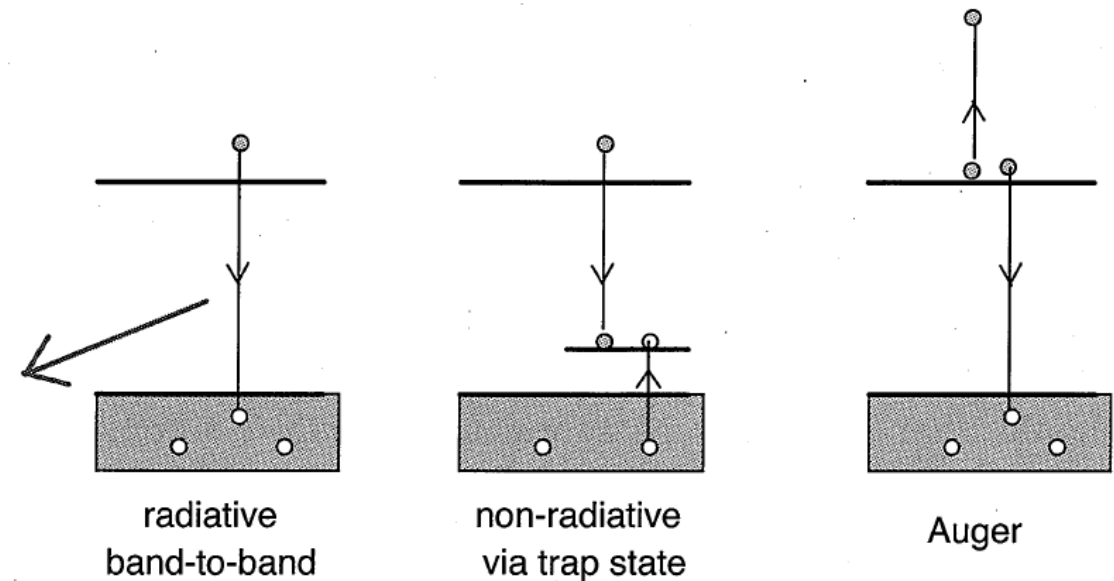
$$\frac{dn}{dt} = -An - Bn^2 - Cn^3$$

Hypothesis: $n=p$

A: recombination coefficient on defects $A \sim \frac{1}{\tau_{non-rad}}$

B: radiative recombination coefficient $B \sim \frac{1}{\tau_{rad}}$

C: non-radiative recombination coefficient, type Auger



Time-resolved PL

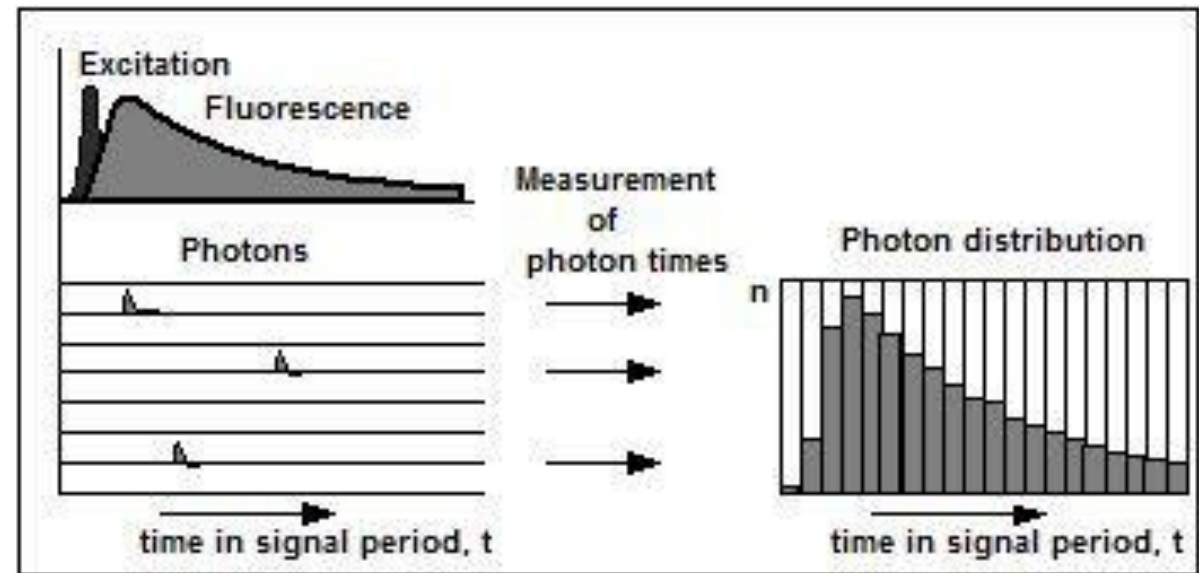
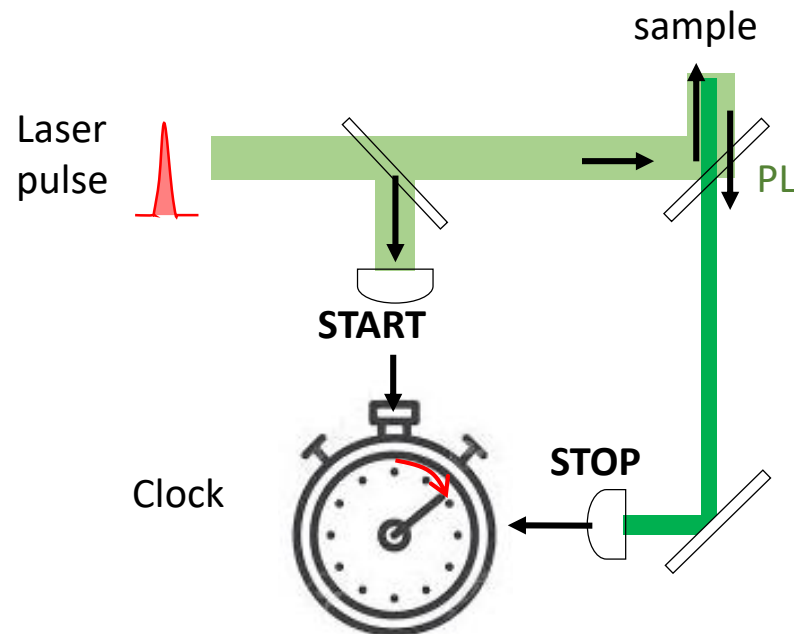
increased
time
resolution



- ☐ With fast-detection electronics (ns and slower)
- ☐ With intensified CCD (ICCD) cameras (≥ 500 ps)
- ☐ With Avalanche Photodiode (SPAD, APD)
or Time Correlated Single Photon Counting TCSPC (ps and slower)
- ☐ With a Streak camera ($< \text{ps}$ and slower)
- ☐ With optical gating (fs and slower)

Time-resolved PL

With time-correlated single photon counting system (TCSPC)



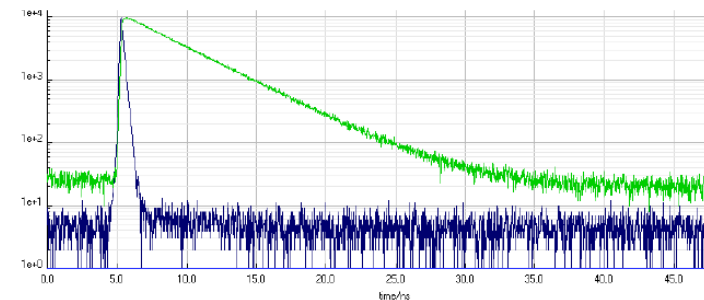
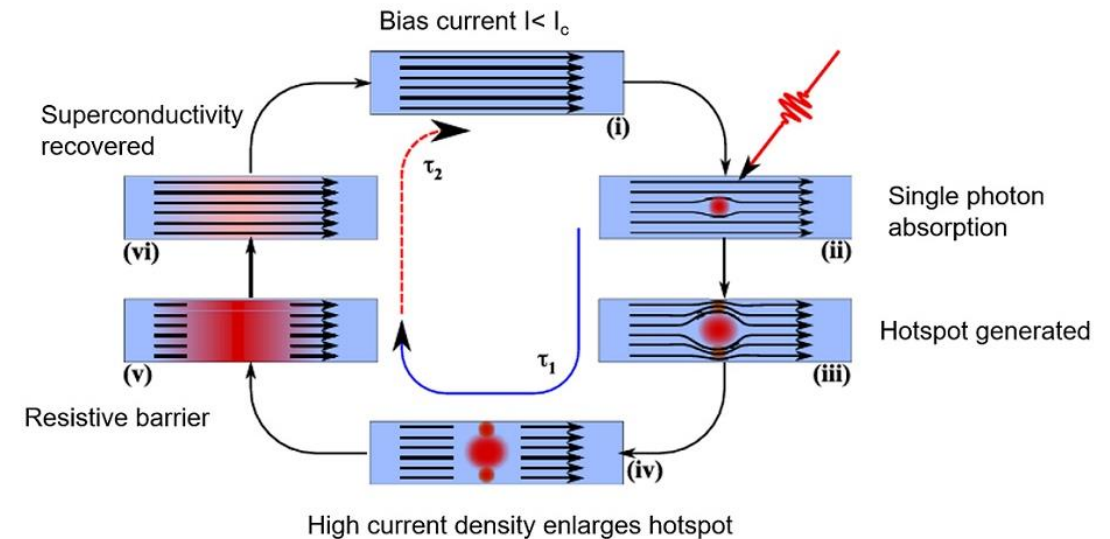
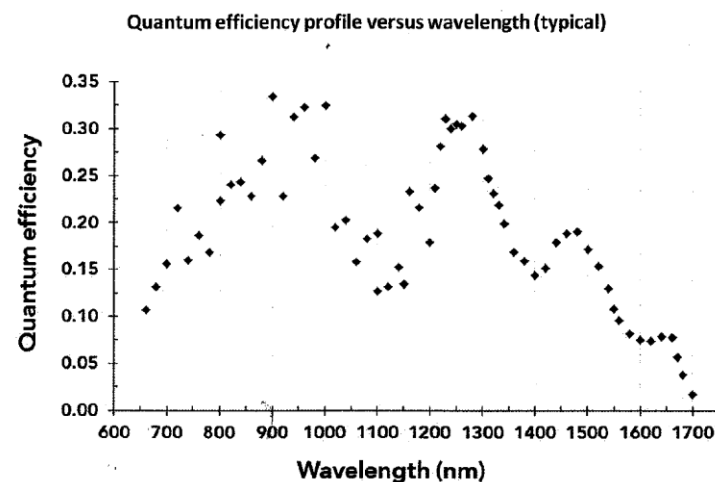
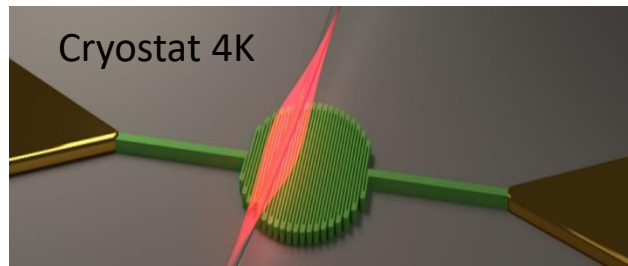
Photons can be detected by photodetectors with intrinsically high gain such as photomultipliers / micro-channel plate / single photon avalanche photodiodes (SPAD)

TCSPC can be used to measure lifetimes ranging from <50 ps to over 50 ms

- + Cheap
- single-channel
- single photon : to be repeated a large number of times to get good S/N signal

Time-resolved PL

With superconducting nanowire single photon detector (SNSPD)



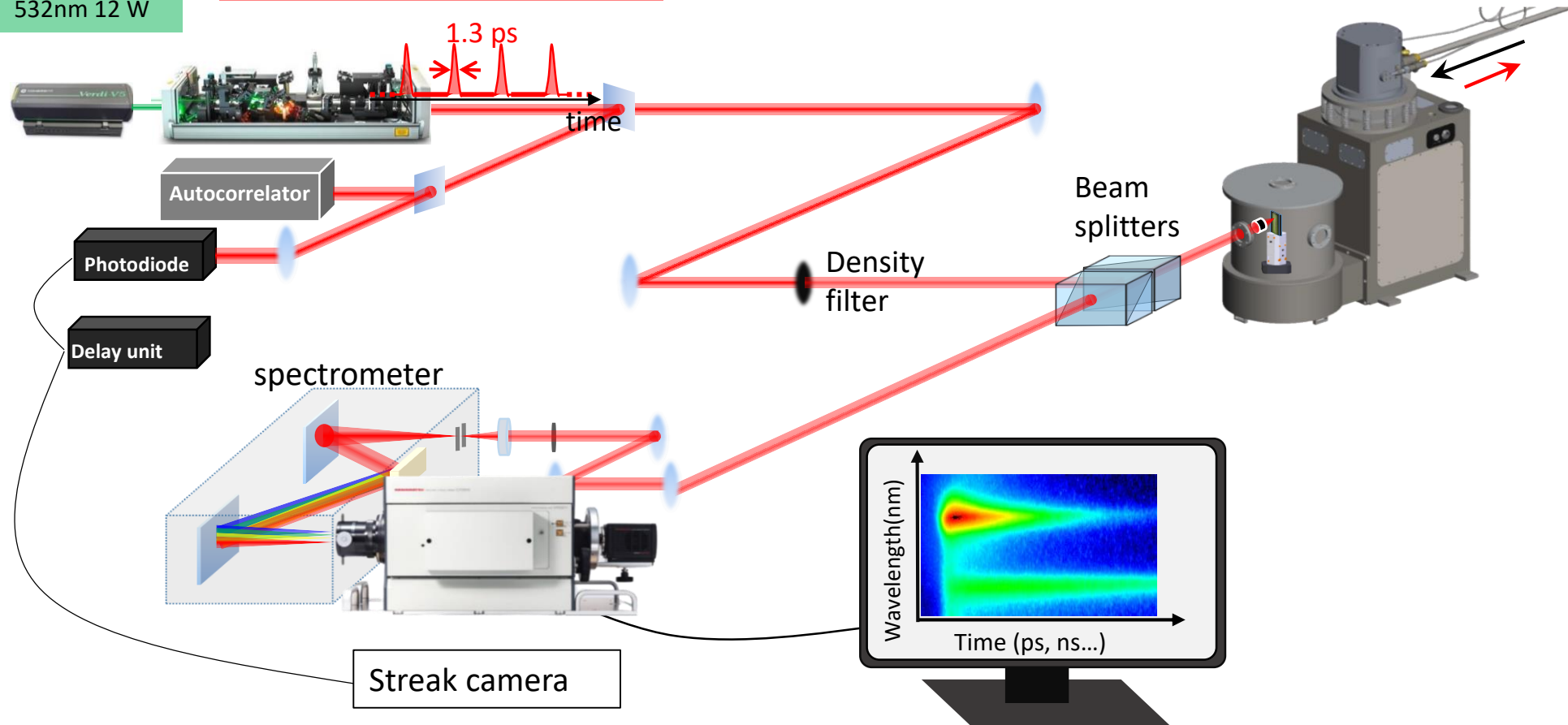
- ☐ Can detect single photons with high efficiency over a broad spectral range (even NIR)
- ☐ High timing resolution of less than 50 ps

Time-resolved PL

With streak camera

Continuous
Pump laser
532nm 12 W

ps-Pulsed laser
Repetition rate 80 MHz



Closed-cycle Cryostat: 4K to RT
Achromatic objective (NA = 0.82, $\theta=55^\circ$)
Piezo-based nano-positioners (x, y, z)

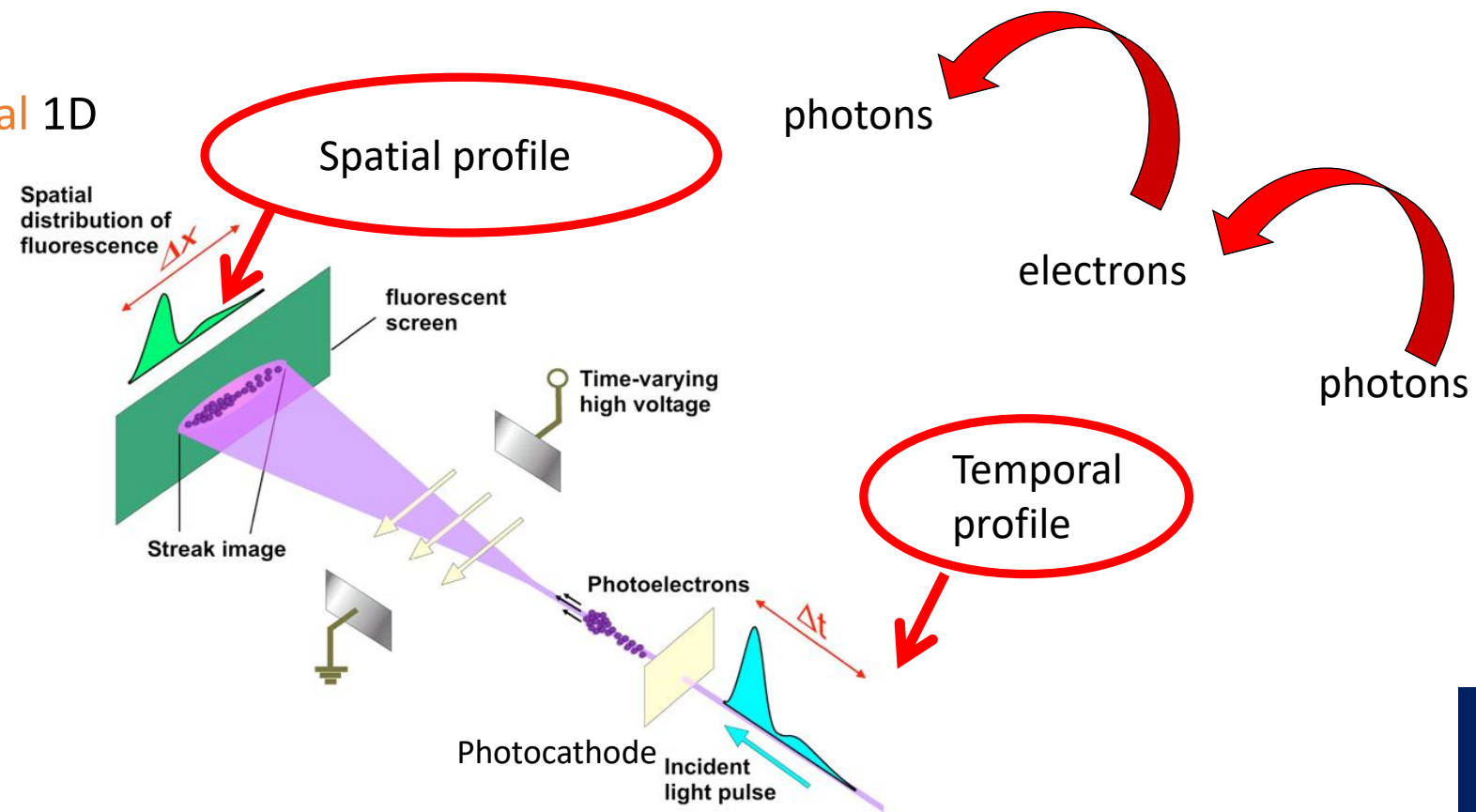
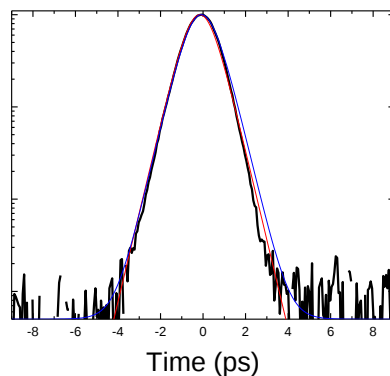
PL intensity as a function of
wavelength and time

Time-resolved PL

With streak camera

- ✓ **Time** resolution (few ps)
 - Rep rate up to 80MHz
- ✓ **Spectral** resolution or **spatial** 1D
- ✓ **Luminescence** flux
- ✓ **Time resolution** <1 ps

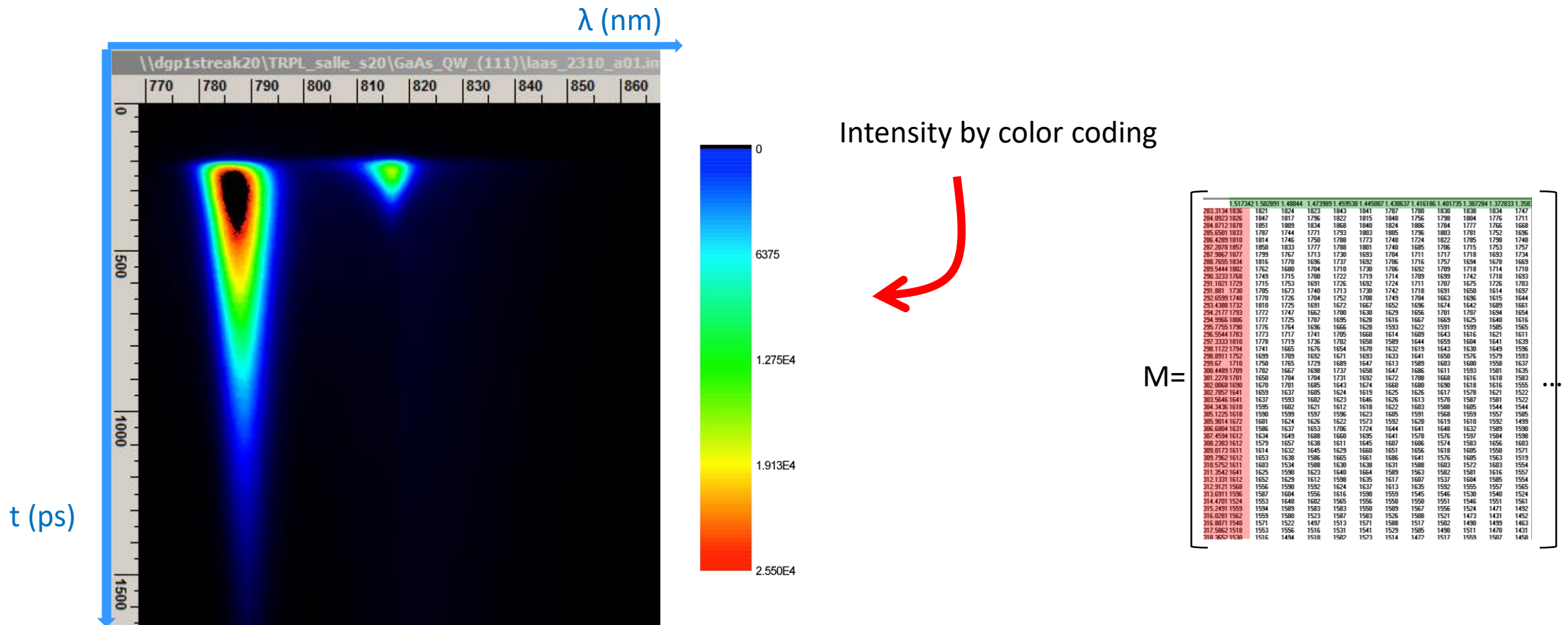
Instrument response function (IRF)



Time-resolved PL

With streak camera

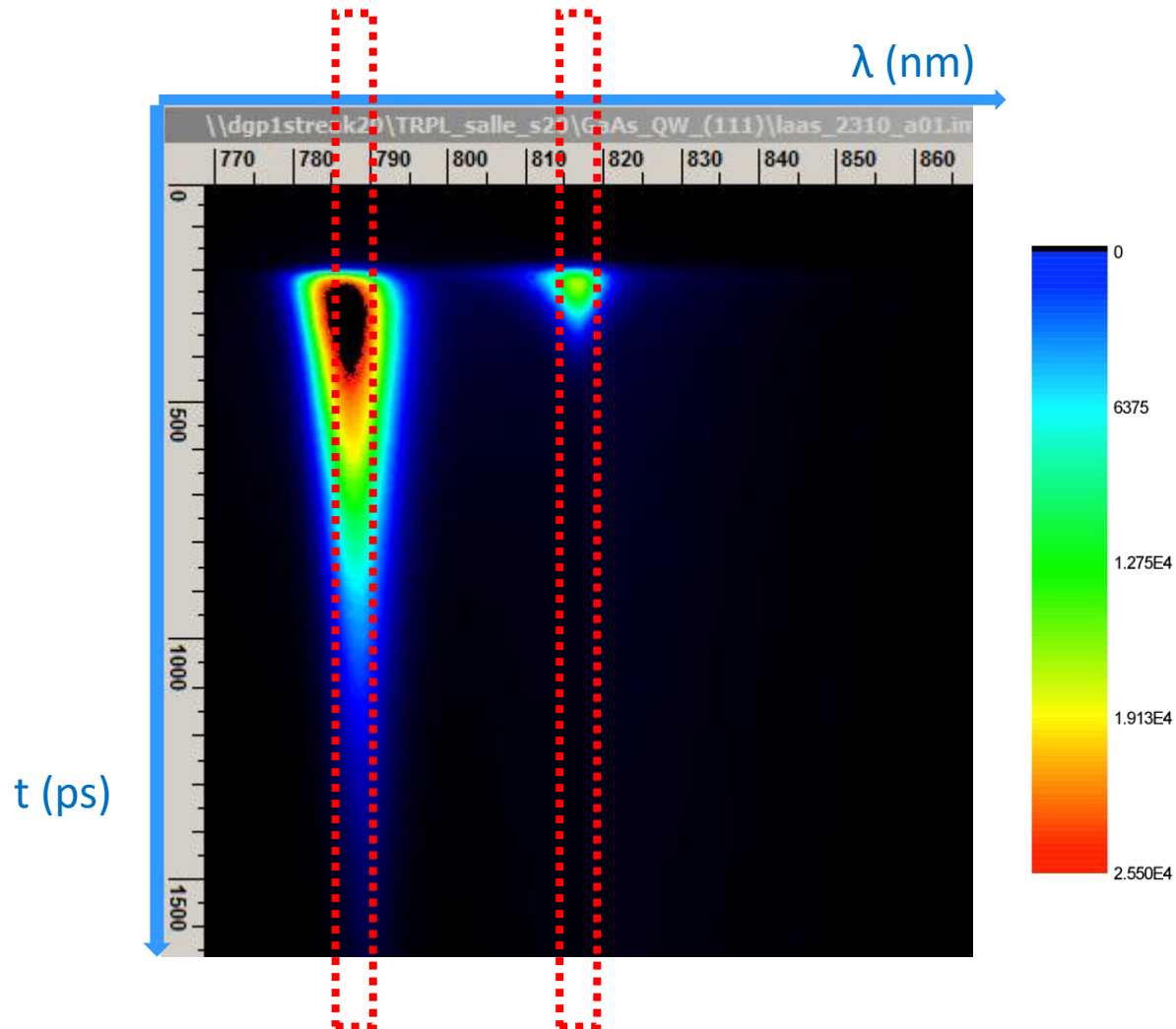
« tridimensionnal » images = time + intensity + spectrum



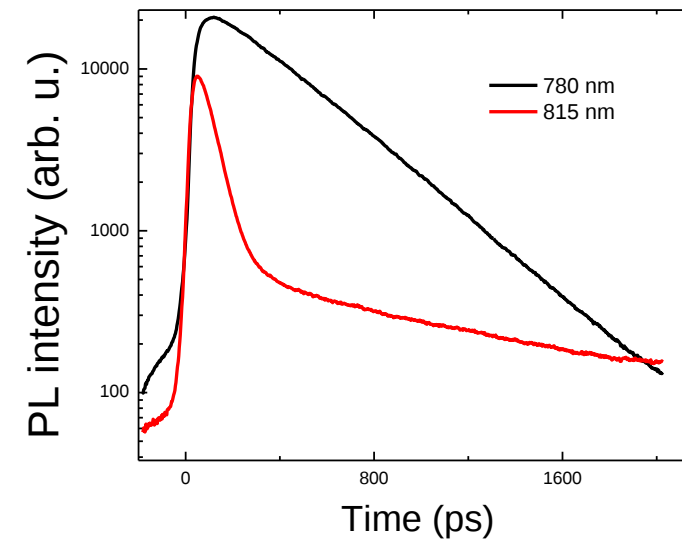
Time-resolved PL

With streak camera

« tridimensionnal » images = time + intensity + spectrum



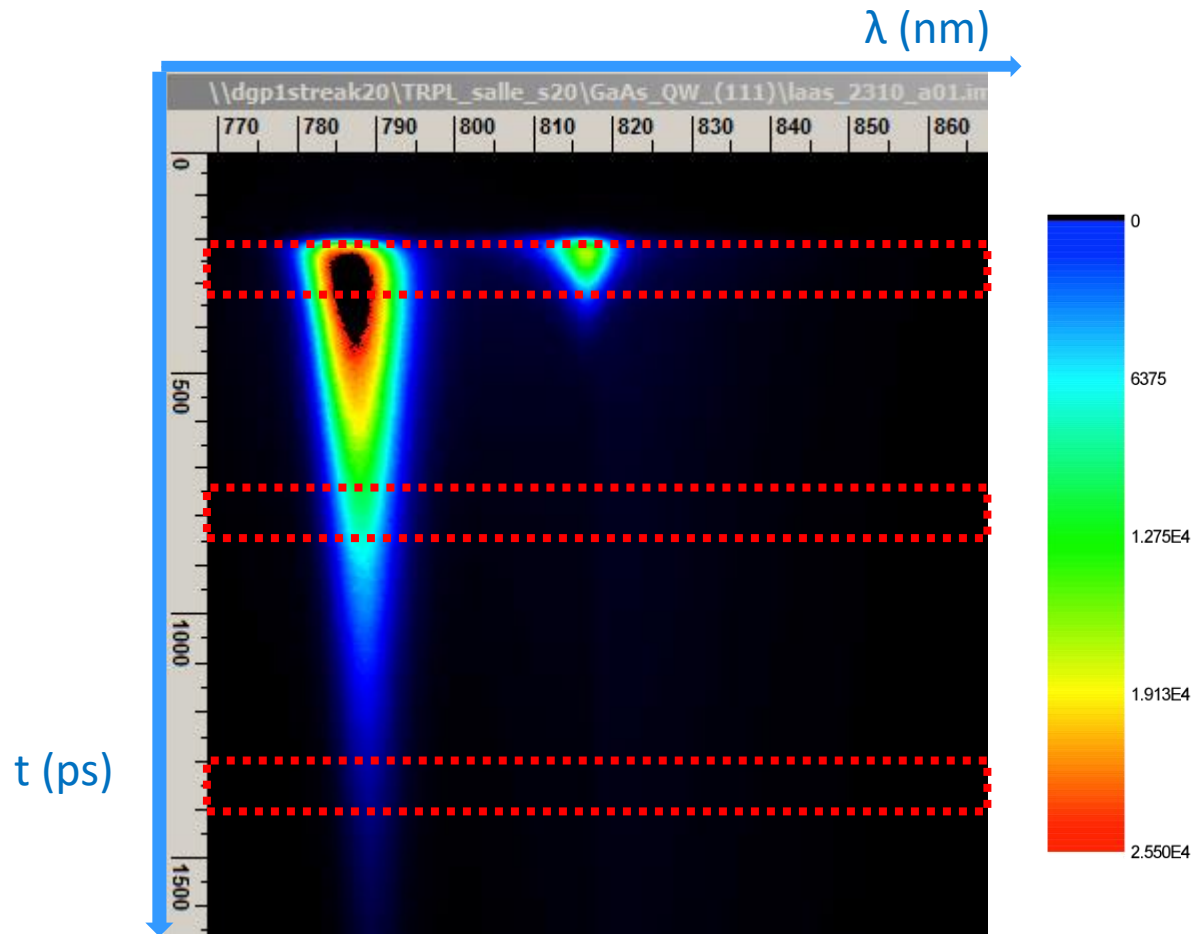
Spectral selection



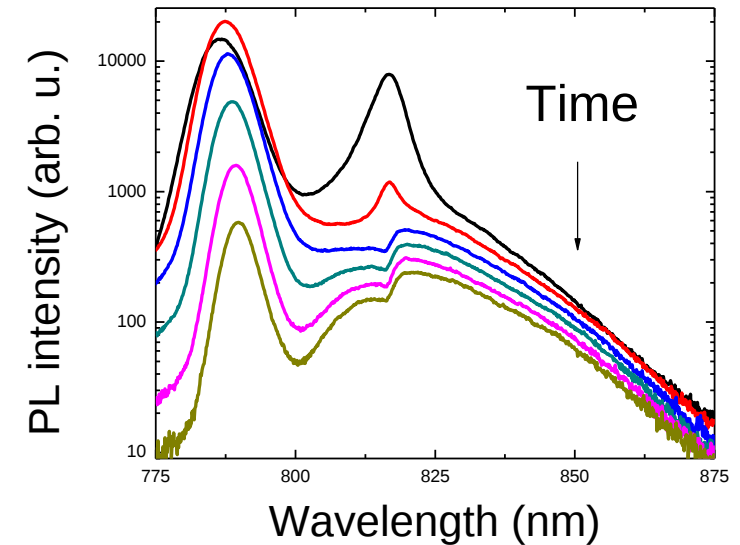
Time-resolved PL

With streak camera

« tridimensionnal » images = time + intensity + spectrum



« Slices » of time delay

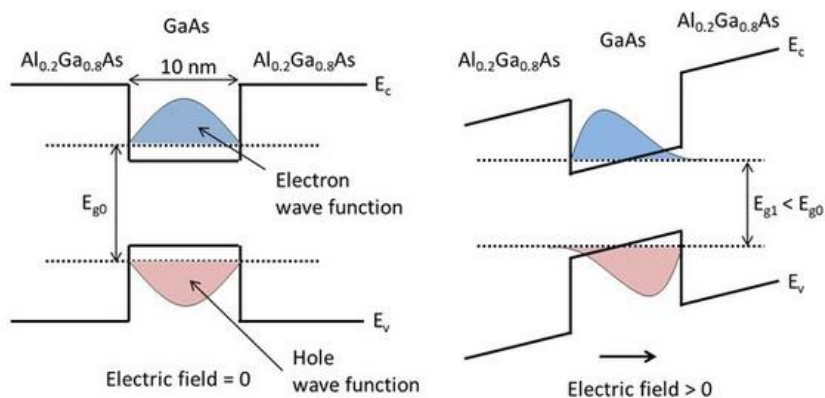


- Carrier cooling
- Diffusion
- ...

Time-resolved PL

Ex: Quantum wells

Internal or external electric field on QW:



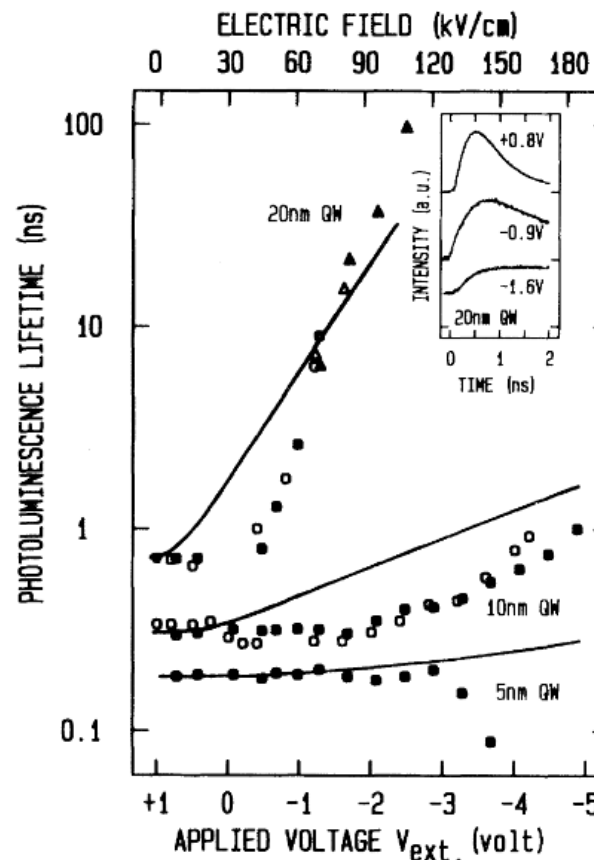
$$B_{vc} = \frac{pq^2}{2\hbar} \cdot \left| \langle Y_c(r) | E \cdot \hat{r} | Y_v(r) \rangle \right|^2 \times d(E_c - E_v - \hbar\omega)$$

Einstein
coefficient

Coupling of Ψ_v to Ψ_c

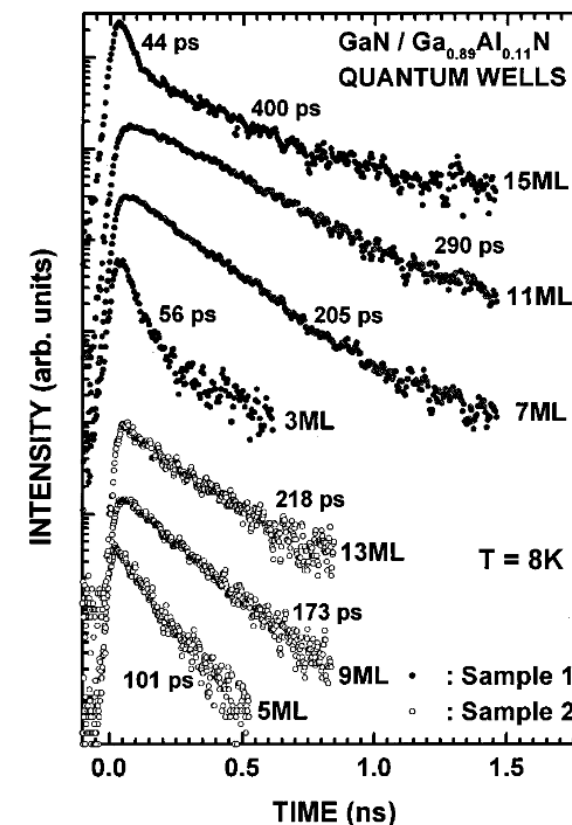
- *Plays with the electron and hole wavefunction overlap
- *Increase of QW decay times with E

GaAs QW



Phys. Rev. Lett. 55, 2610 (1985)

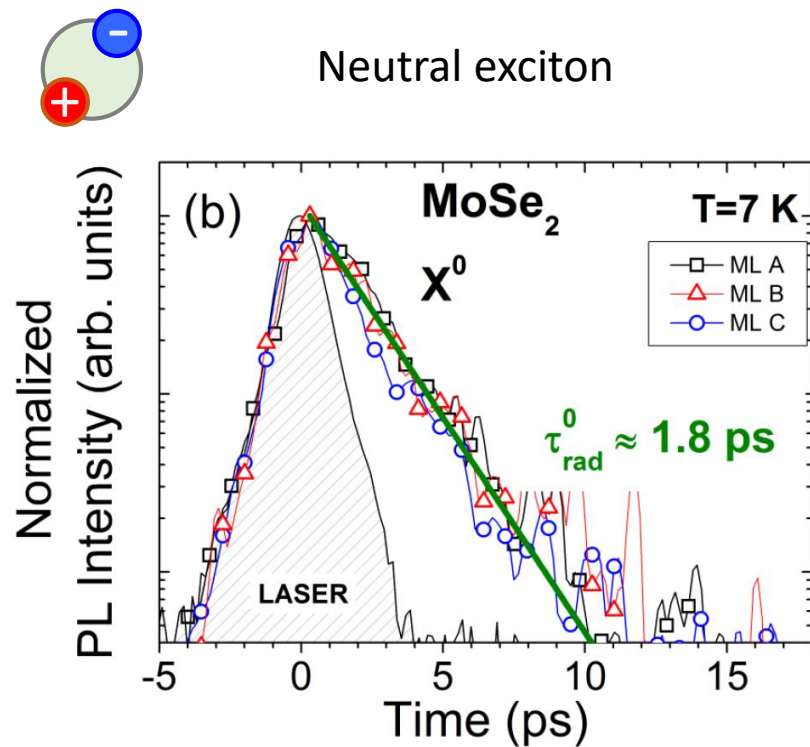
GaN QW



Phys. Rev. B 59, 15363 (1999)

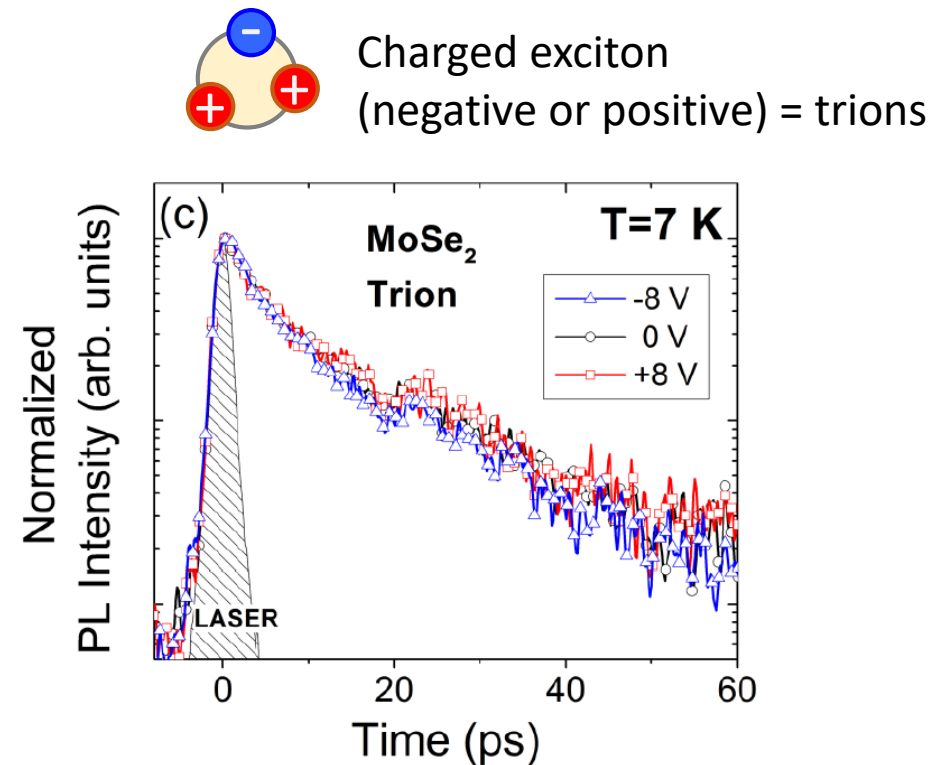
Time-resolved PL

Ex: TMDC



Exciton radiative lifetime $\sim 1.8\text{ ps}$

Exciton Bohr radius $\sim 0.5\text{ nm}$

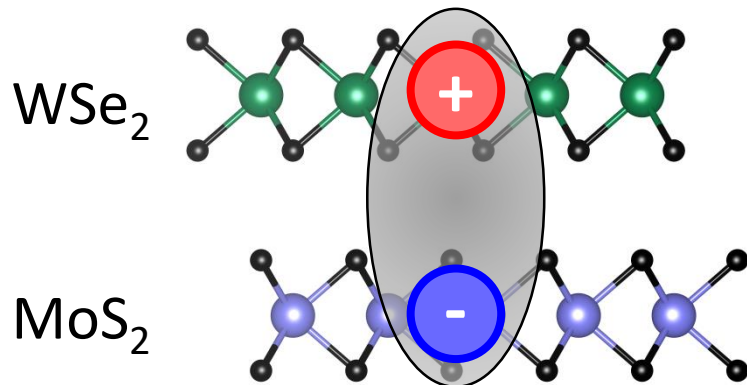


Trion decay $\sim 15\text{ ps}$

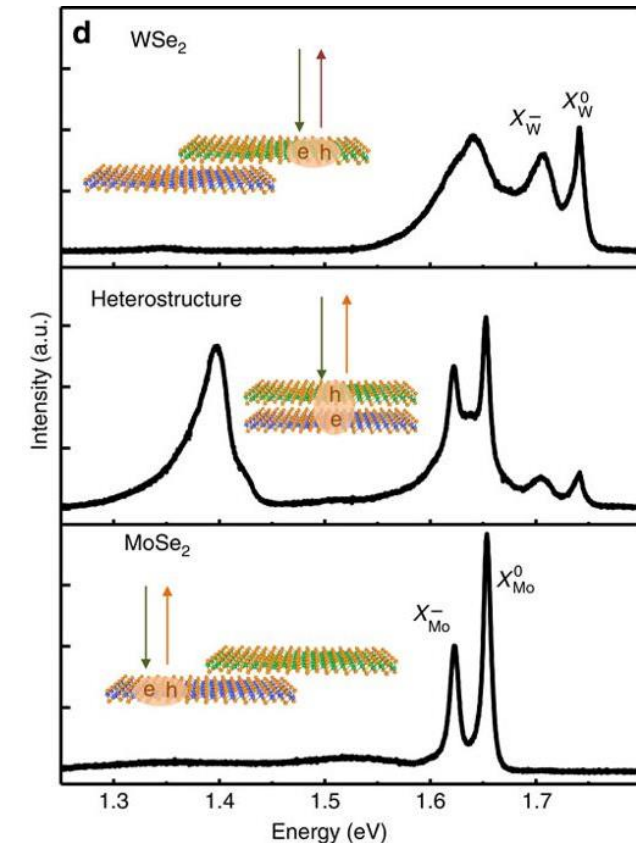
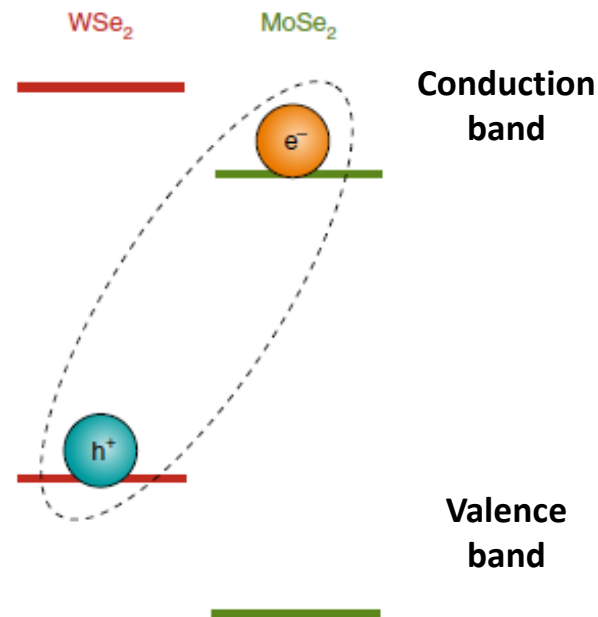
Time-resolved PL

Ex: TMDC heterostructure

Spatially indirect excitons



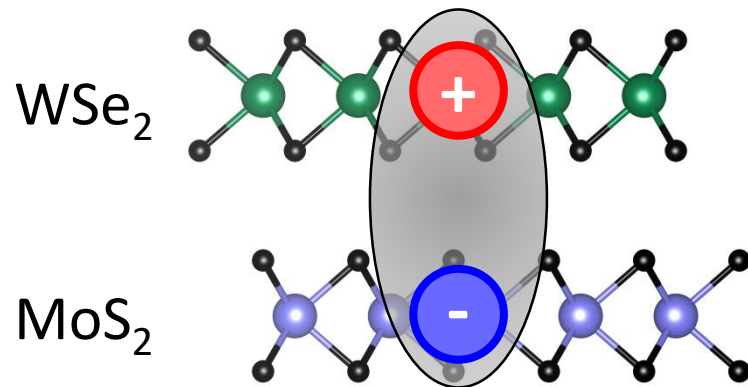
Type II band alignment



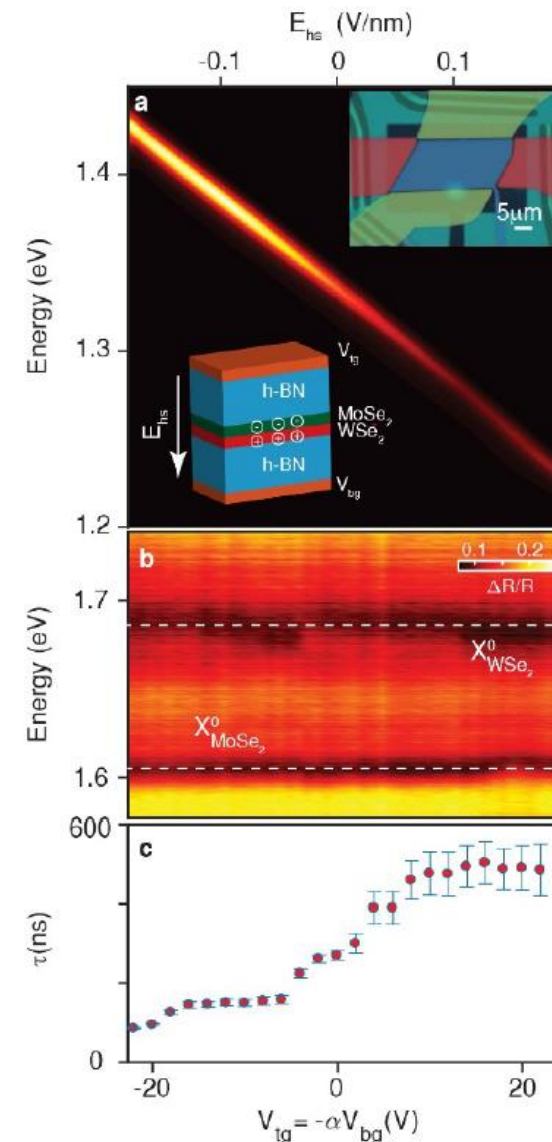
Nat. Comm. **6**, 6242 (2015)

Time-resolved PL

Ex: TMDC heterostructure



Tuning of indirect exciton lifetime with electric field
From ~ 250 ns to 600 ns for $E_{hs} > 0.1$ V/nm

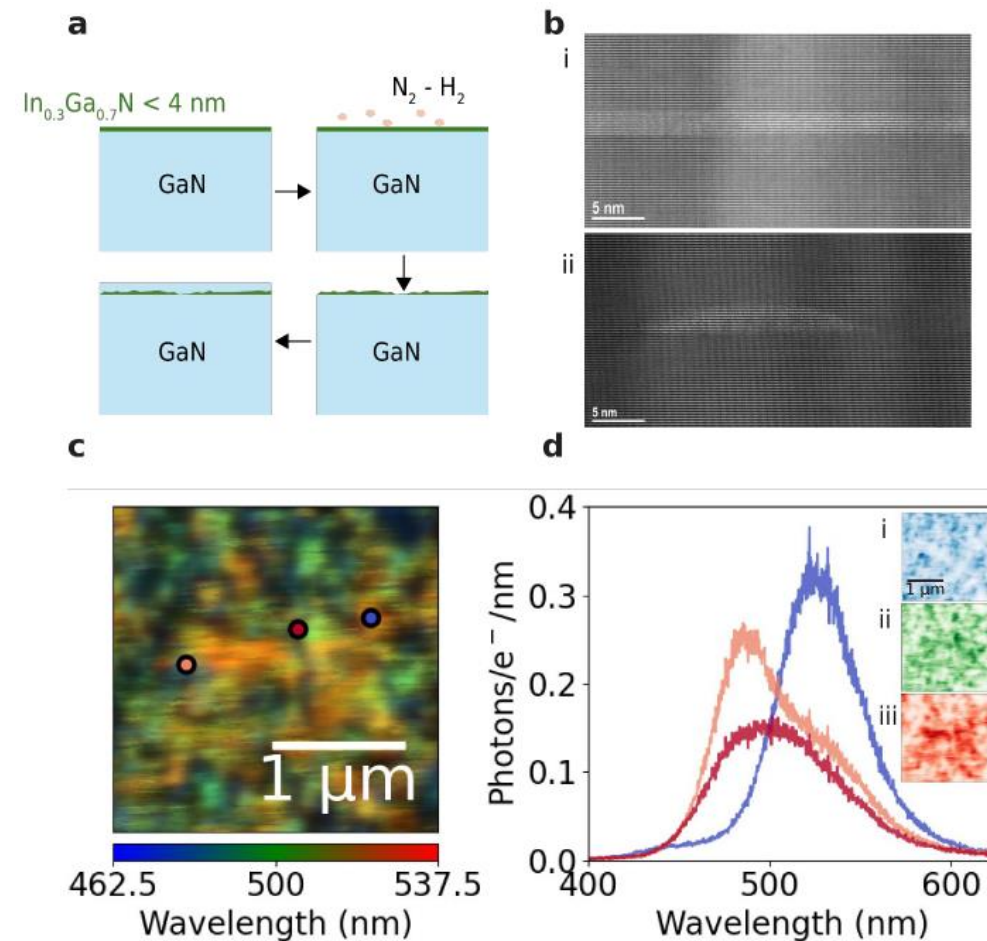
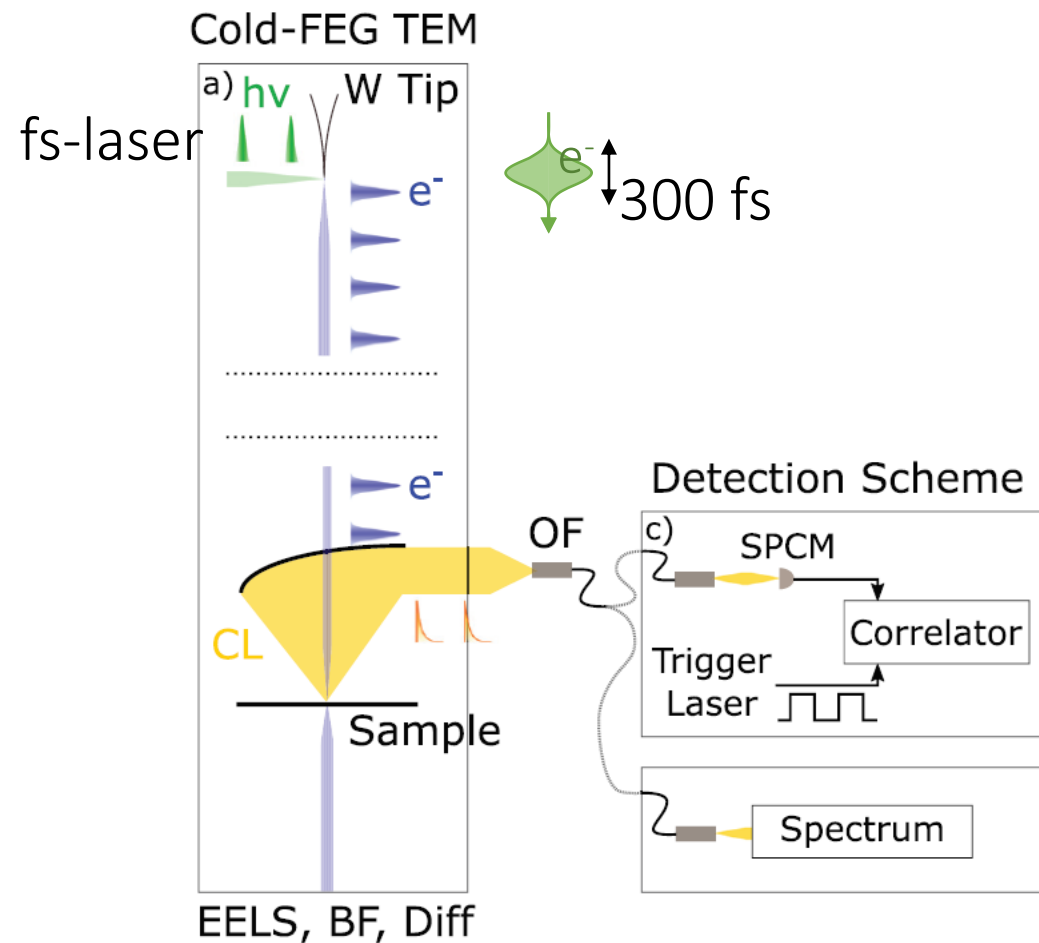


Stark effect

Science 366, 870 (2019)

Time-resolved Cathodoluminescence

Ex: InGaN quantum wells/quantum dots



➔ Diffusion of Indium creates QW/QD areas

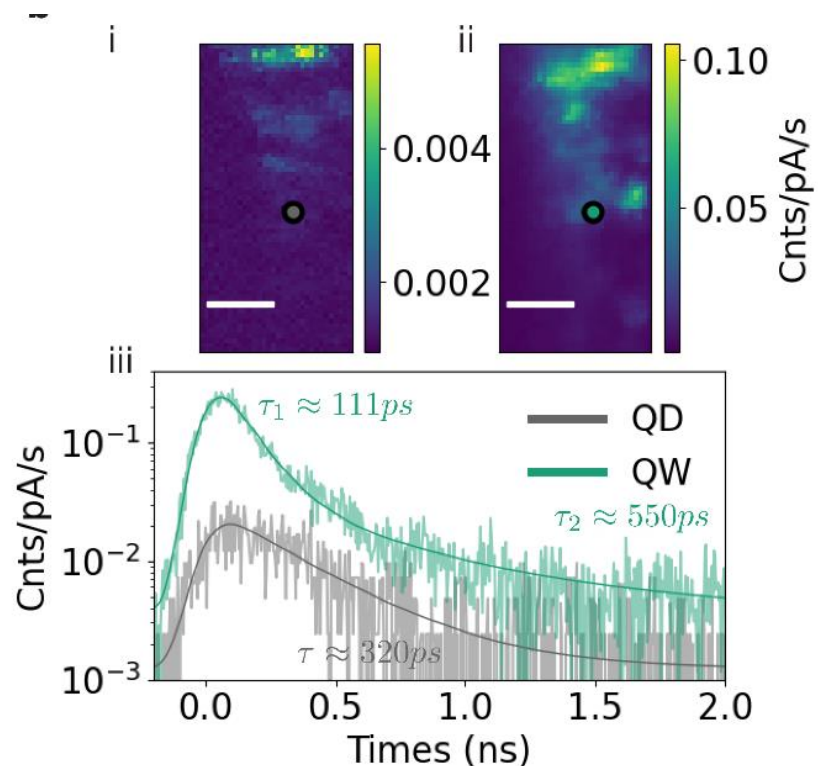
Time-resolved Cathodoluminescence

Ex: InGaN quantum wells/quantum dots

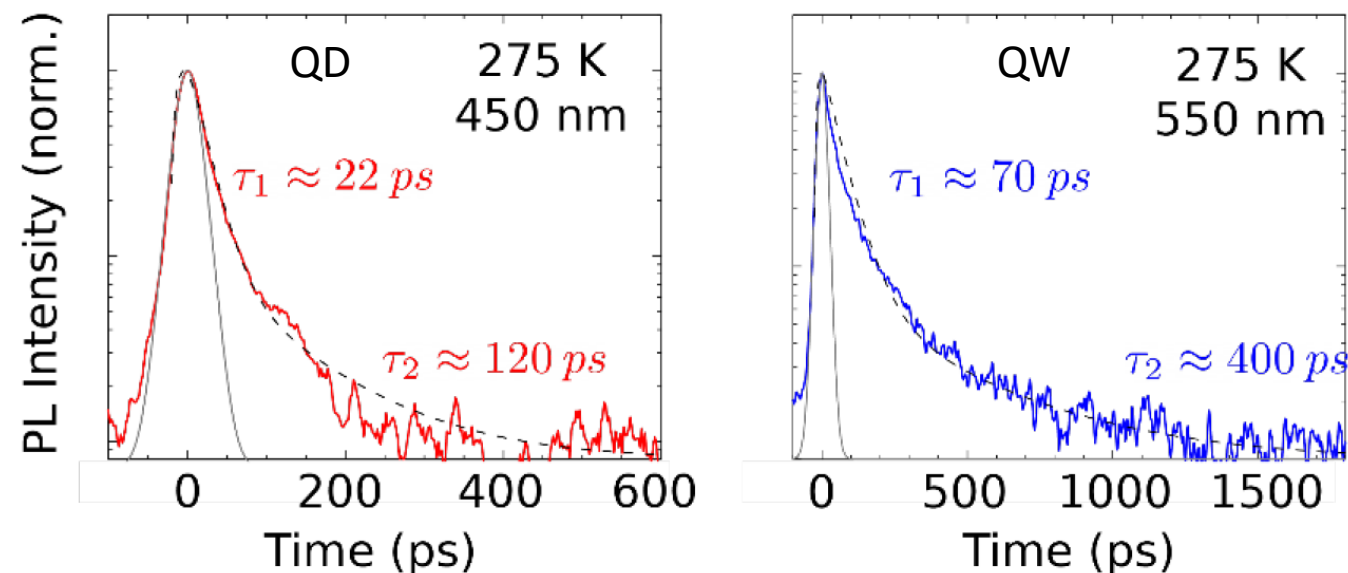
→ QD → 450 nm and short lifetimes

→ QW → 550 nm and longer lifetimes

TR-CL



TR-PL



→ Comparable decays measured by both TR techniques

Outline

Basics of SC physics

- * SC band structure
- Density of states
- Doping
- * Interaction with light
- * Dimensionality

Optical characterization of SC structures

- * Absorption
- * Photoemission
- * Reflectivity
- * Photoluminescence (PL)
 - PL Eexcitation (PLE)
- * Time-Resolved PL
 - TR-Cathodoluminescence

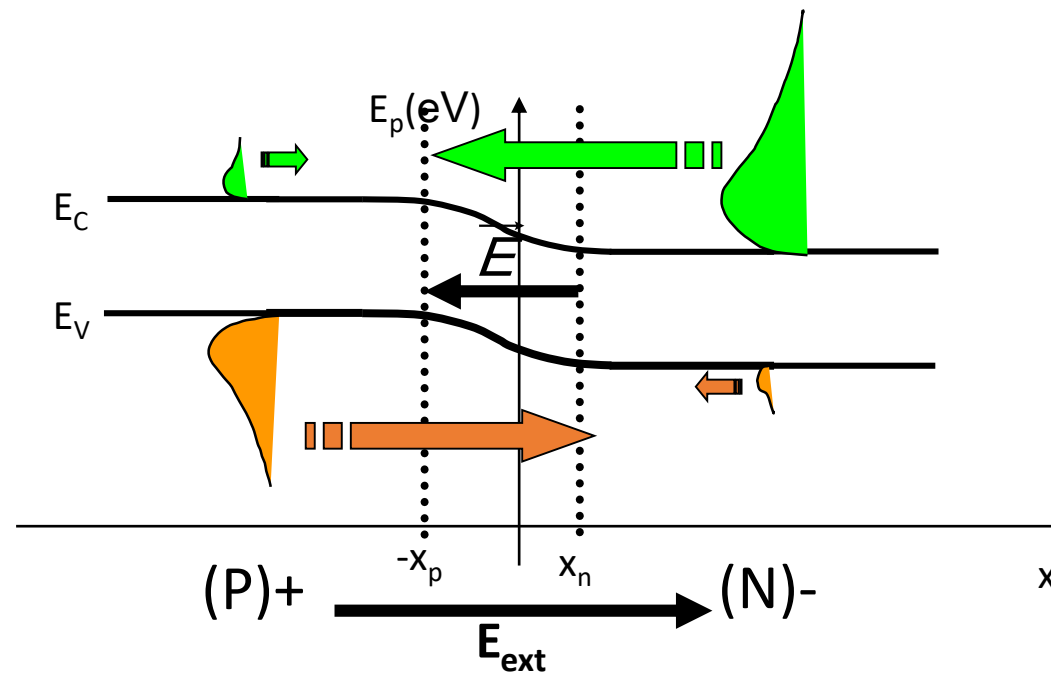
Electric/Optical characterization of SC structures

- * Electroluminescence (EL)
- * Photocurrent (PC)

Raman spectroscopy

Electro-optical measurements

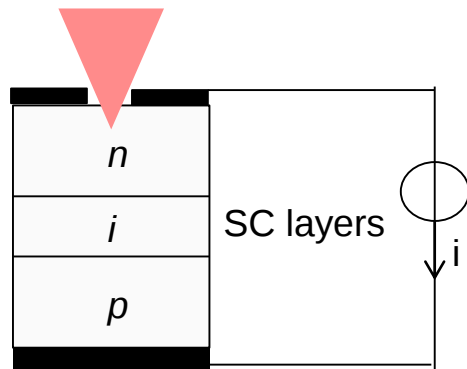
P(i)N junction



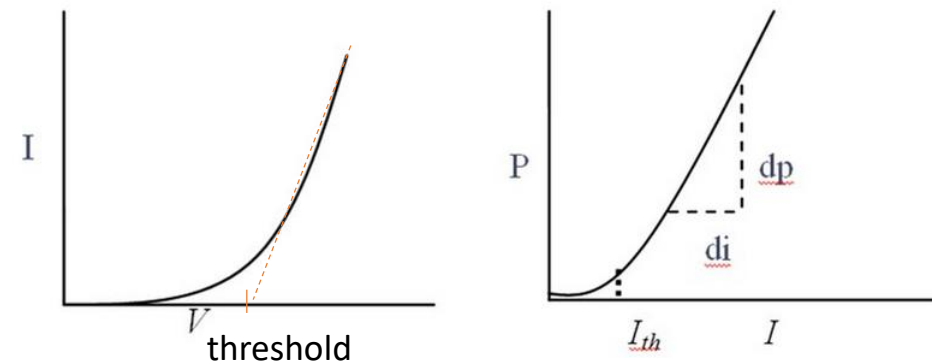
Electroluminescence

- *No need for optical excitation
- *Need for conductive samples/electric contacts

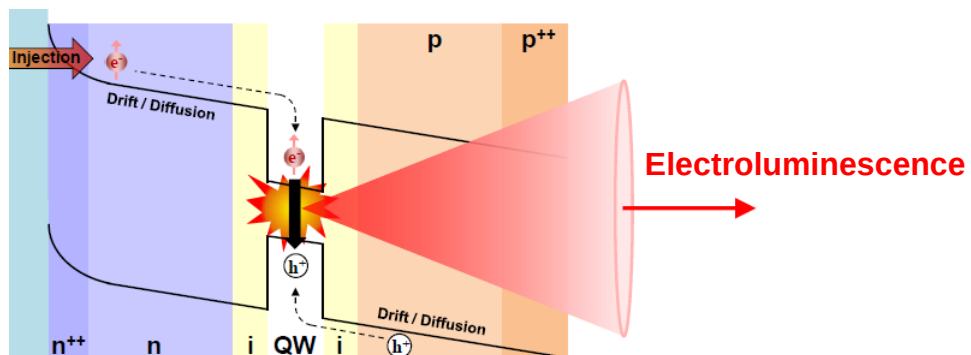
Electroluminescence



Sample



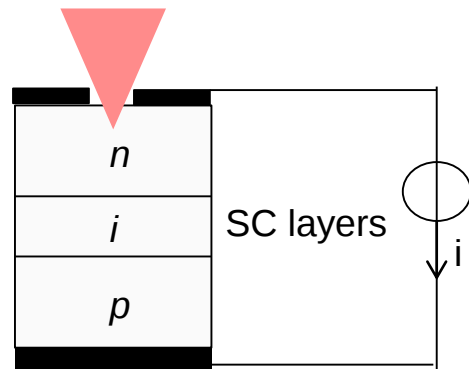
Towards optical devices: LED, laser diode,...



Electroluminescence

- *No need for optical excitation
- *Need for conductive sample/contacts

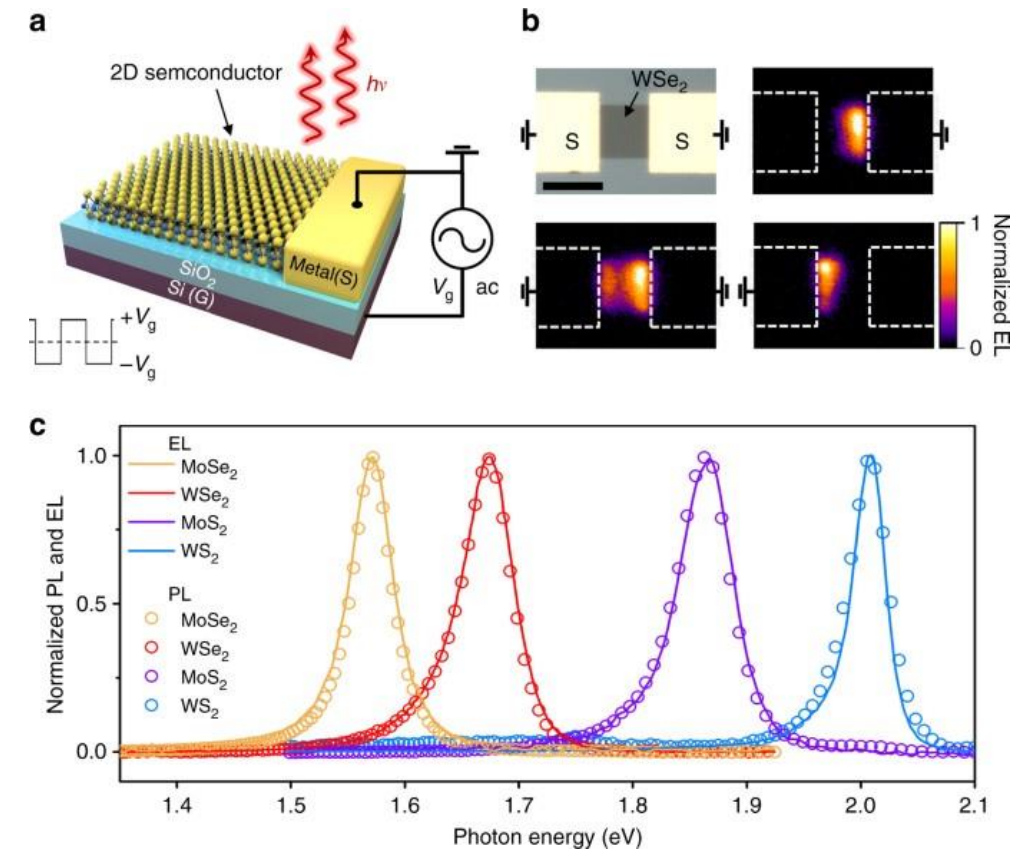
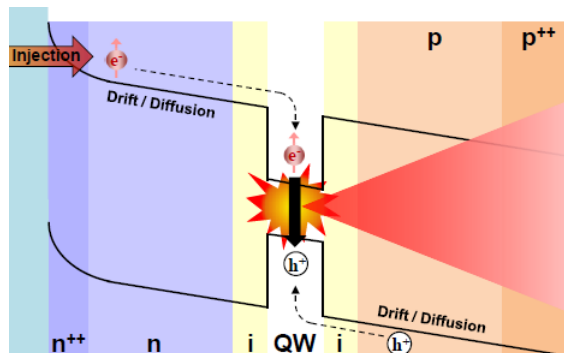
Electroluminescence



Sample



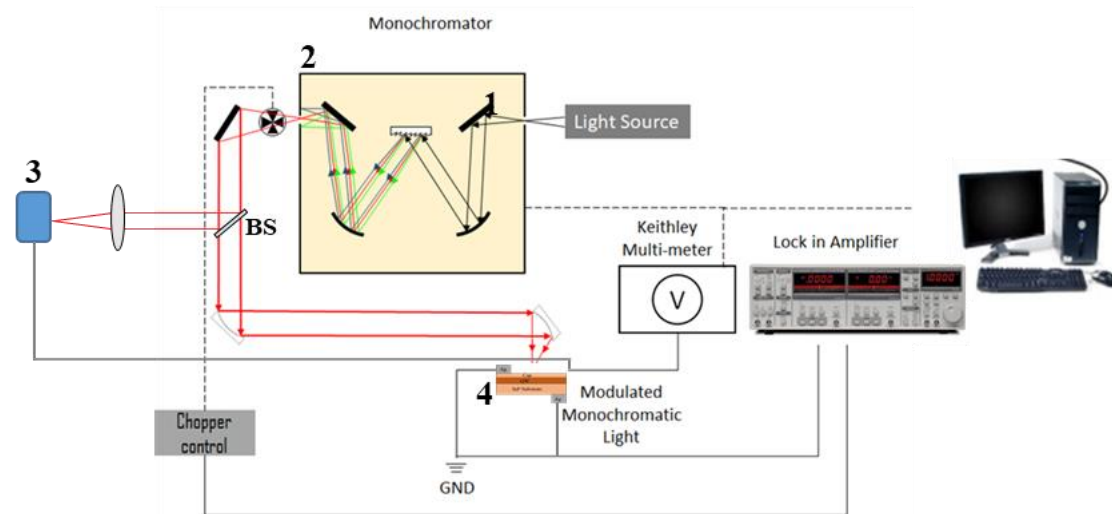
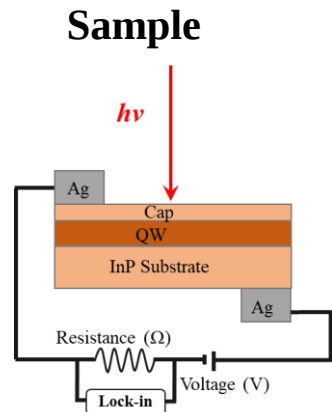
Electroluminescence



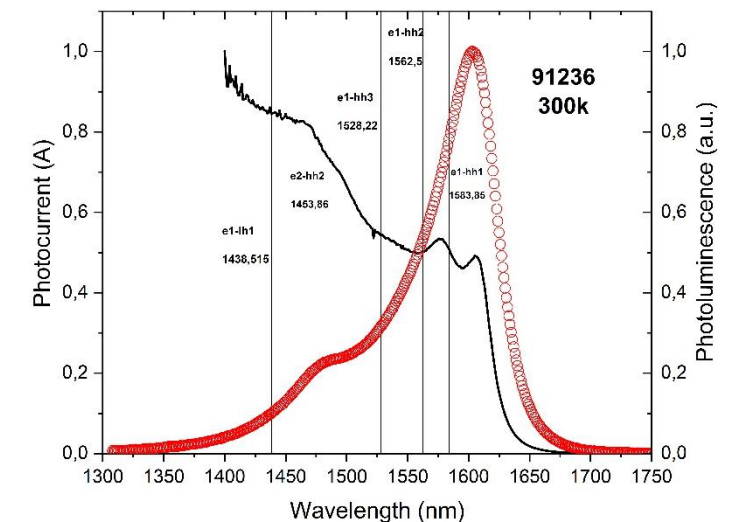
Nature (2016)

Photocurrent spectroscopy

- *No need for optical detector
- *Need for conductive samples/contacts



Comparison of PL and PC with calculated transitions



Ex: Semiconductor Optical Amplifier based on InGaAsP QWs

Even forbidden transitions can be observed due to symmetry breaking by internal electric field

Outline

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- * Photoemission
- * Reflectivity
- * Photoluminescence (PL)
 - PL Eexcitation (PLE)
- * Time-Resolved PL
 - TR-Cathodoluminescence

Electric/Optical characterization of SC structures

- * Electroluminescence (EL)
- * Photocurrent (PC)

Raman spectroscopy

Raman spectroscopy

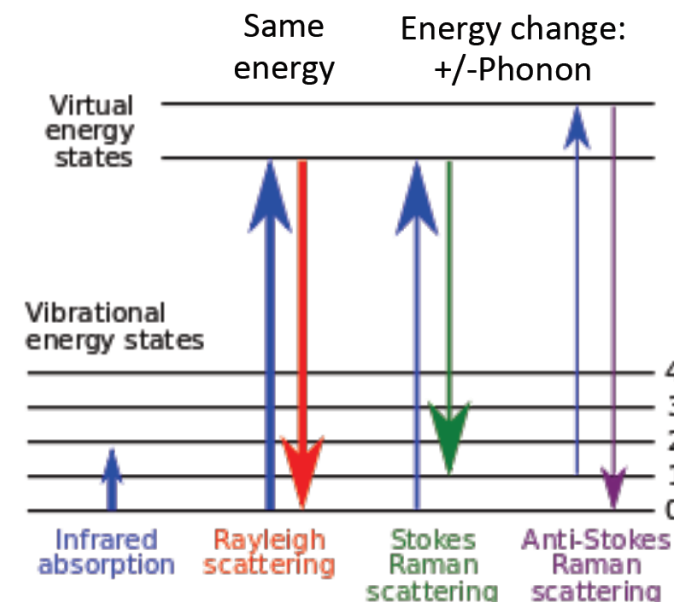
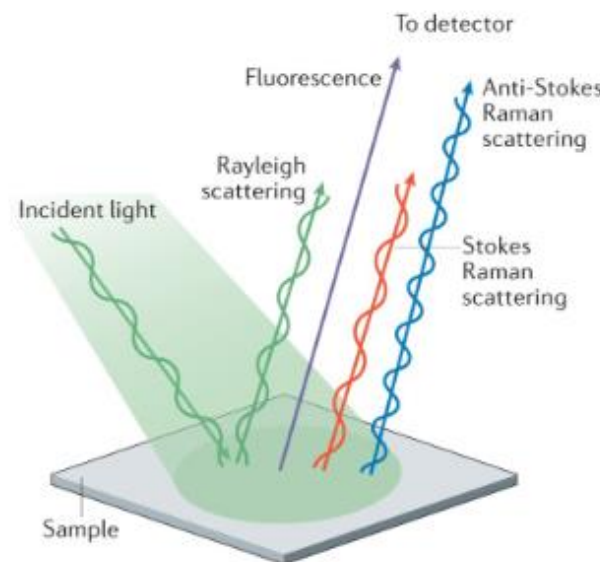
Raman effect:

Inelastic scattering of light under emission or absorption of phonons

Benefits of Raman spectroscopy

Useful in solid state physics, chemistry, biology:

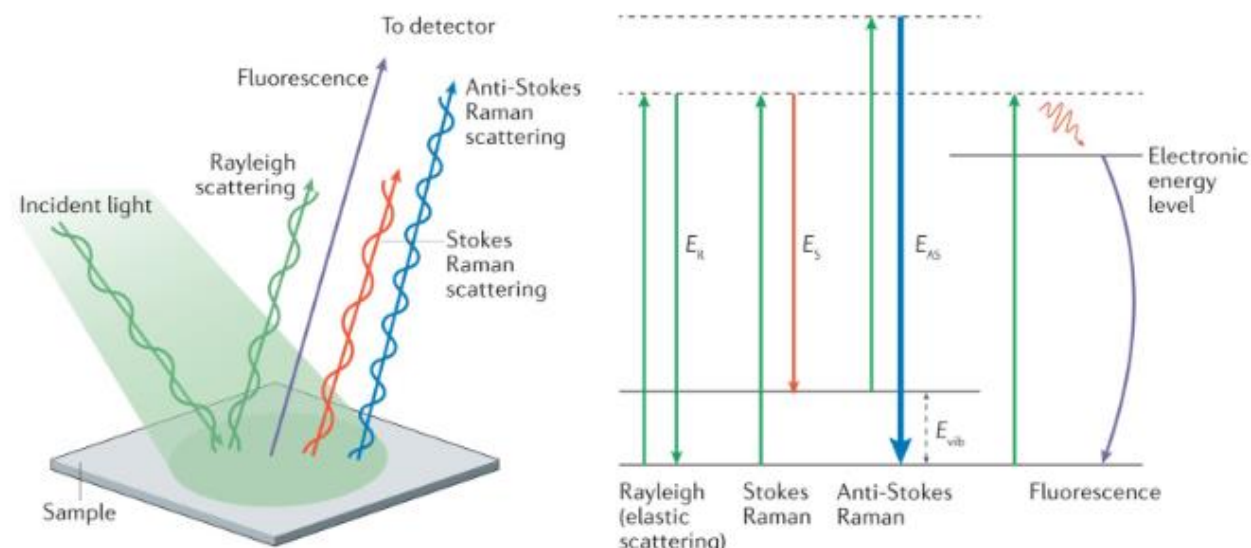
- Crystallographic phase
- Molecule identification
- Chemical bonds analysis
- Characteristic phonon modes
- ...



See Anna Fontcuberta i Morral presentation – PULSE Summer school 2015
Introductory Raman spectroscopy 2003

Raman spectroscopy

Experimental setup

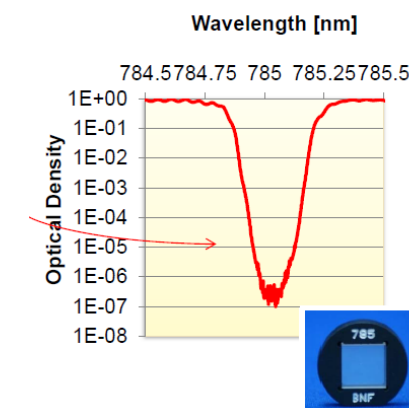
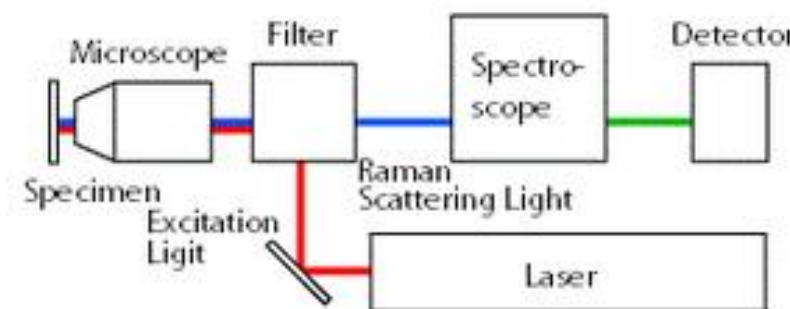
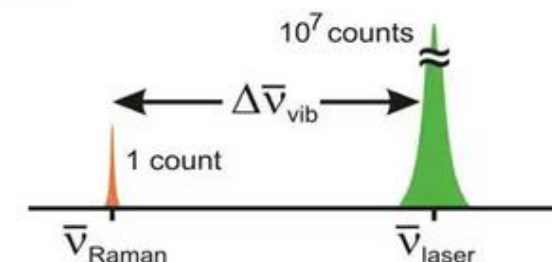


Difficulties:

- *Weak Raman signal
- *Raman signal needs to be filtered out of laser light

Unit: cm^{-1}

$$\frac{\text{Stokes Raman scattering}}{\text{Rayleigh scattering}} \approx 10^{-7}$$



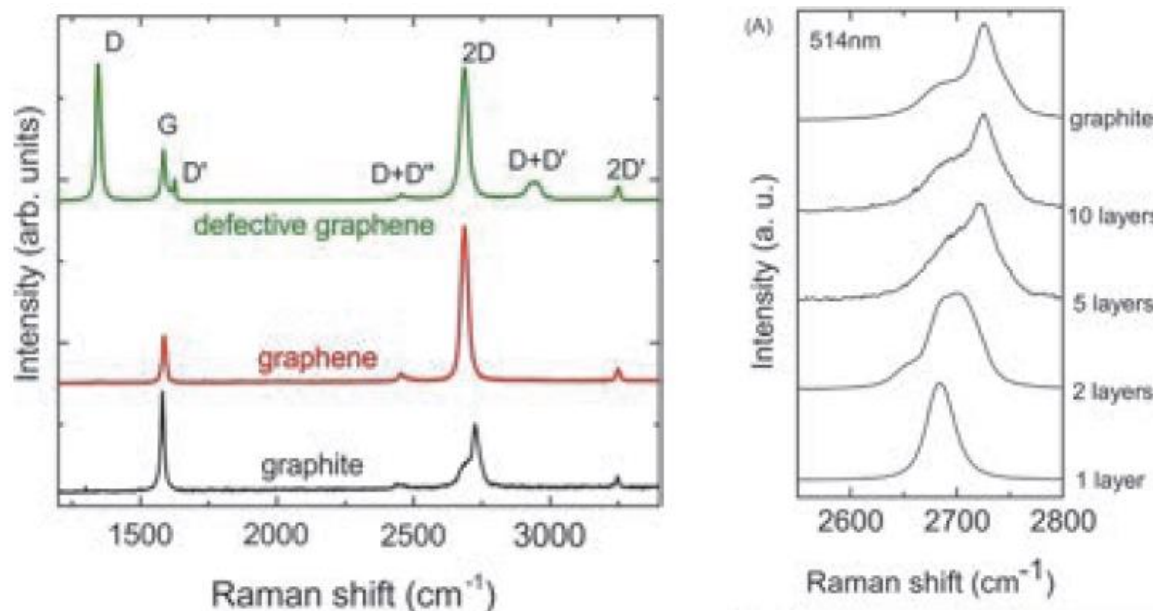
Laser: narrow line (atomic/cavity filtered)

Notch Filter/Spectrometer
with high rejection rates

Raman spectroscopy

On 2D materials

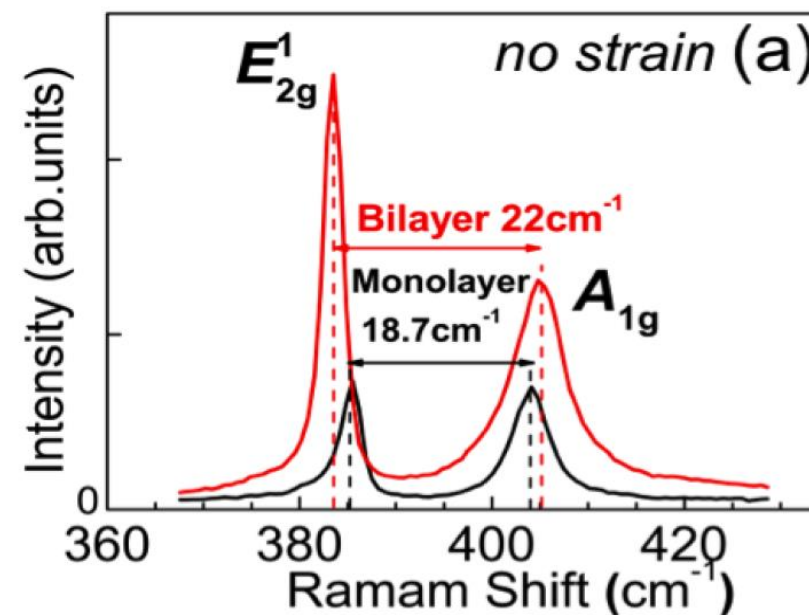
On graphene



Ferrari et al., Nature Nano 8, 235 (2013)

➡ Determination of number of graphene layers

On TMDC



Appl. Phys. Lett. 99, 102109 (2011)

Conclusions

Review of some common techniques from the « optical spectroscopy toolbox »:

Absorption, Reflectivity, Transmission, Photoluminescence

Identify intrinsic and extrinsic optical properties of a SC structure playing with tuning knobs

- Temperature
- Laser energy (=PLE)
- Laser power
- Time (=time-resolved photoluminescence)
- Doping (resident electrons or hole densities)
- Electric field
- Light polarization (excitation and detection)
- Magnetic field (in-plane, out-of-plane)
- Dielectric environment
- Strain
- Direction of propagation of light

How to exploit them for optical applications → see Fauzia Jabeen presentation on Friday

Thank you !