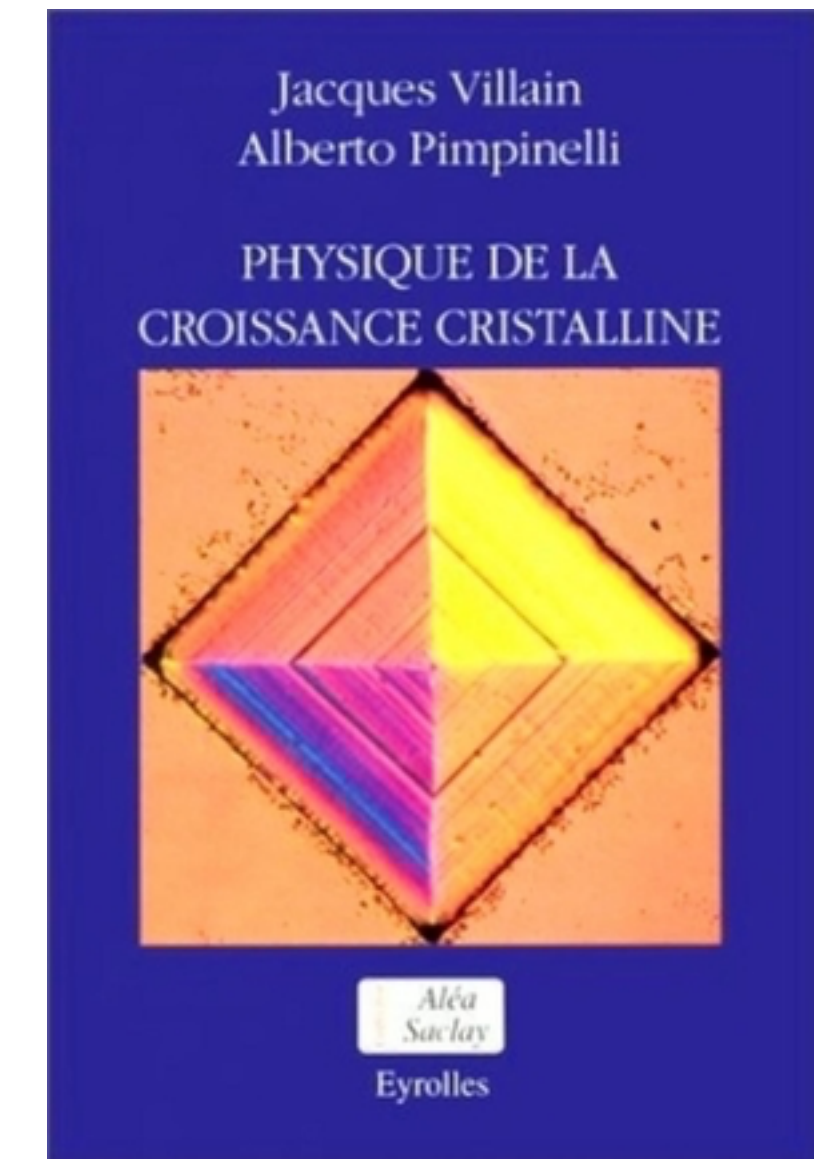
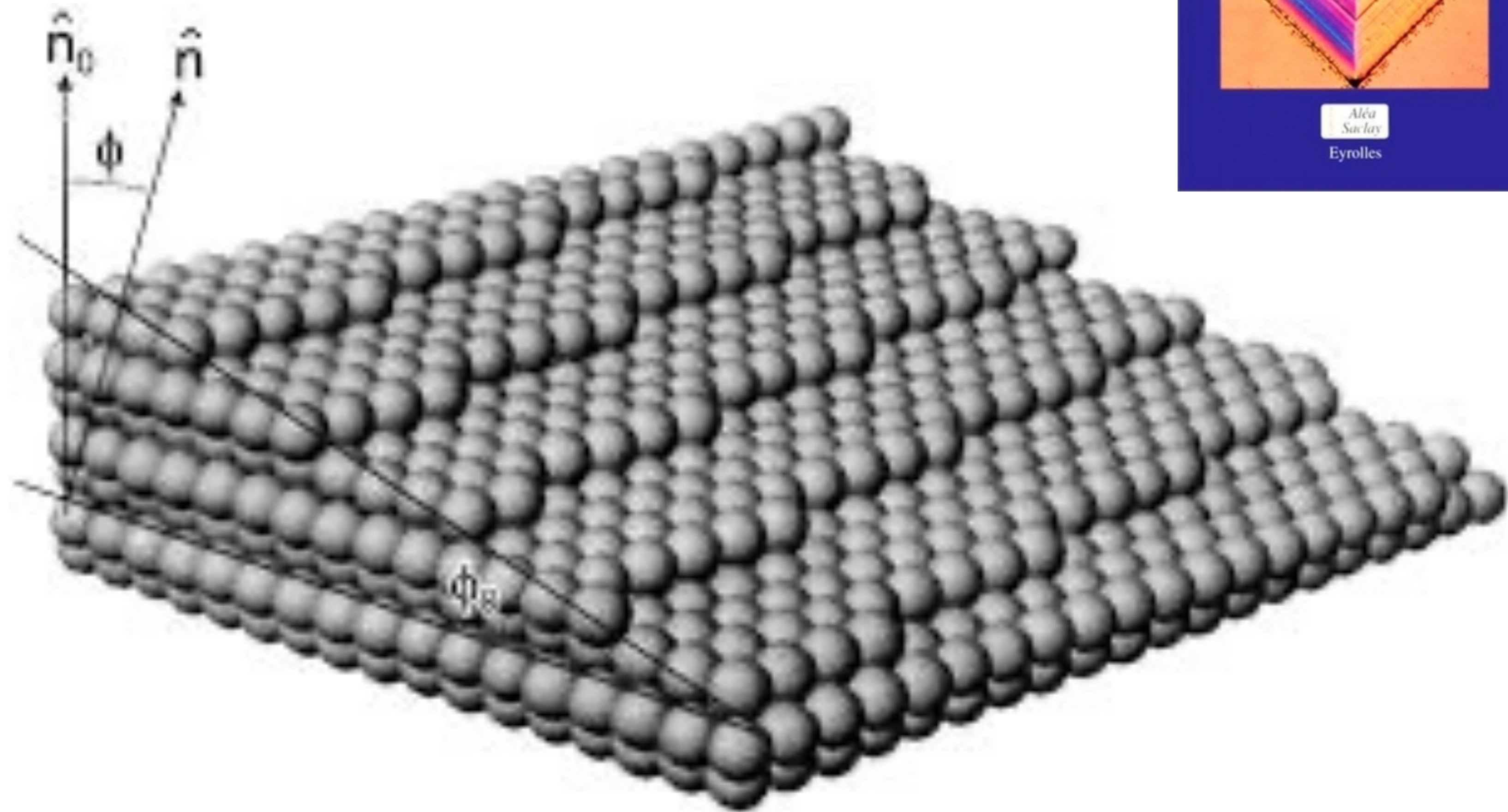
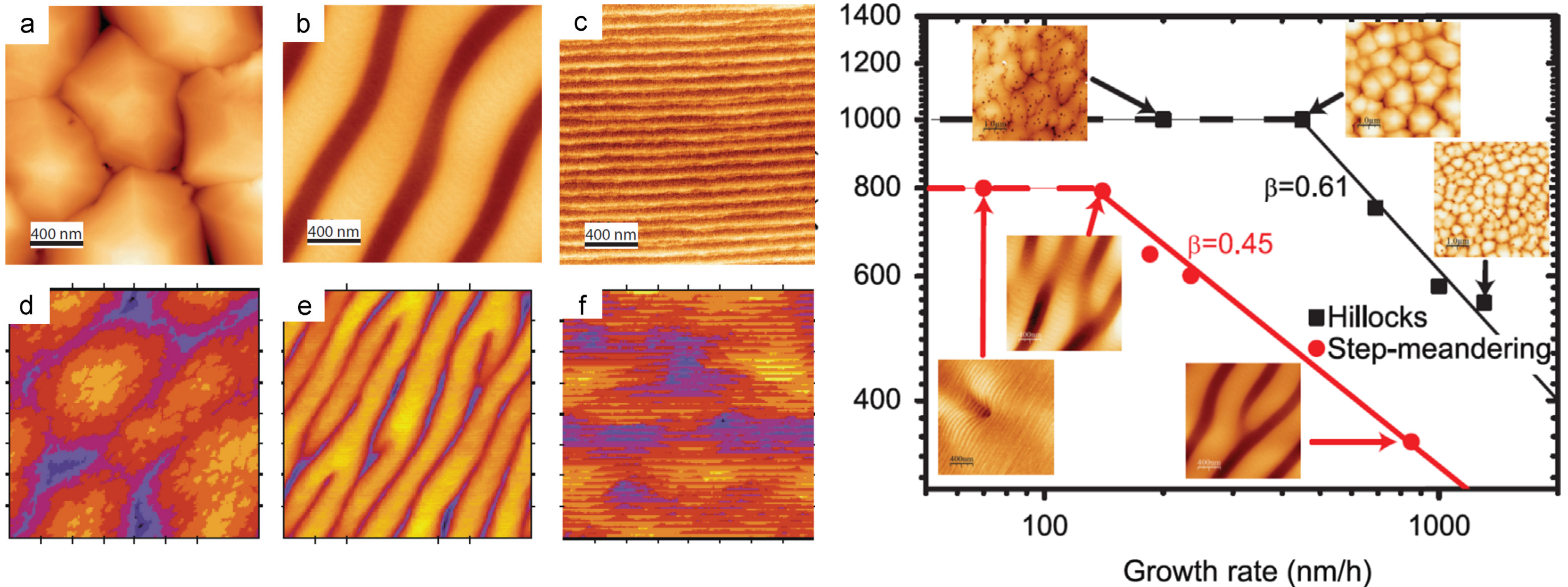


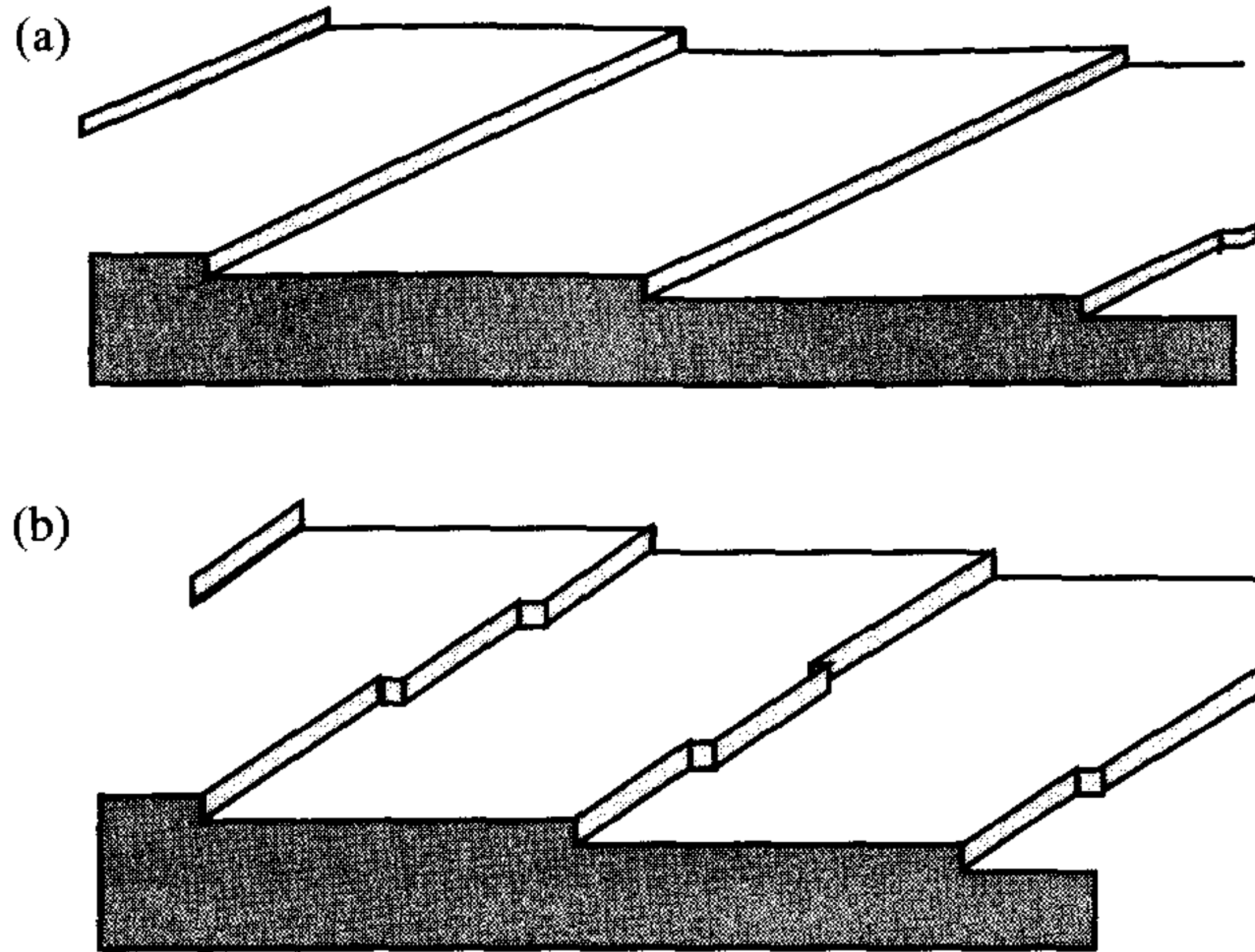
Steps, terraces, and step flow growth



Epitaxial Growth – Morphological Instabilities



Steps, terraces, and the Burton, Cabrera and Frank (BCF) model of step flow growth



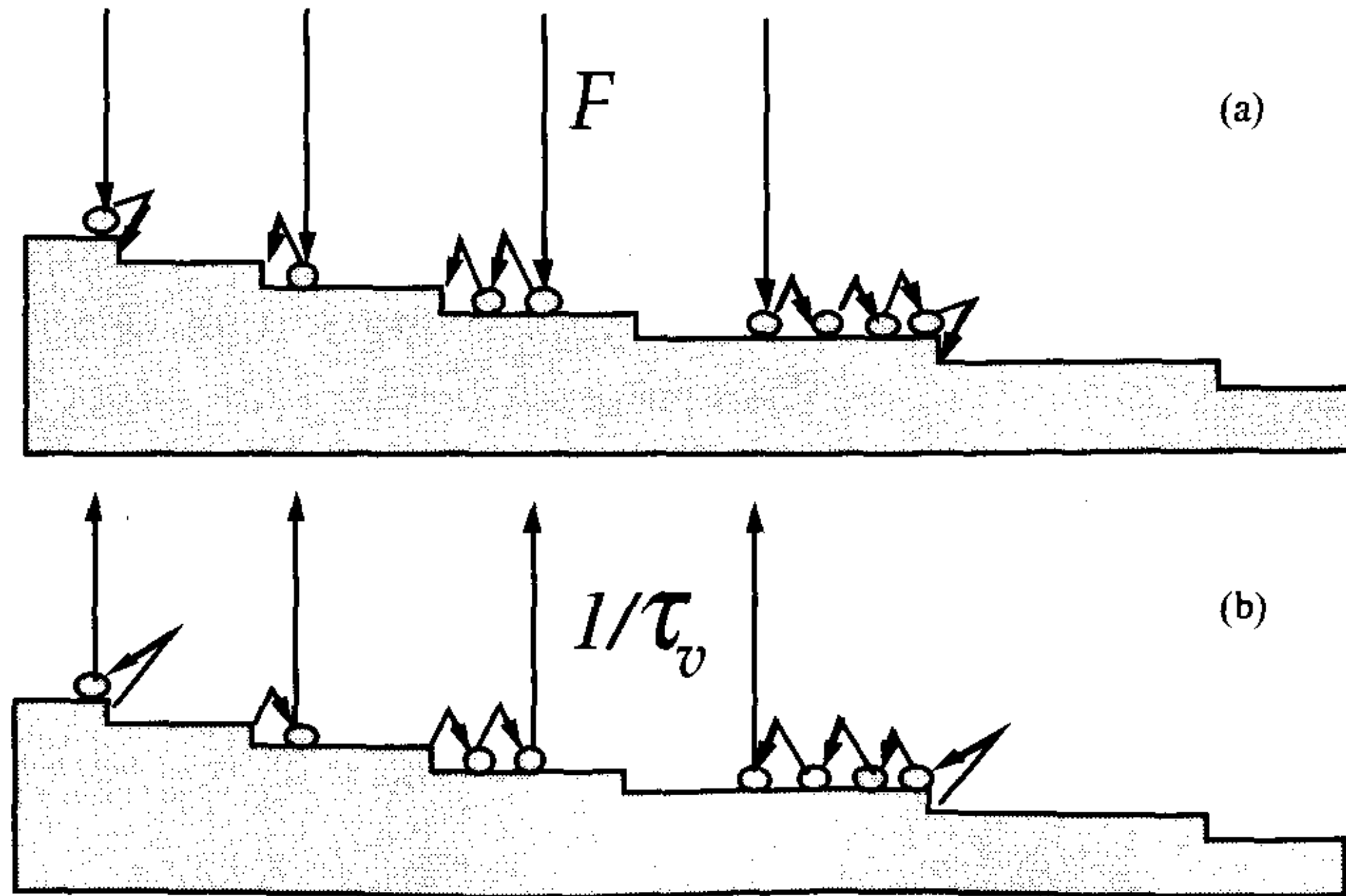


Fig. 6.1. One-dimensional picture of the growth (a) and evaporation (b) of a stepped surface.

$$\frac{\partial n}{\partial t} = D \nabla^2 n + F - \frac{n}{\tau_v}$$

$$\vec{J} = -D \nabla n$$

Boundary conditions at step edge :

$$n = n|_{step}; \quad \nabla n = \nabla n|_{step}$$

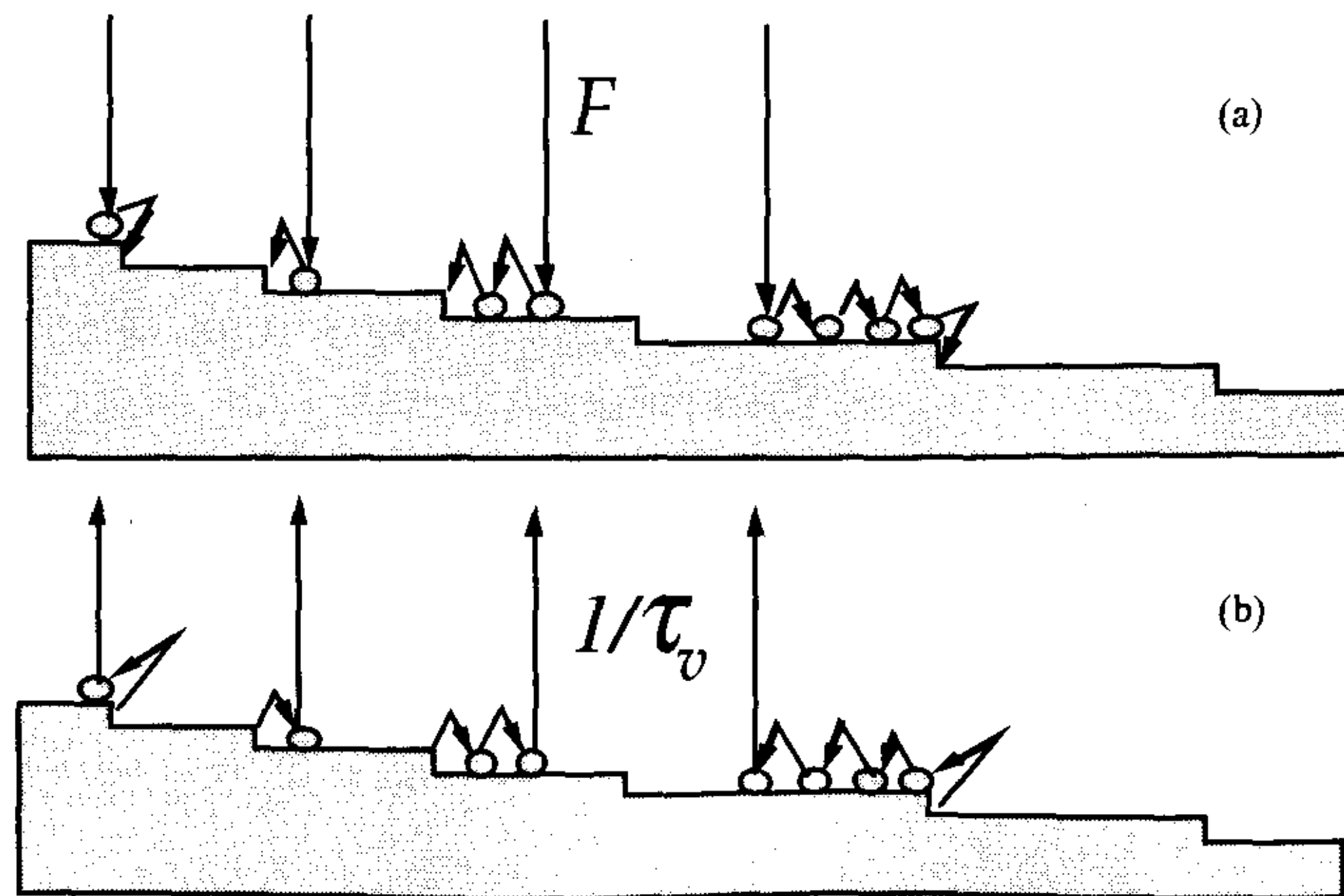


Fig. 6.1. One-dimensional picture of the growth (a) and evaporation (b) of a stepped surface.

Quasi-static approximation

$$\vec{J} = -D \nabla n$$

Adatom density computed for stationary steps

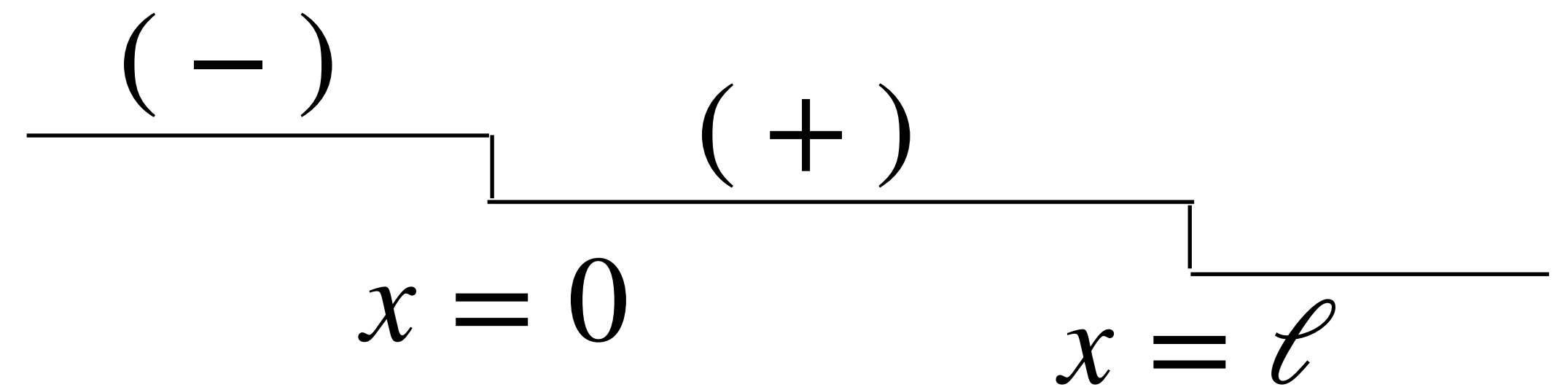
step edge : $n|_{step}$

Example – no evaporation

$$\frac{\partial n}{\partial t} = D \nabla^2 n + F$$

Boundary conditions at step edge :

$$n(0) = n(\ell) = 0$$



$$n(x) = (F/D) \frac{x(\ell - x)}{2}$$

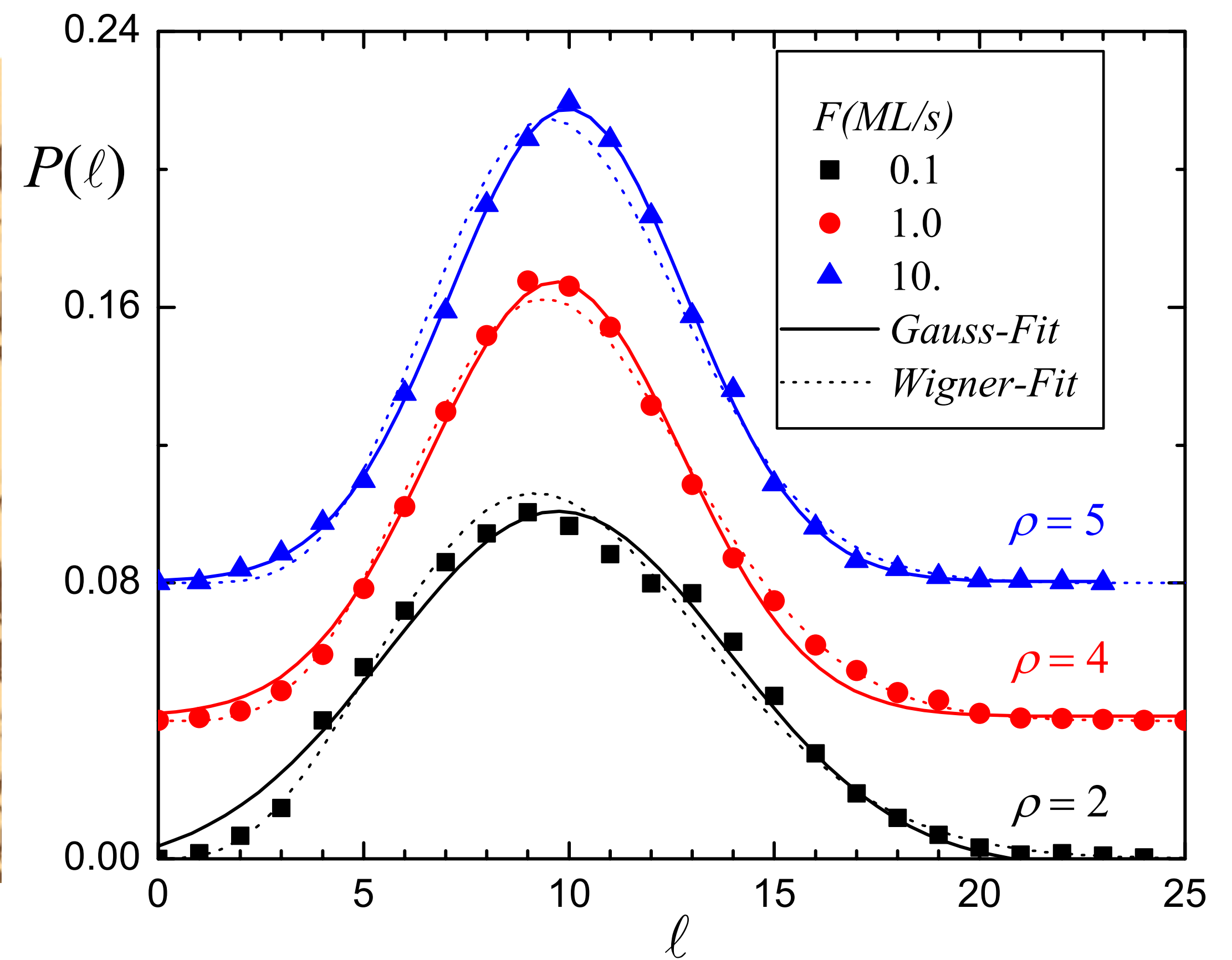
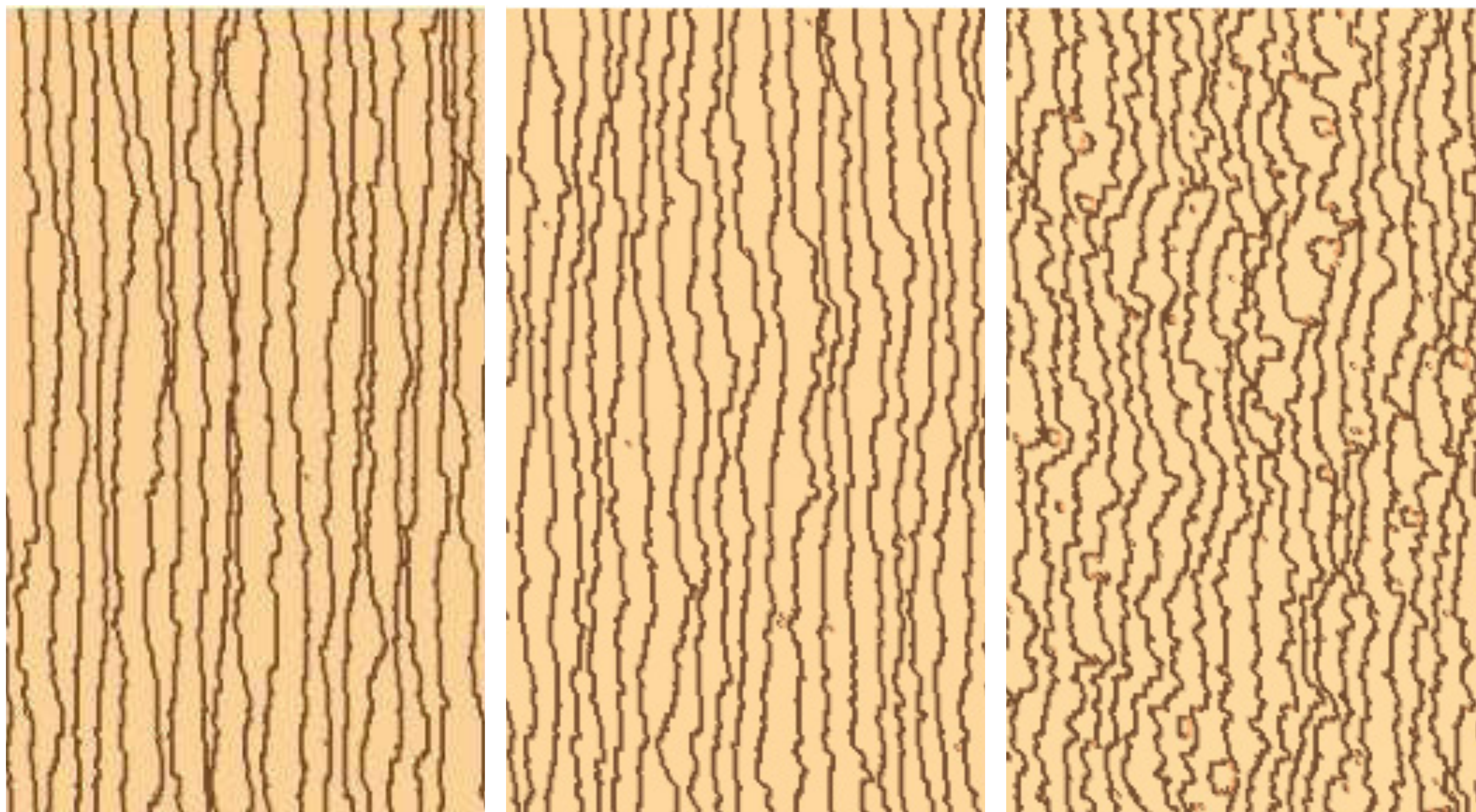
$$v_{\text{step}} = -D \frac{dn^-}{dx} \Big|_{x=0} + D \frac{dn^+}{dx} \Big|_{x=0}$$

$$v_{\text{step}} = F\ell/2 + F\ell/2 = F\ell$$

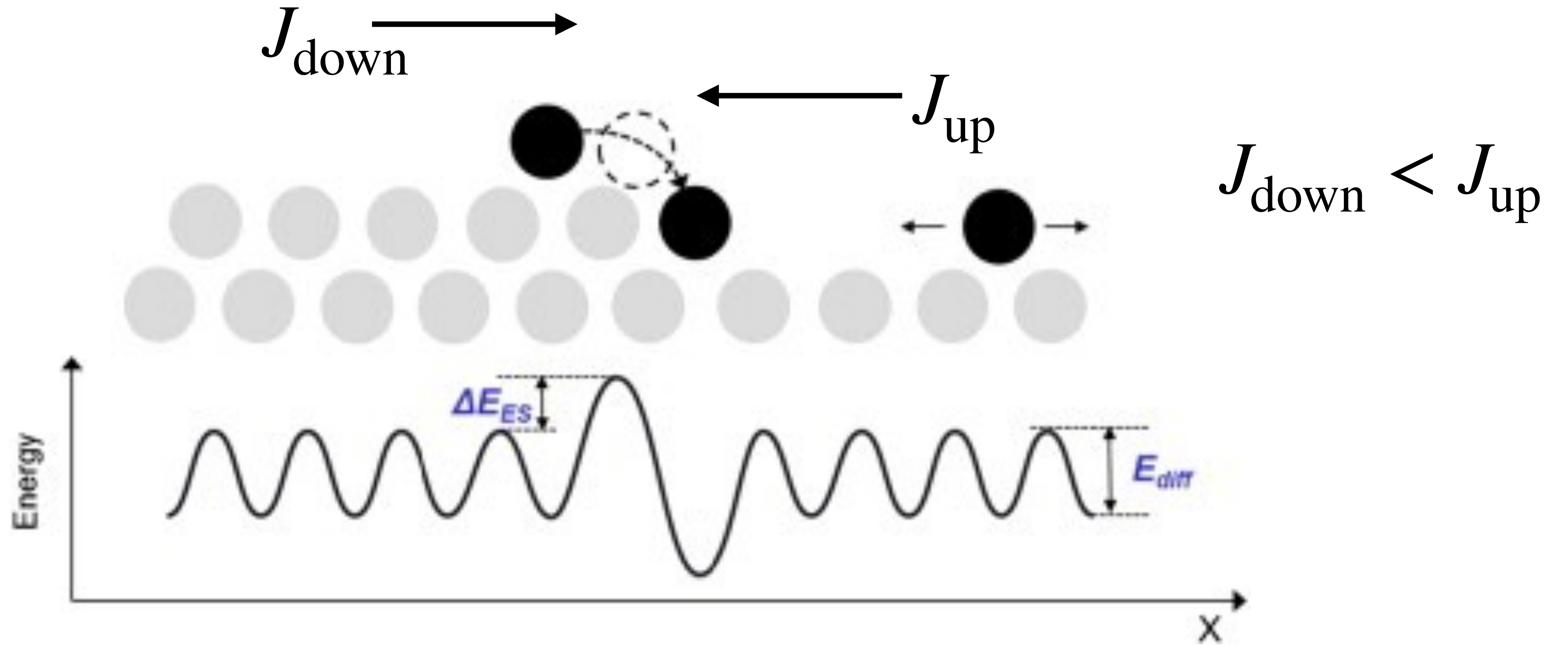
Growth rate = v_{step} $\times$ step density

$$G = F\ell \cdot 1/\ell = F$$

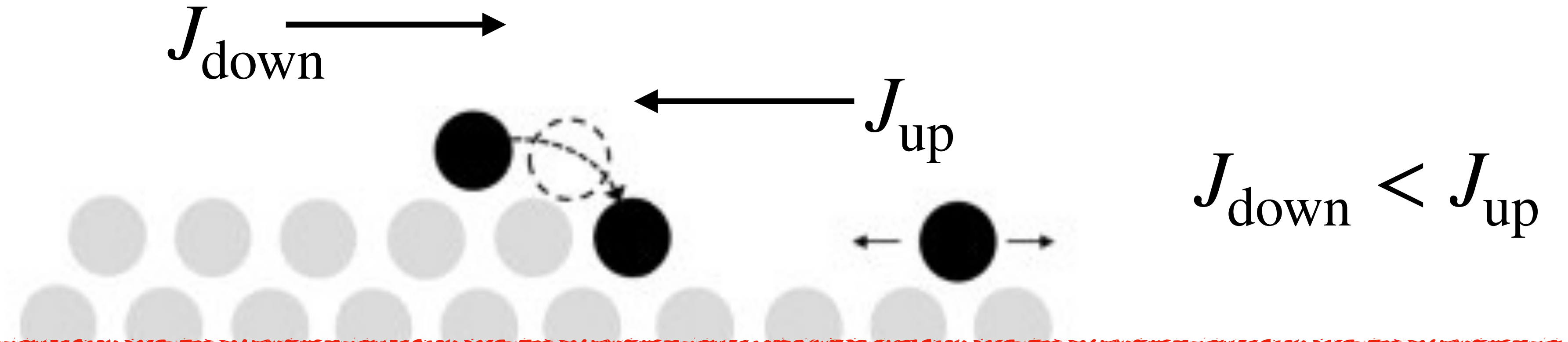
kMC — Moving Steps



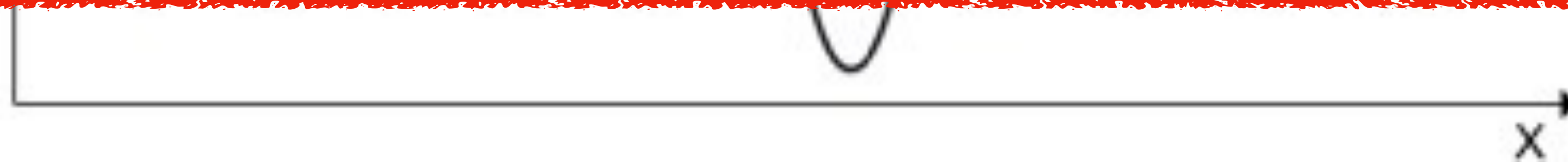
Ehrlich-Schwoebel effect



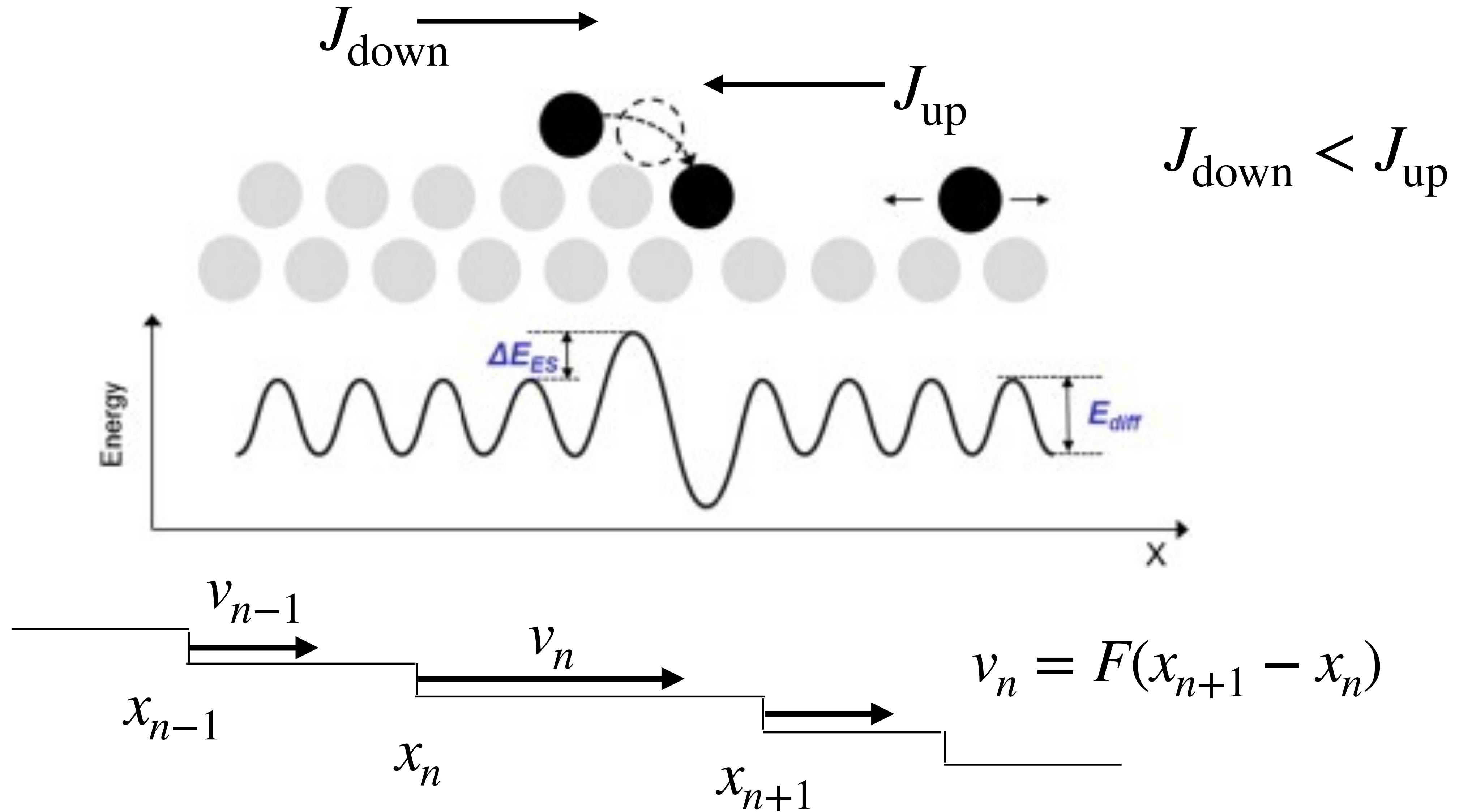
Ehrlich-Schwobel effect



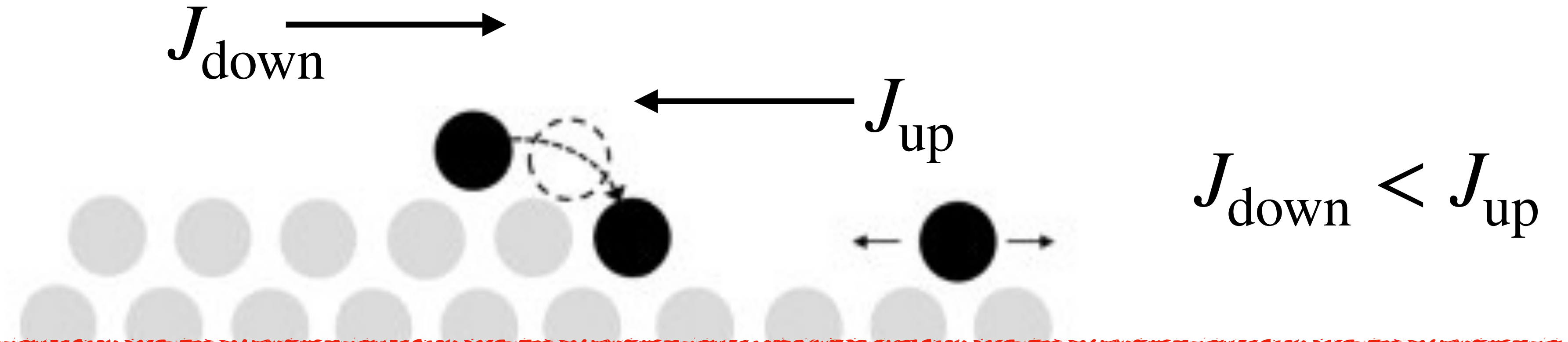
Extreme case: infinite ES effect



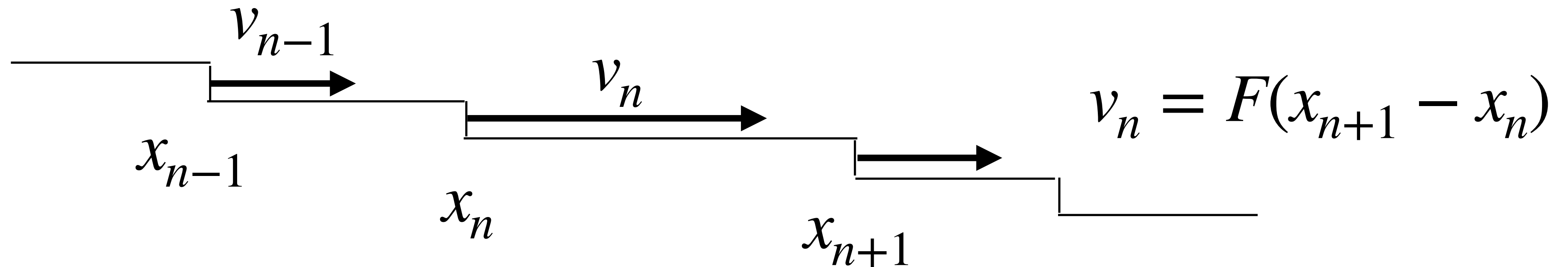
Ehrlich-Schwoebel effect



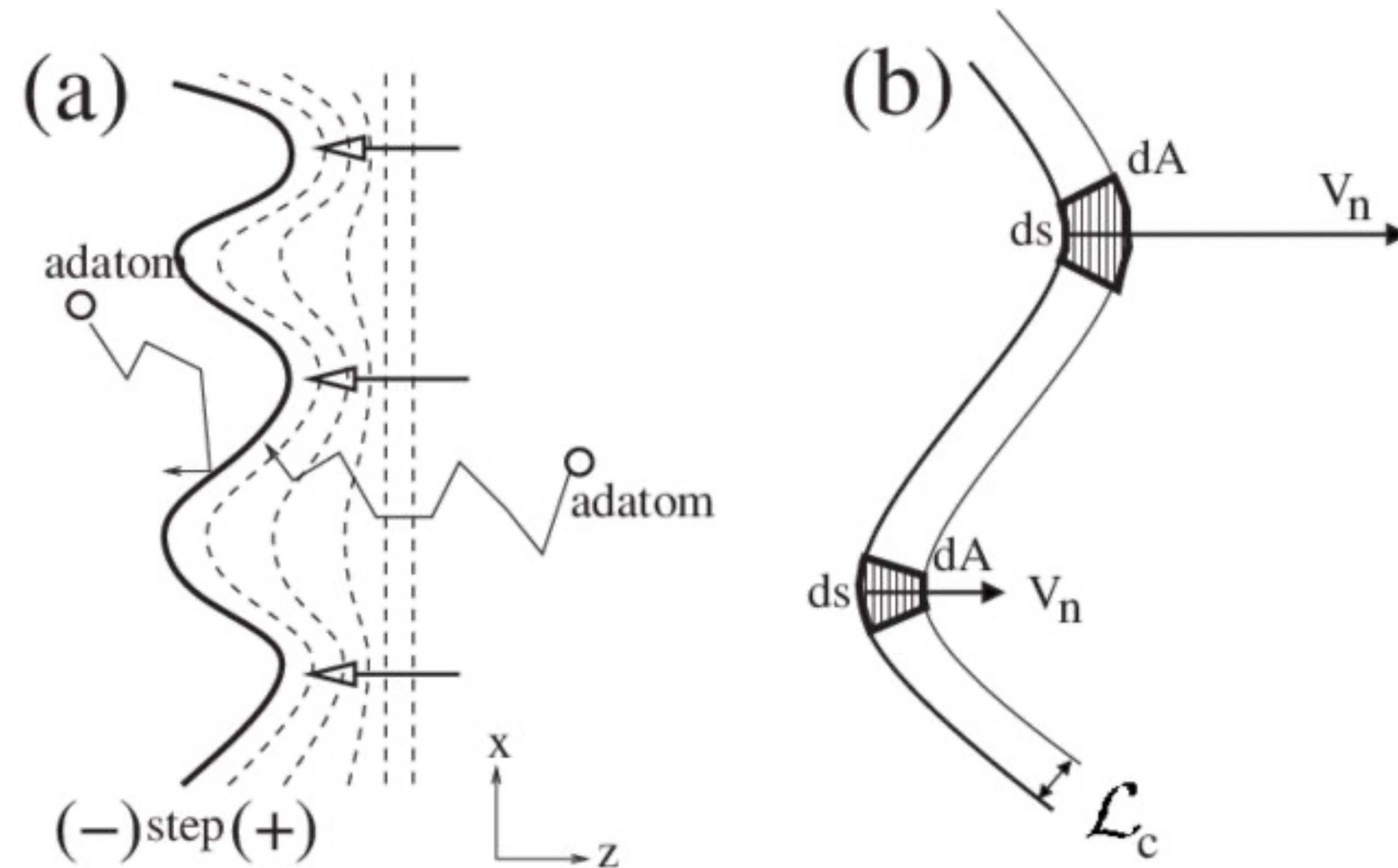
Ehrlich-Schwobel effect



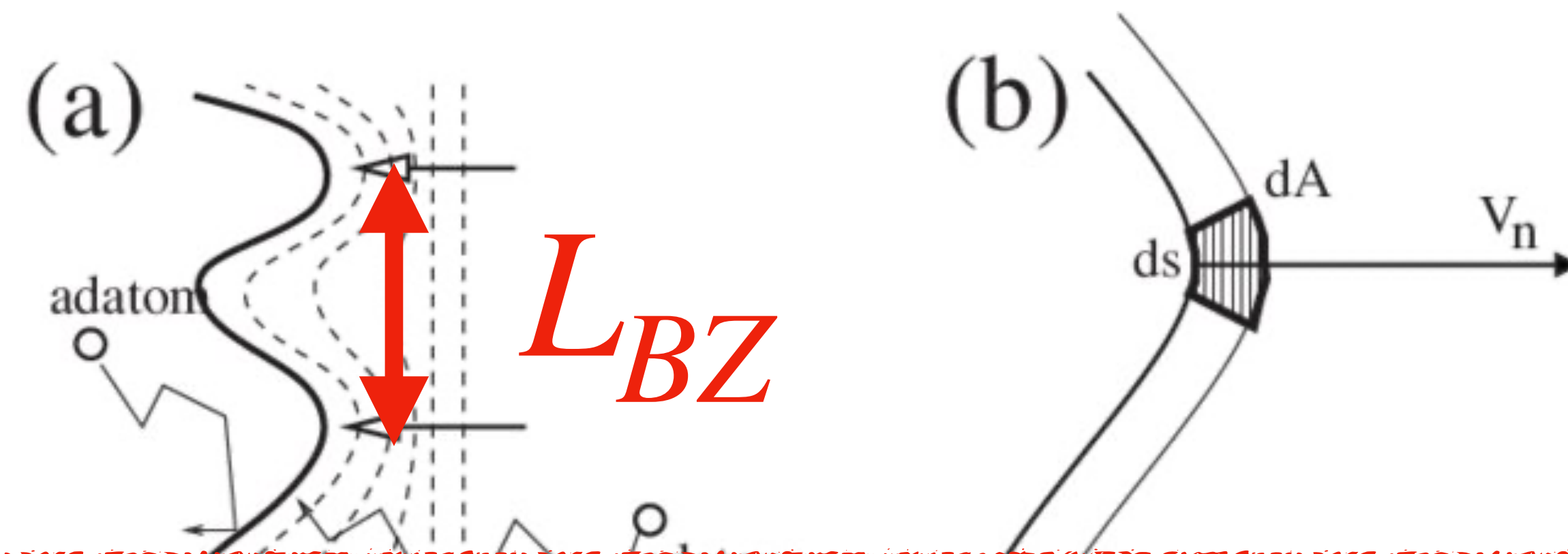
Kinetic repulsive interaction



Diffusional instability ("Schwoebel gives meanders")



Diffusional instability ("Schwoebel gives meanders")

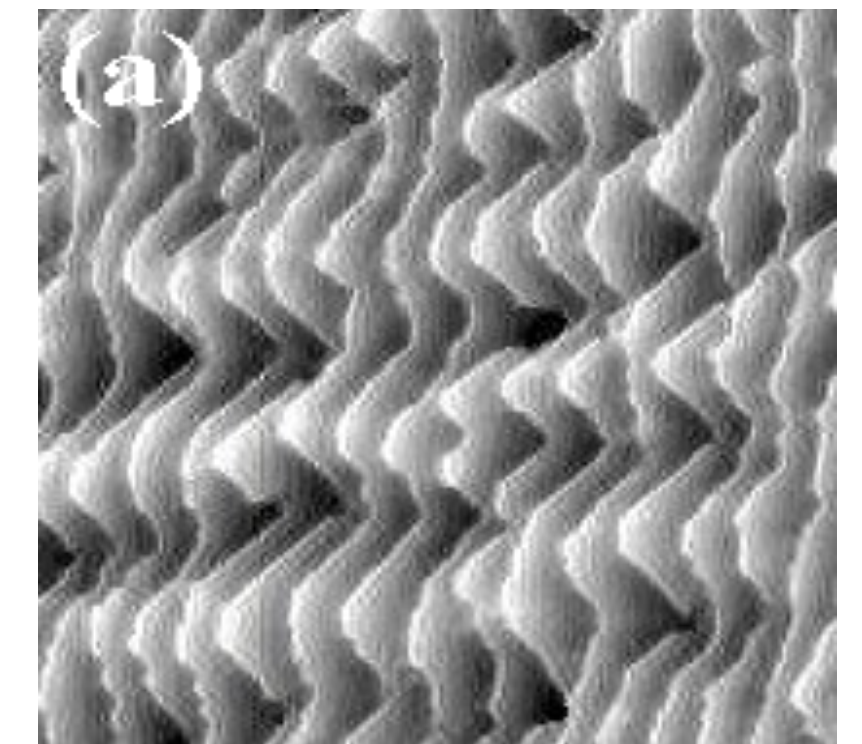
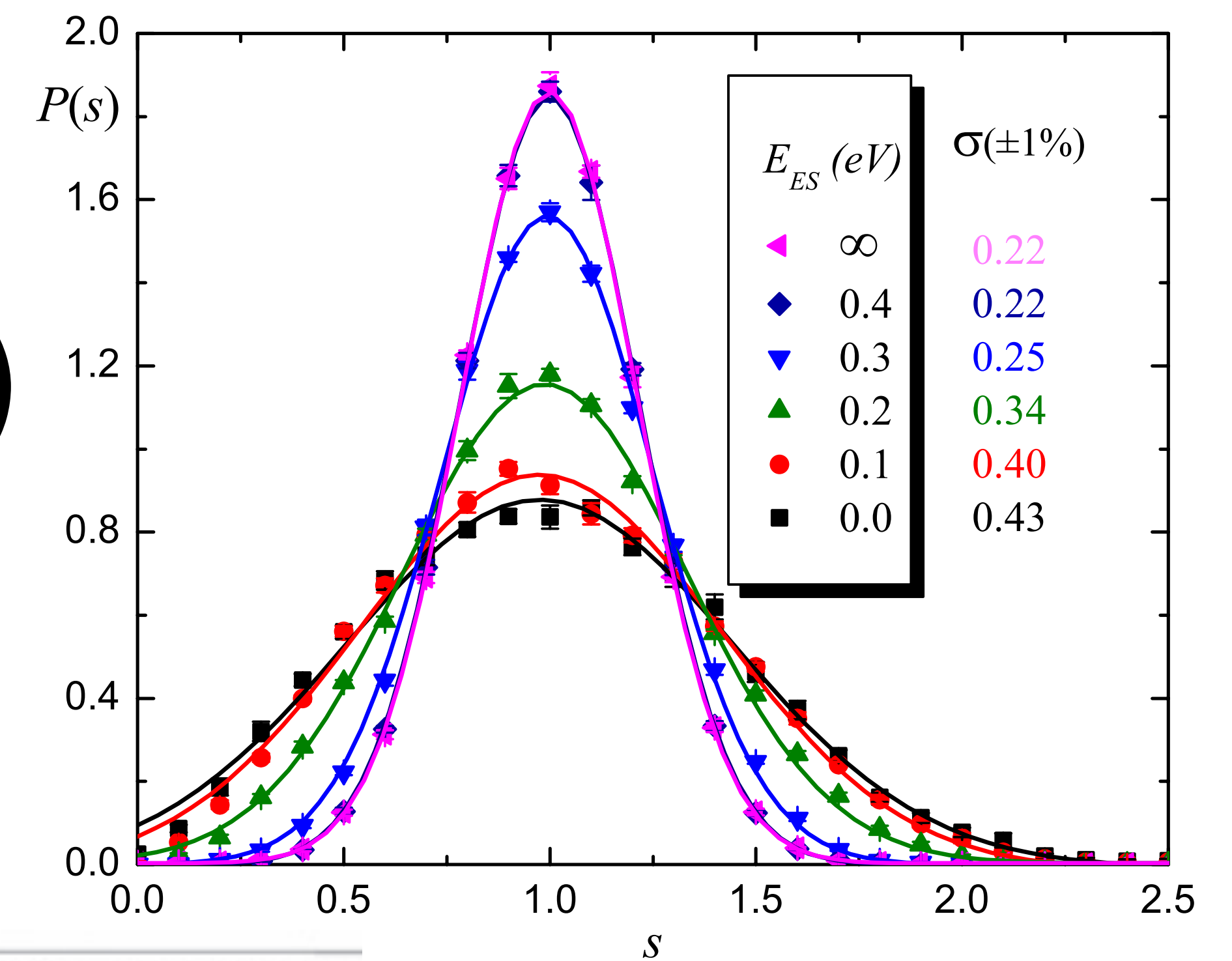
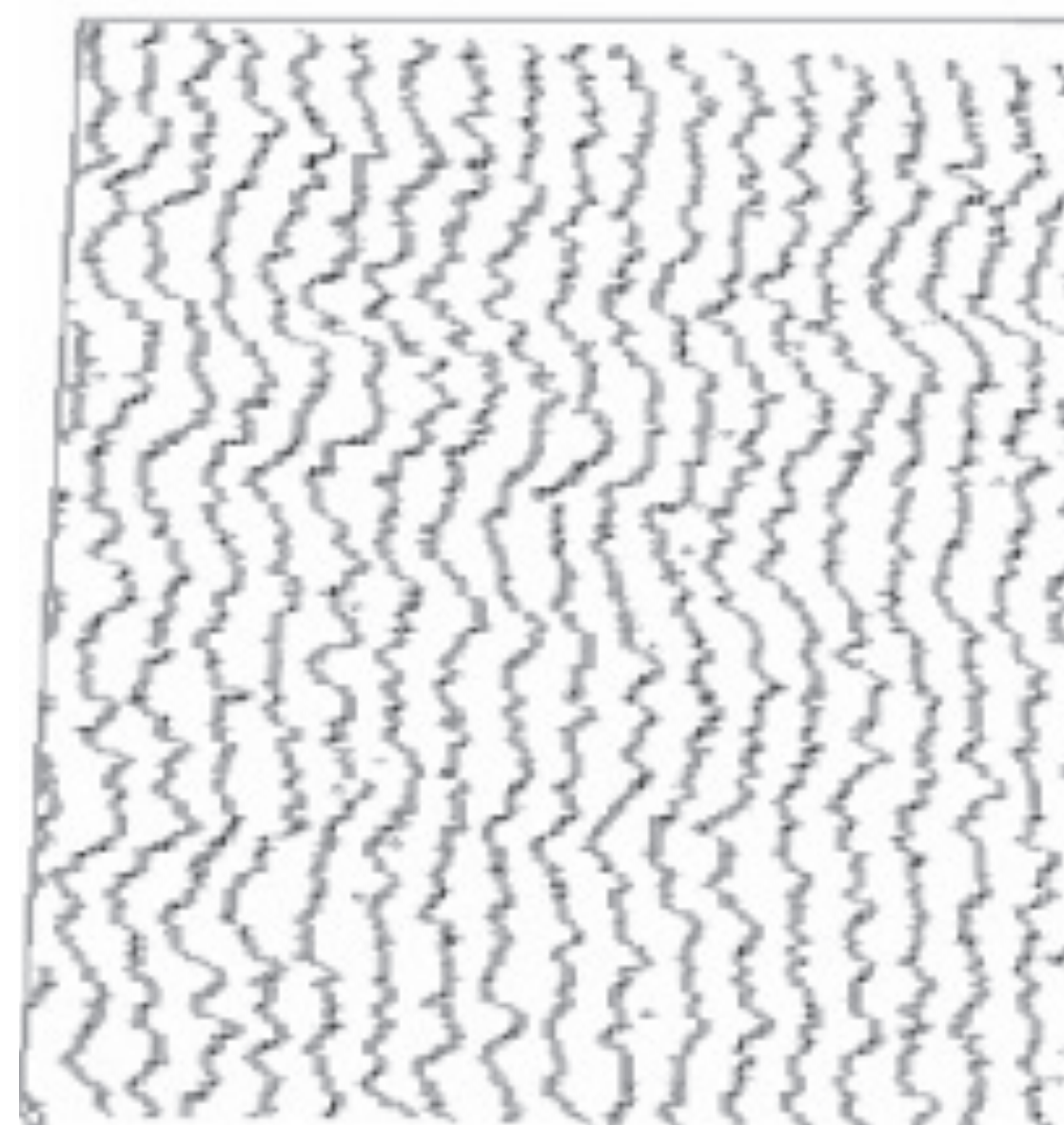
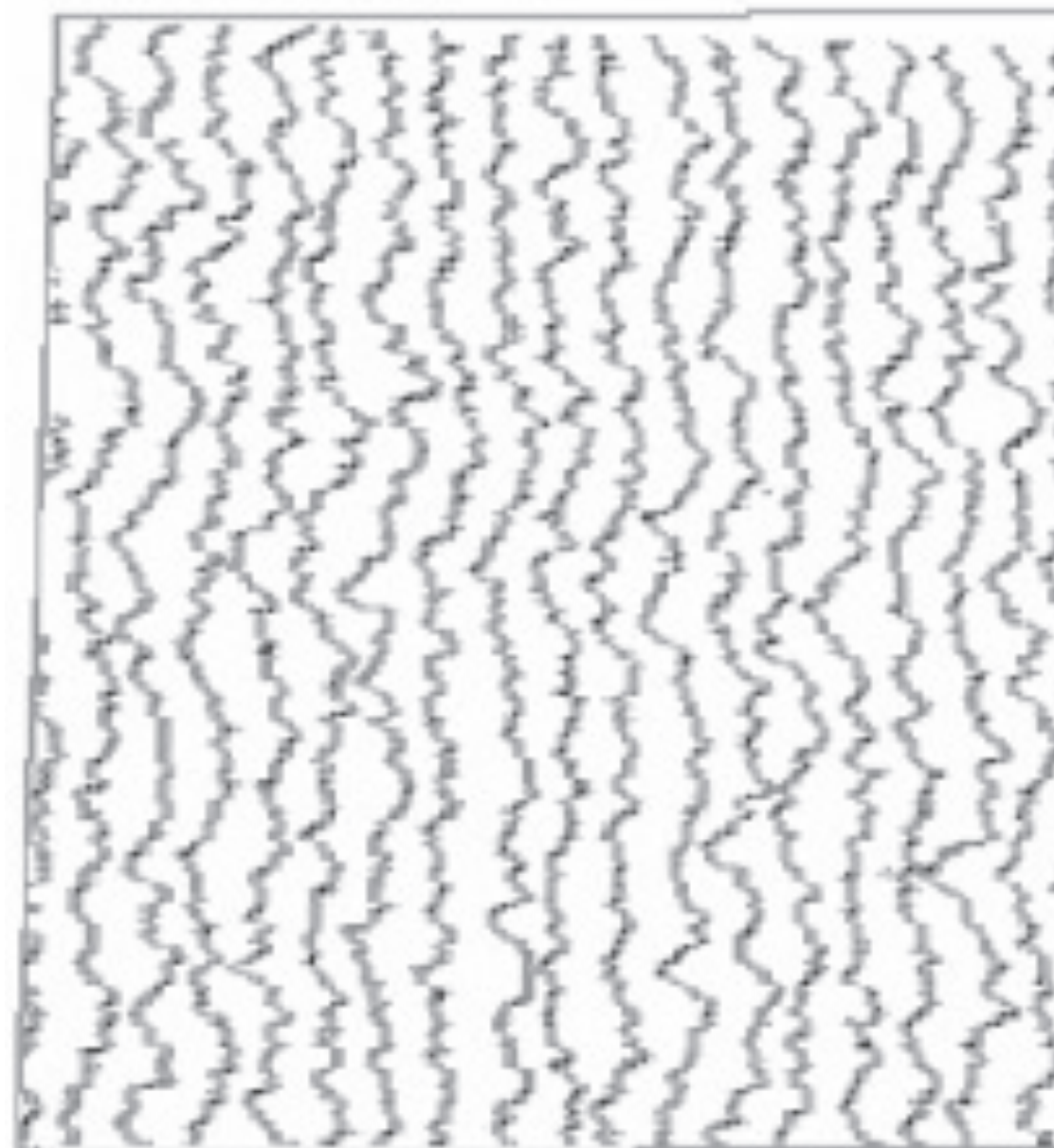
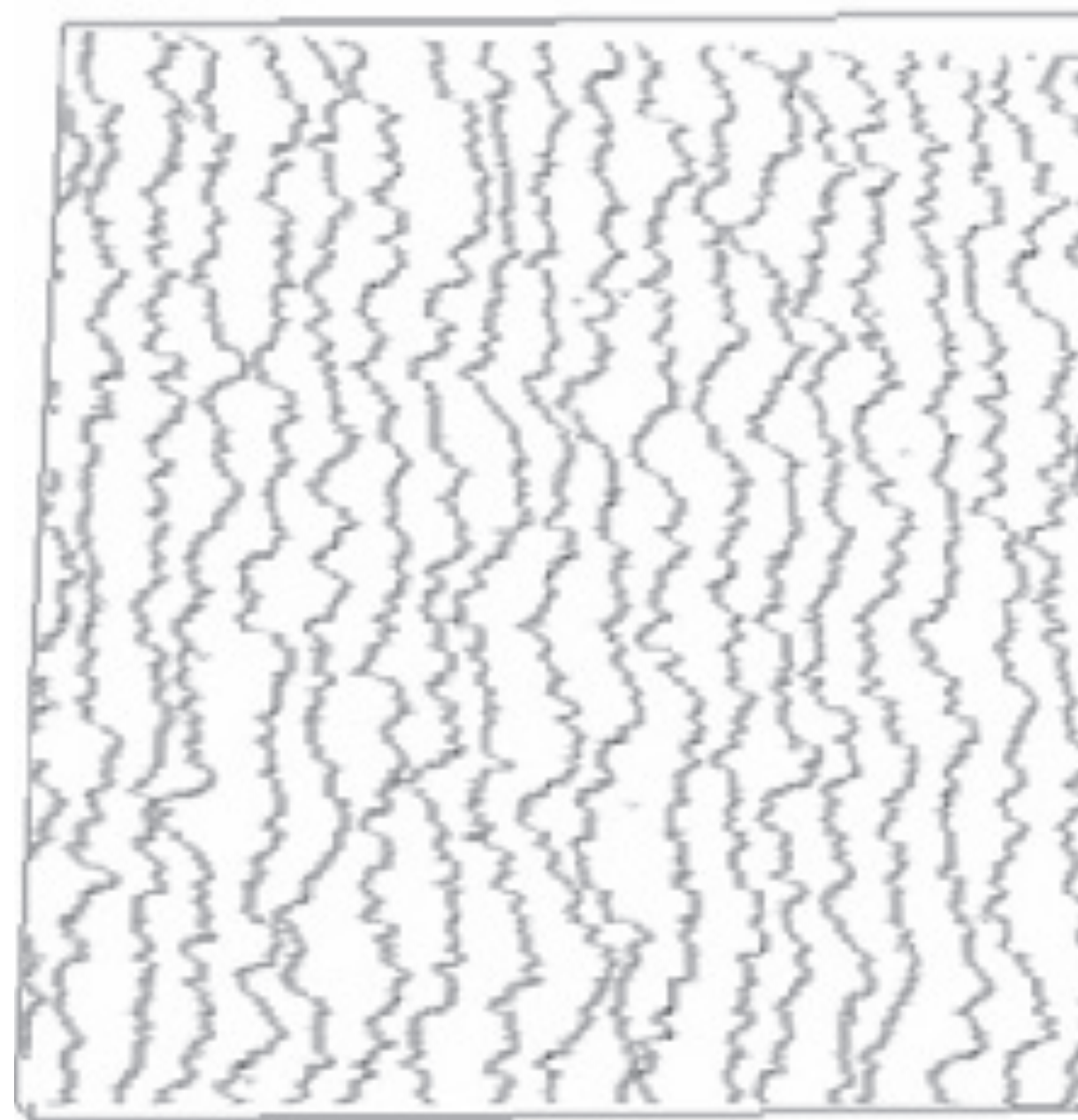


$$L_{BZ} \approx (D/F)^{1/2}$$

Bales & Zangwill, 1990

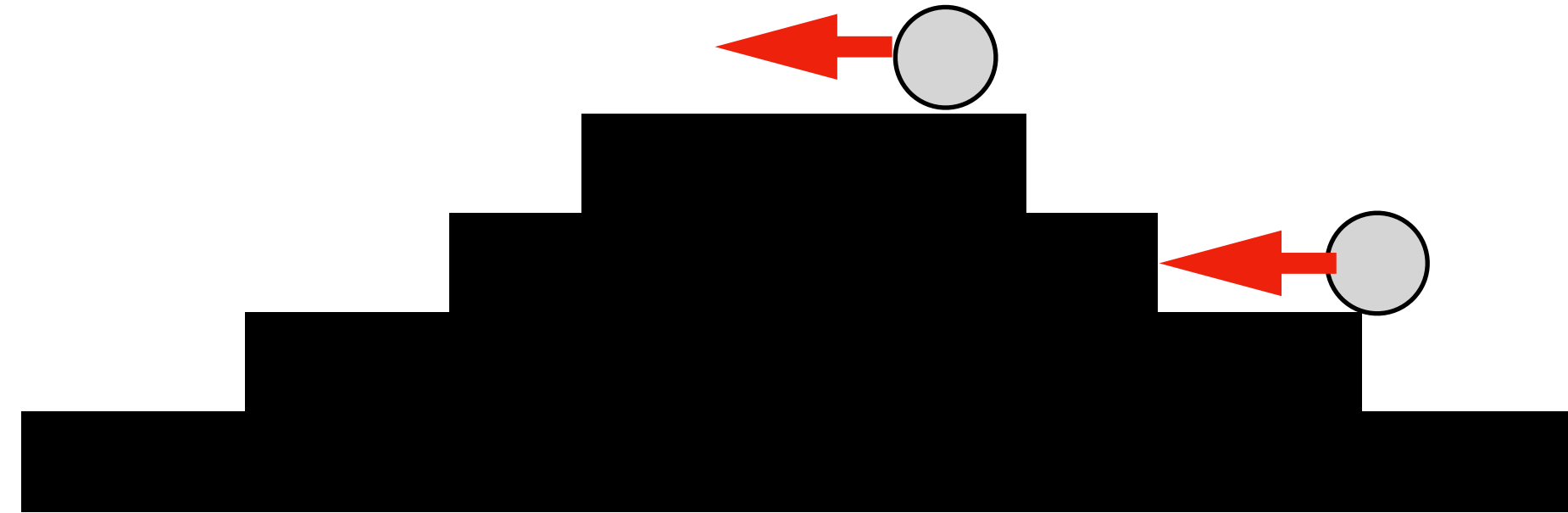
Morphological instability of a terrace edge during step-flow growth

Diffusional instability ("Schwoebel gives meanders")



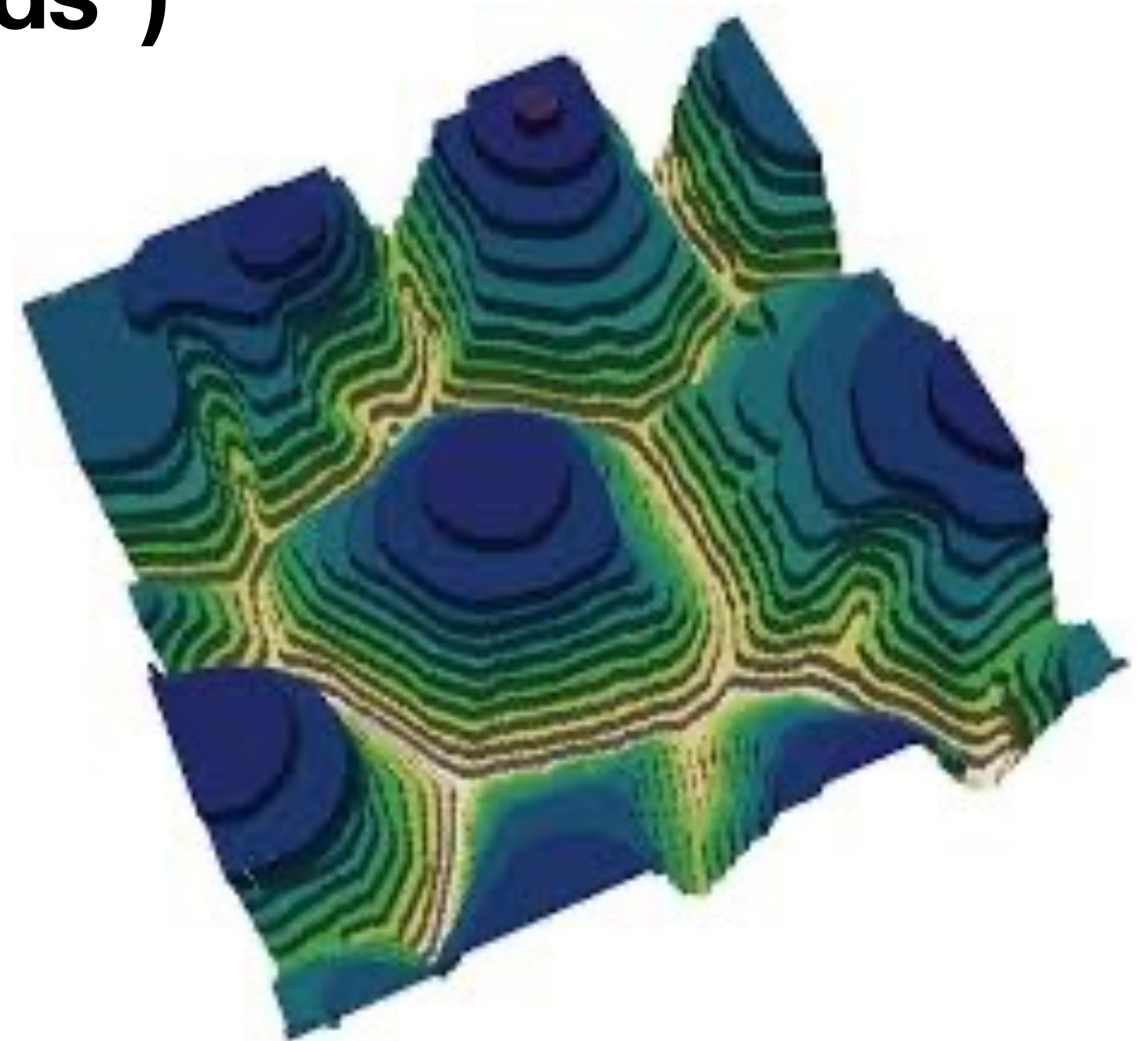
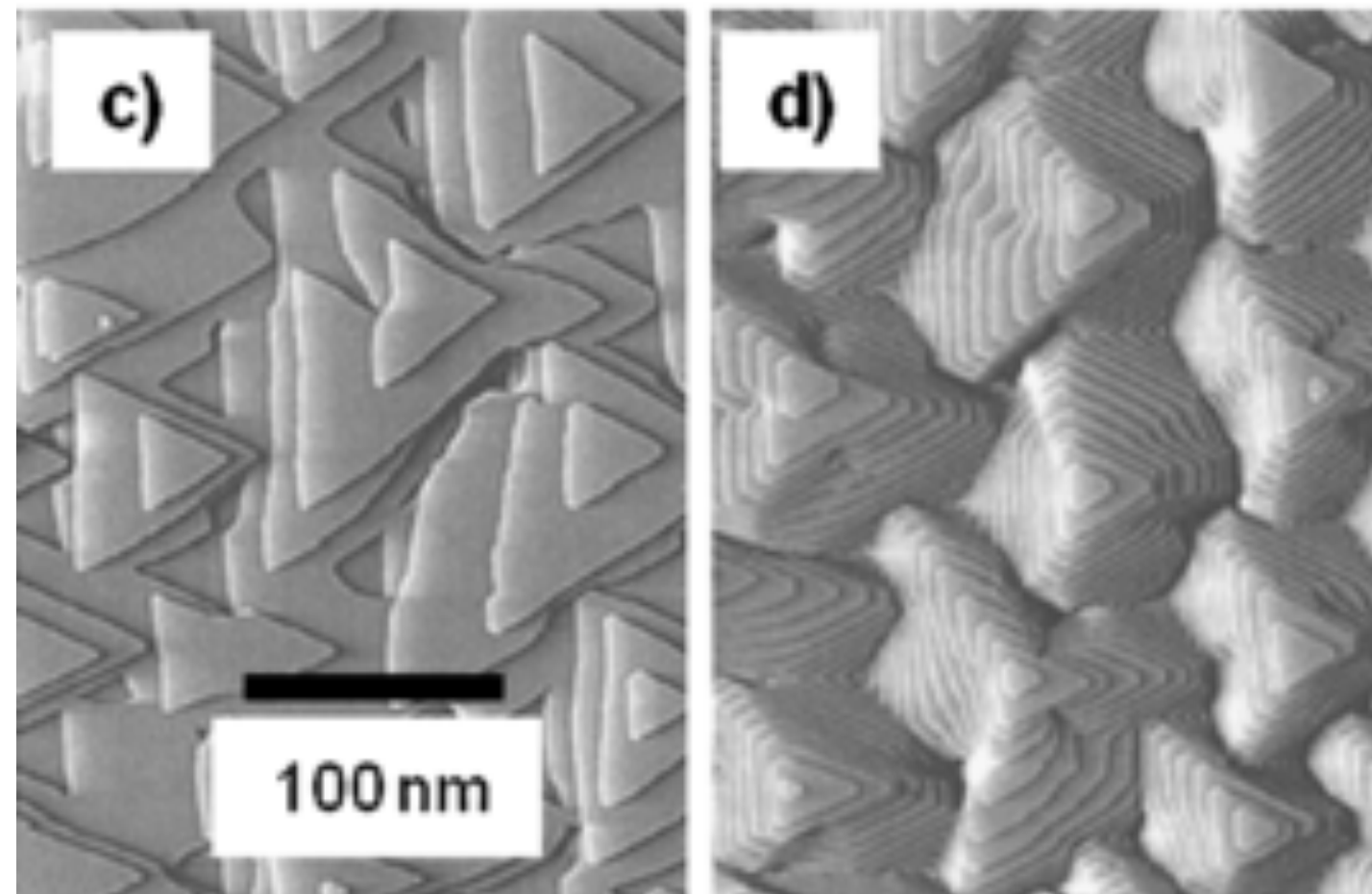
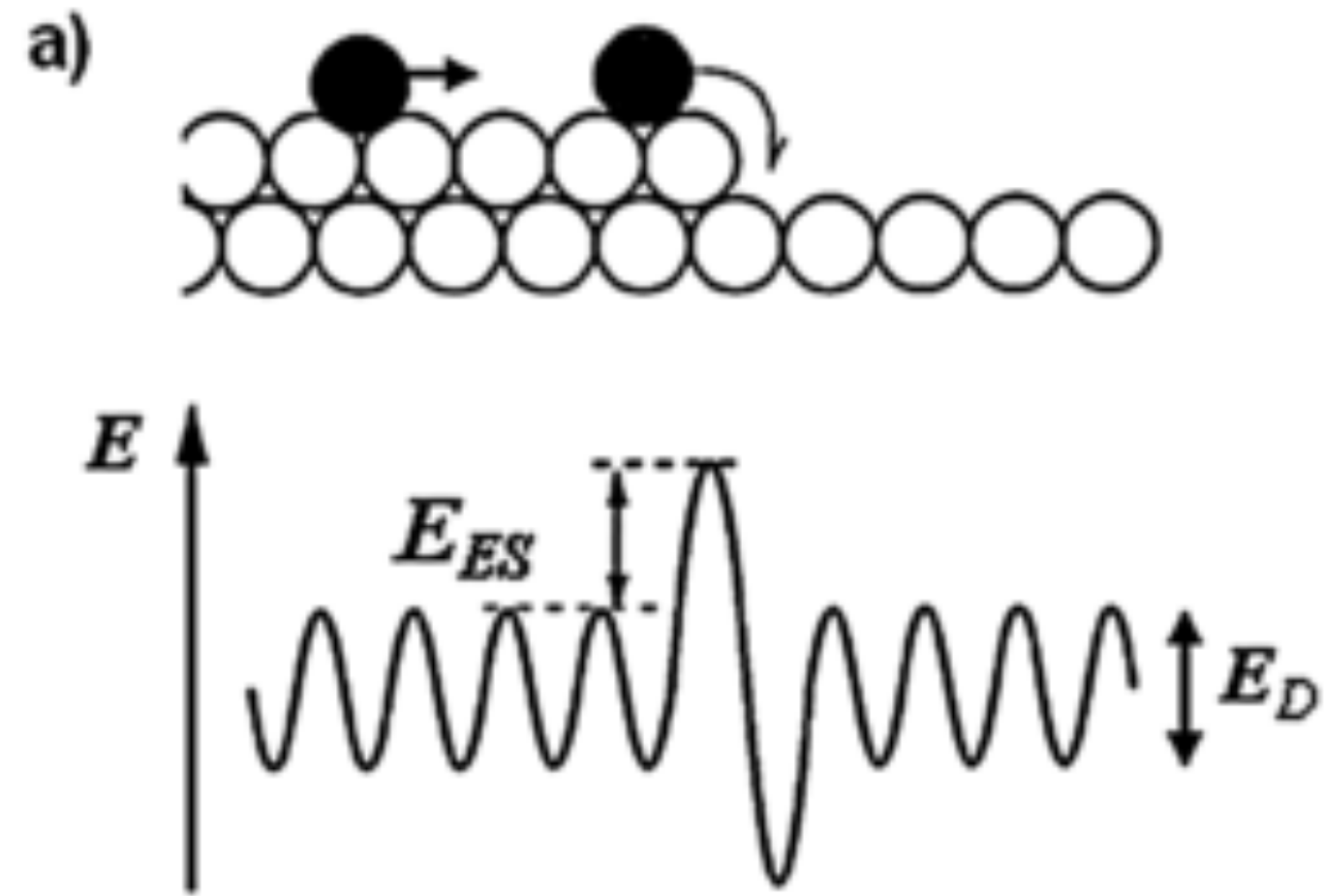
Cu (0,2,24)

2D to 3D: interlayer transport barriers ("Schwoebel gives mounds")

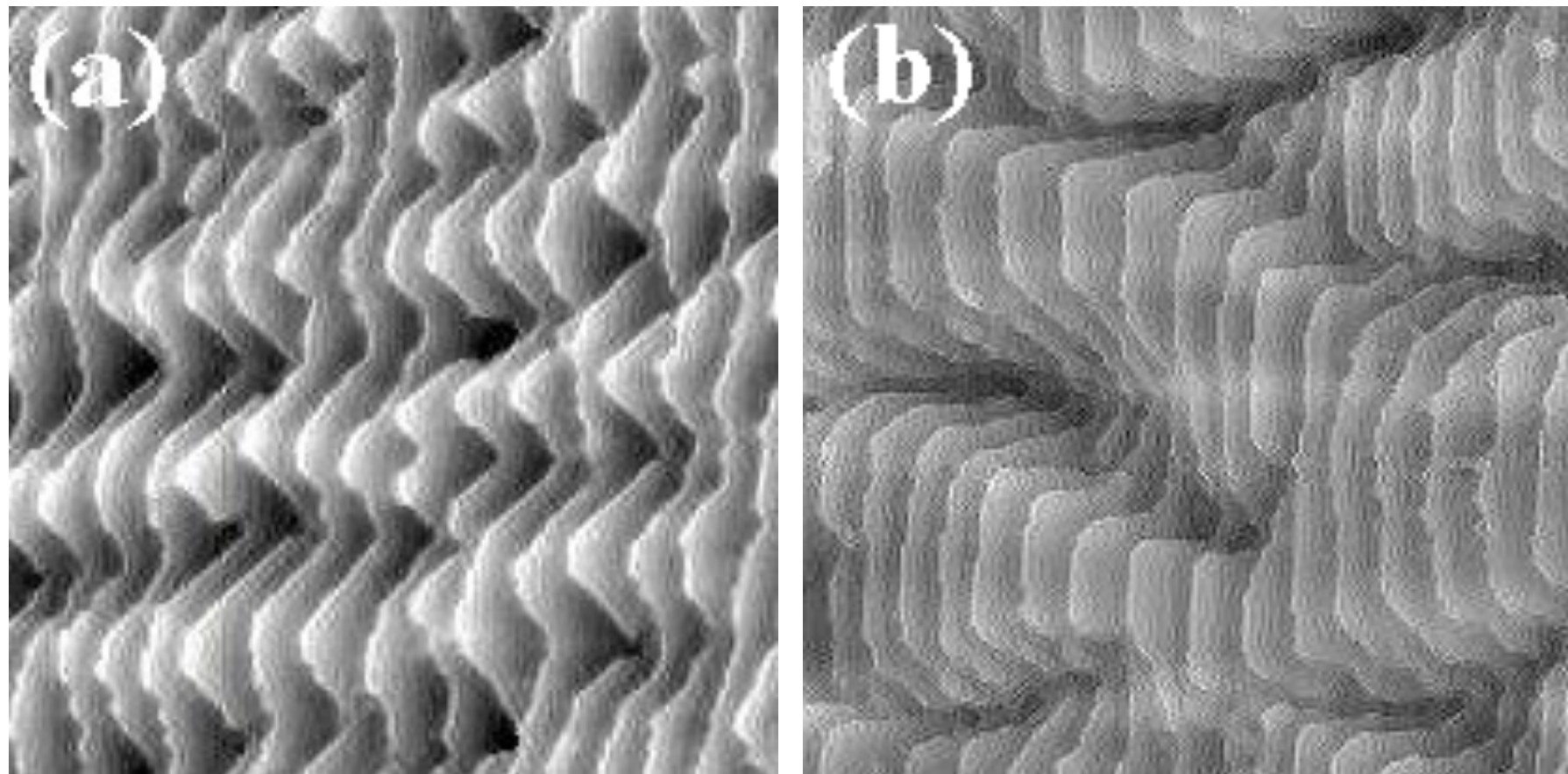


$$J_{\text{down}} < J_{\text{up}}$$

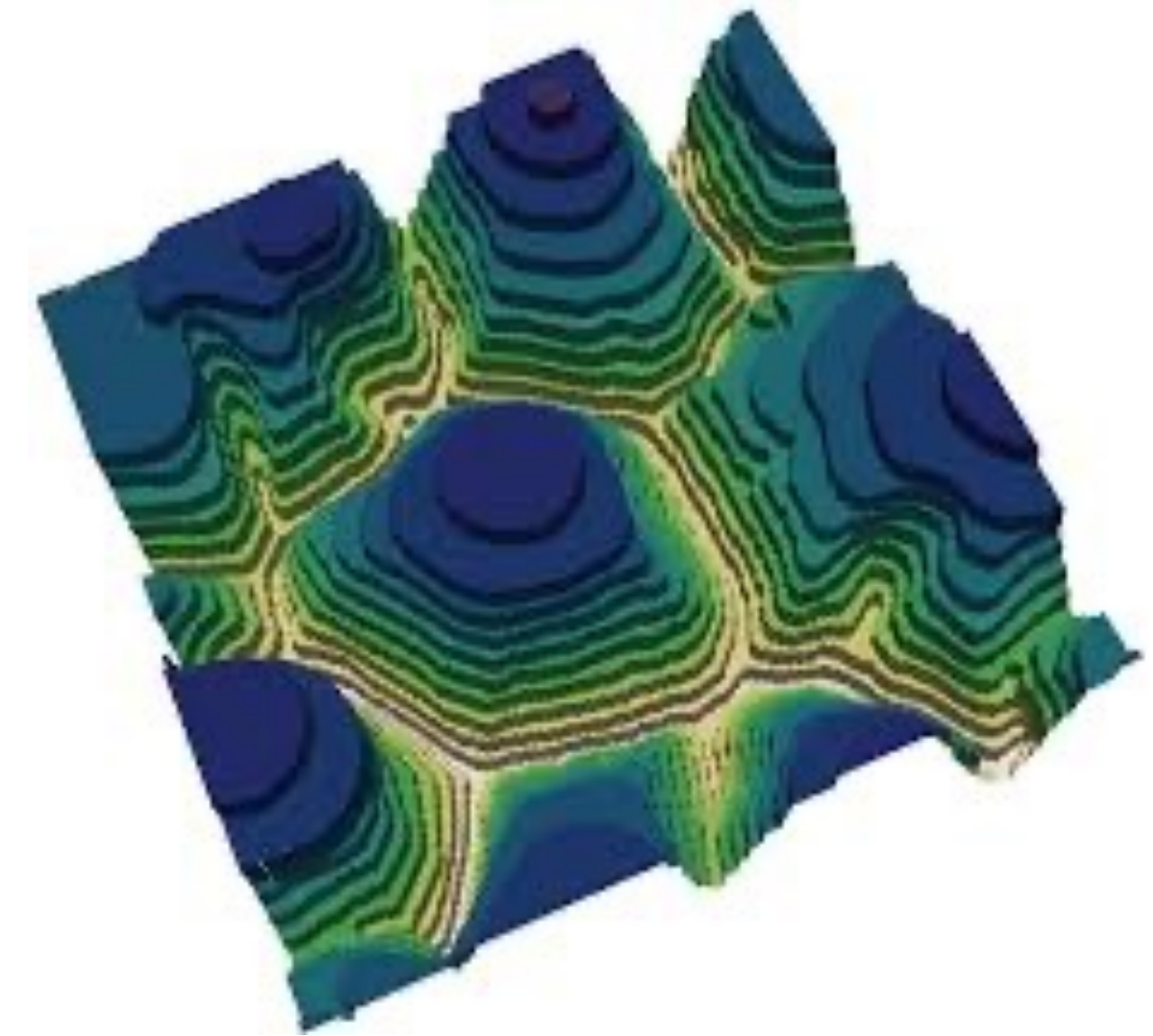
2D to 3D: interlayer transport barriers ("Schwoebel gives mounds")



2D to 3D: interlayer transport barriers ("Schwoebel gives mounds")



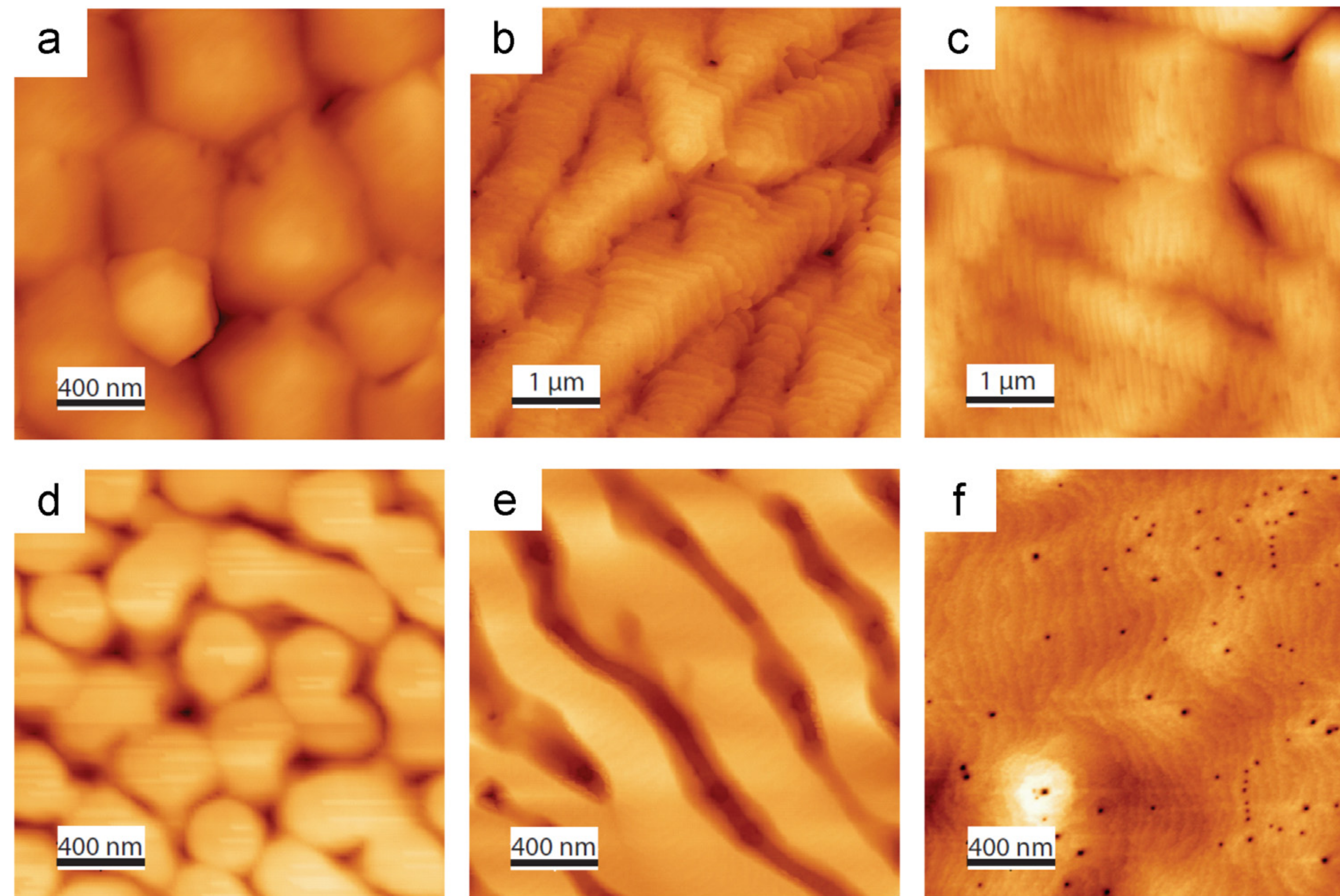
$$J_{\text{down}} < J_{\text{up}}$$



Net current “up”: leads to meandering steps and/or mounds

Unstable morphologies

Nils.A.K. Kaufmann et al. / Journal of Crystal Growth 433 (2016) 36–42



AFM images of NH₃-MBE-grown GaN layers in hillock, step-meandering, and step-flow regime. For (a), (b) and (c) the growth temperatures were 800°C, 865°C and 920°C, respectively. The sample displayed in (d) and (e) was grown by nitrogen-rich PAMBE at 790°C on FS GaN (unintentionally offcut variations of the GaN FS substrate are responsible for the different morphologies). (f) Step-flow growth surface obtained by PAMBE at 740°C under metal-rich conditions.

GaN vicinals — meandering steps or mounds

Unstable morphologies

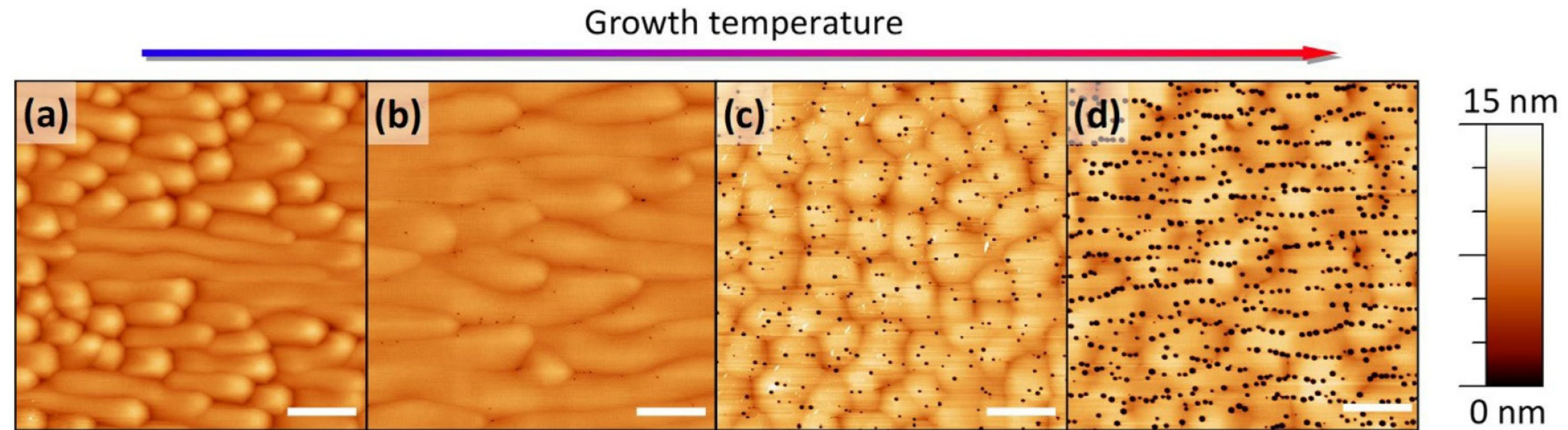


FIG. 1. $10 \times 10 \mu\text{m}^2$ atomic force microscopy images of the AlGaN surface of samples A–D [(a)–(d)]. The growth temperature increases from 743 °C (left) to 818 °C (right). The scale bar corresponds to 2 μm.

GaN vicinals — meandering steps + small mounds or mounds

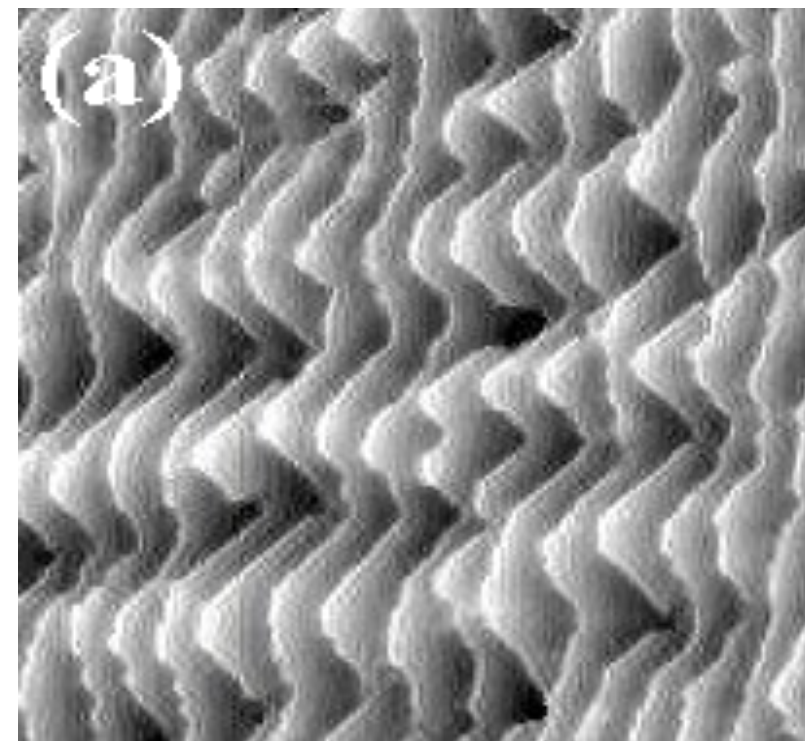
Morphological instabilities in epitaxial growth: a case study

PHYSICAL REVIEW B, VOLUME 64, 165401

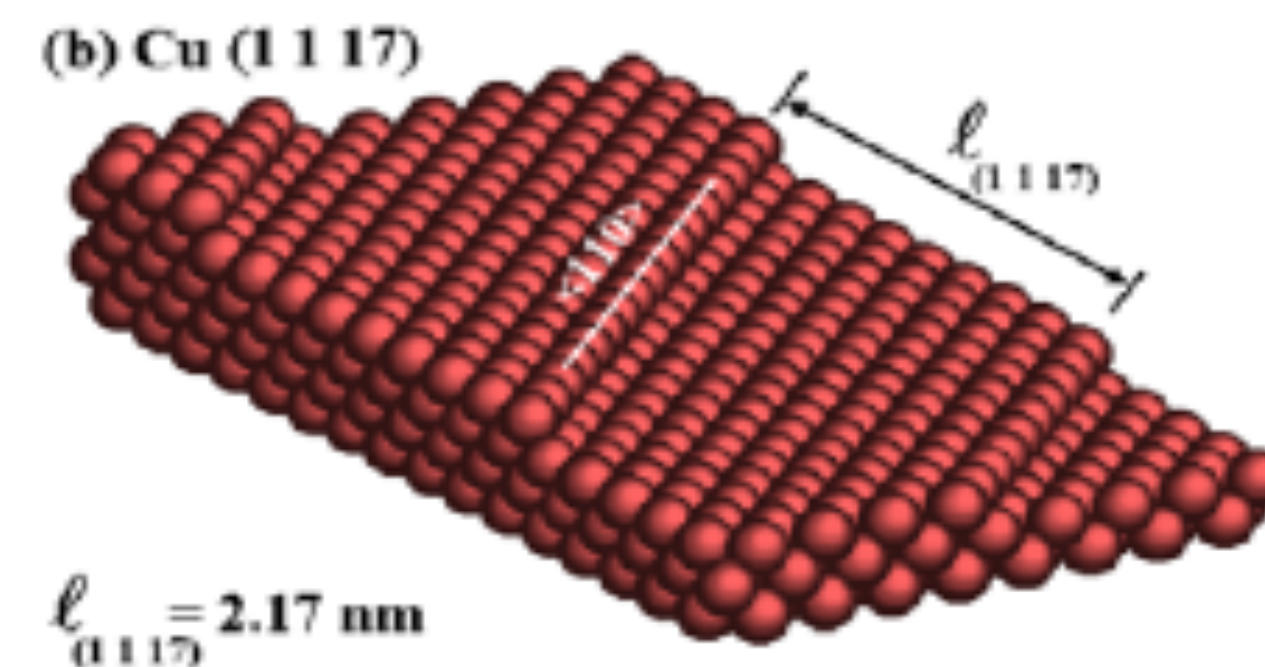
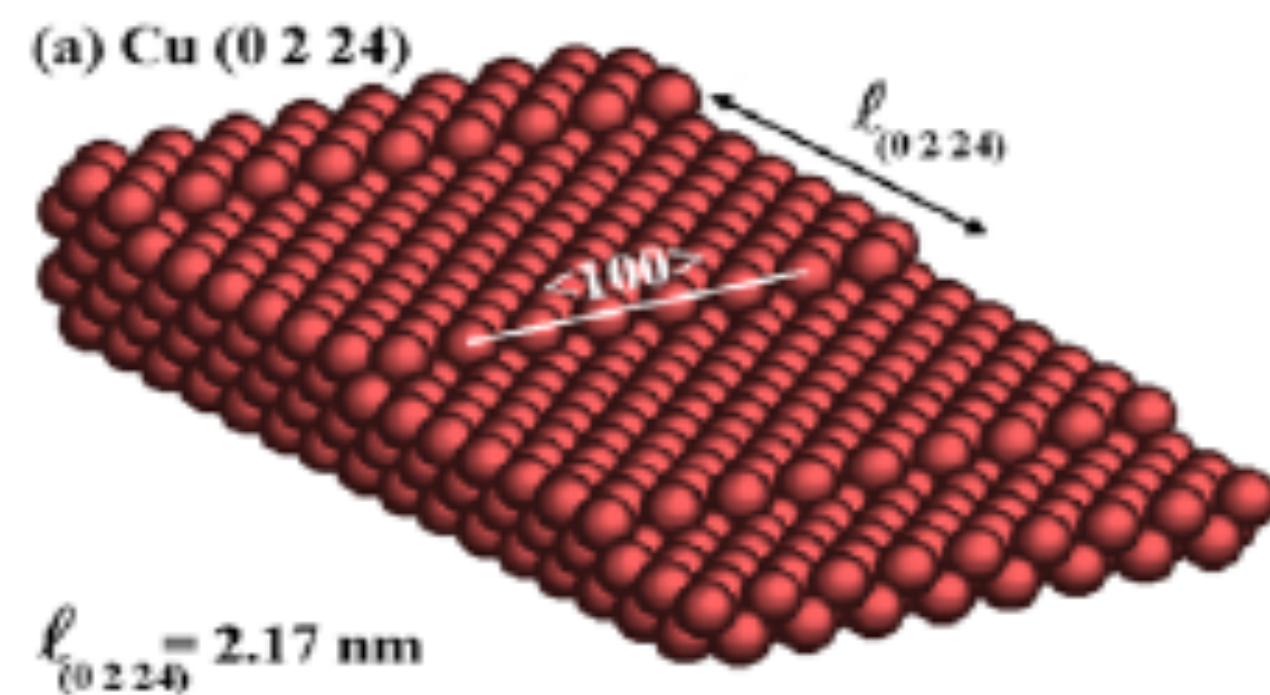
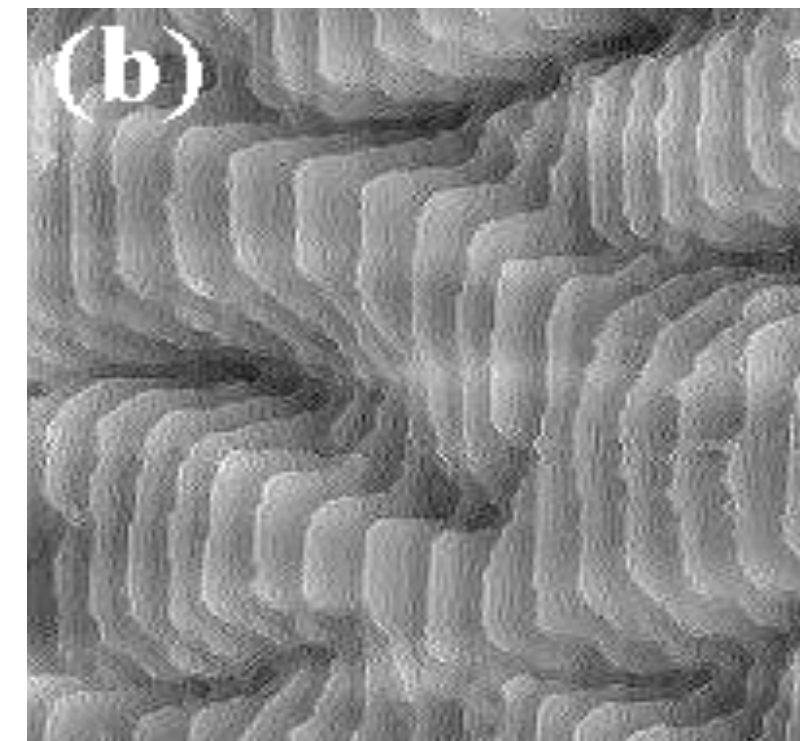
Morphological instability of Cu vicinal surfaces during step-flow growth

T. Maroutian, L. Douillard, and H.-J. Ernst
CEA Saclay, DSM/DRECAM/SPCSI, F-91191 Gif sur Yvette, France
(Received 13 February 2001; published 28 September 2001)

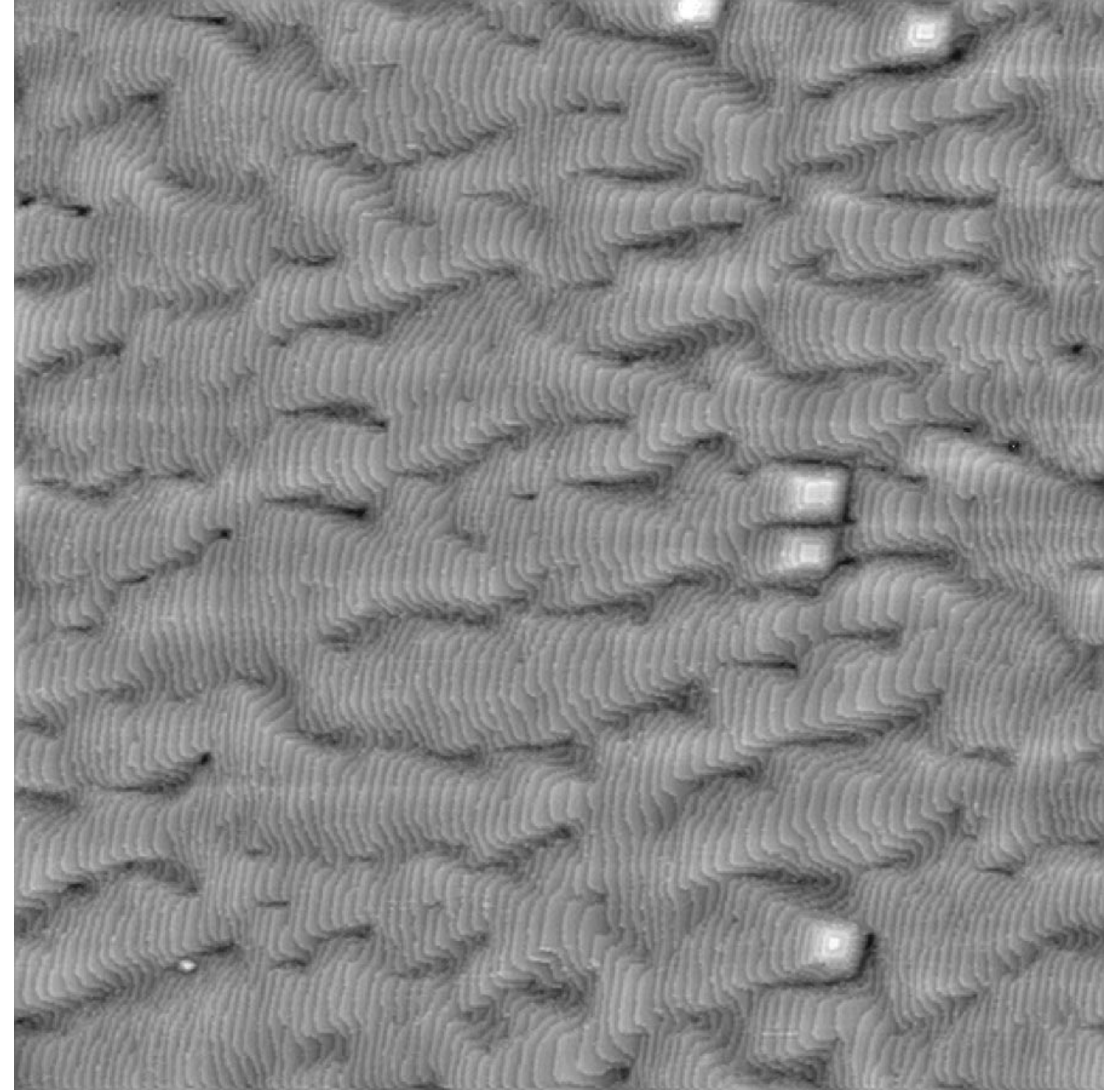
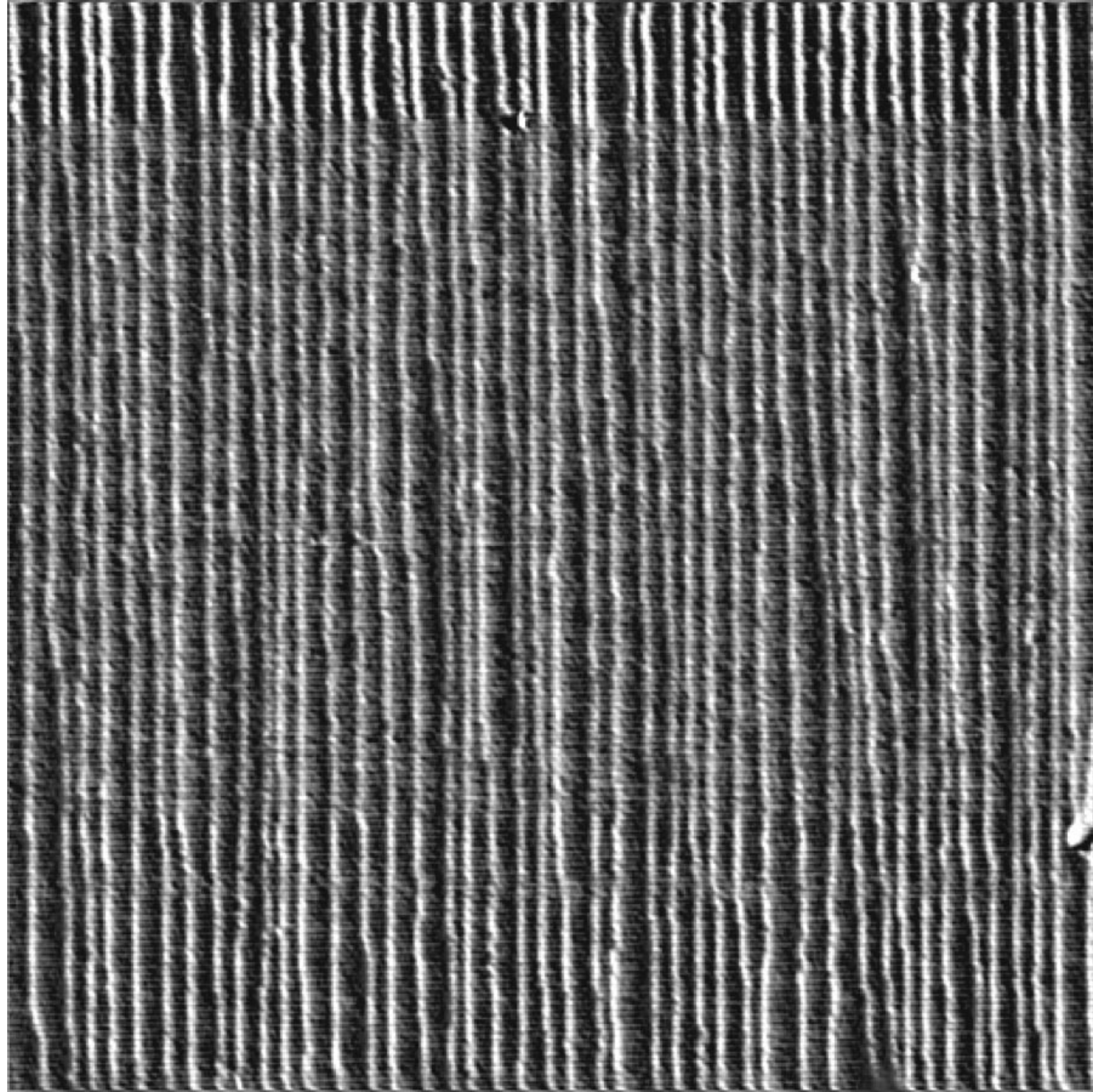
Cu (0,2,24)

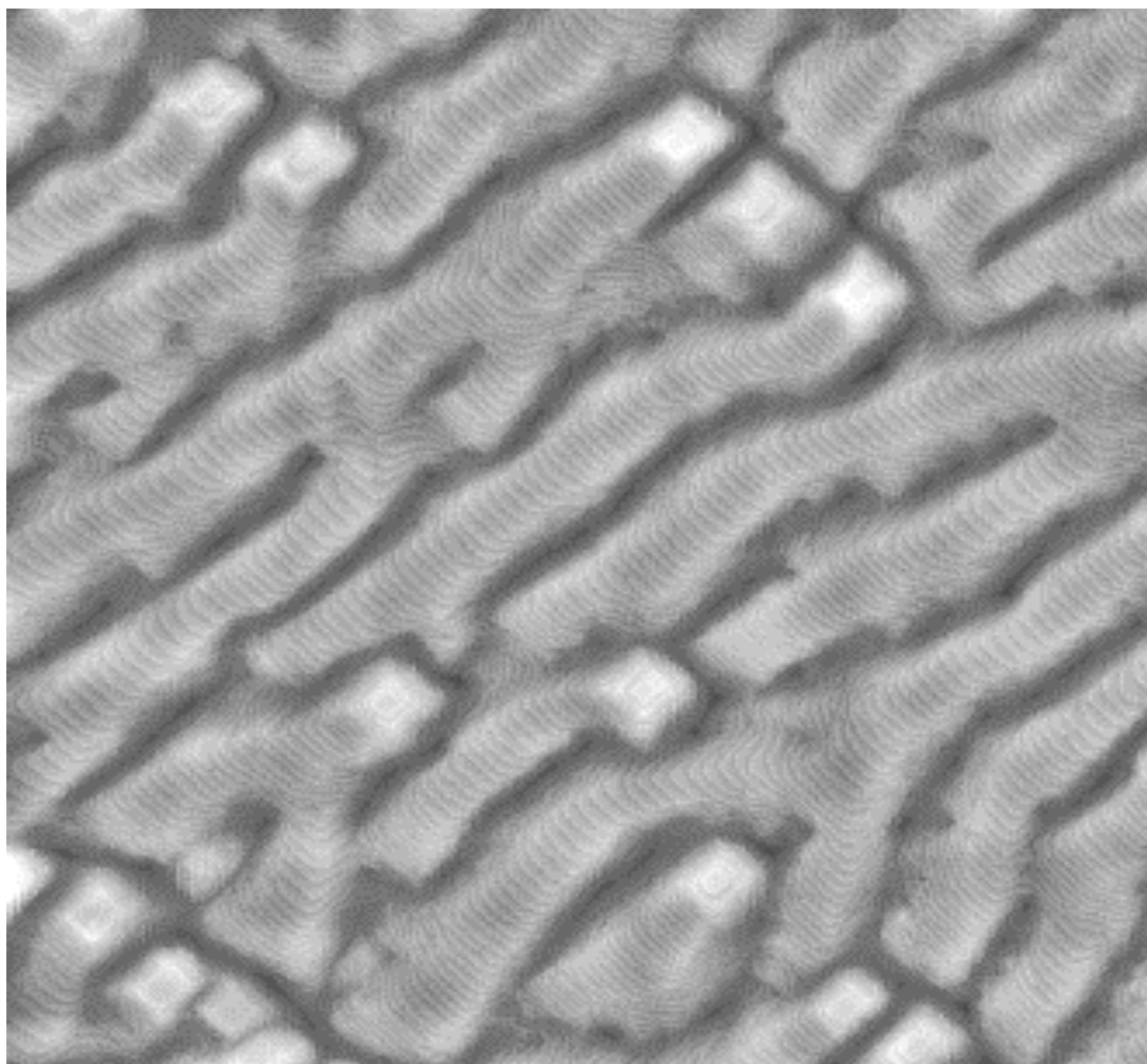


Cu (1,1,17)

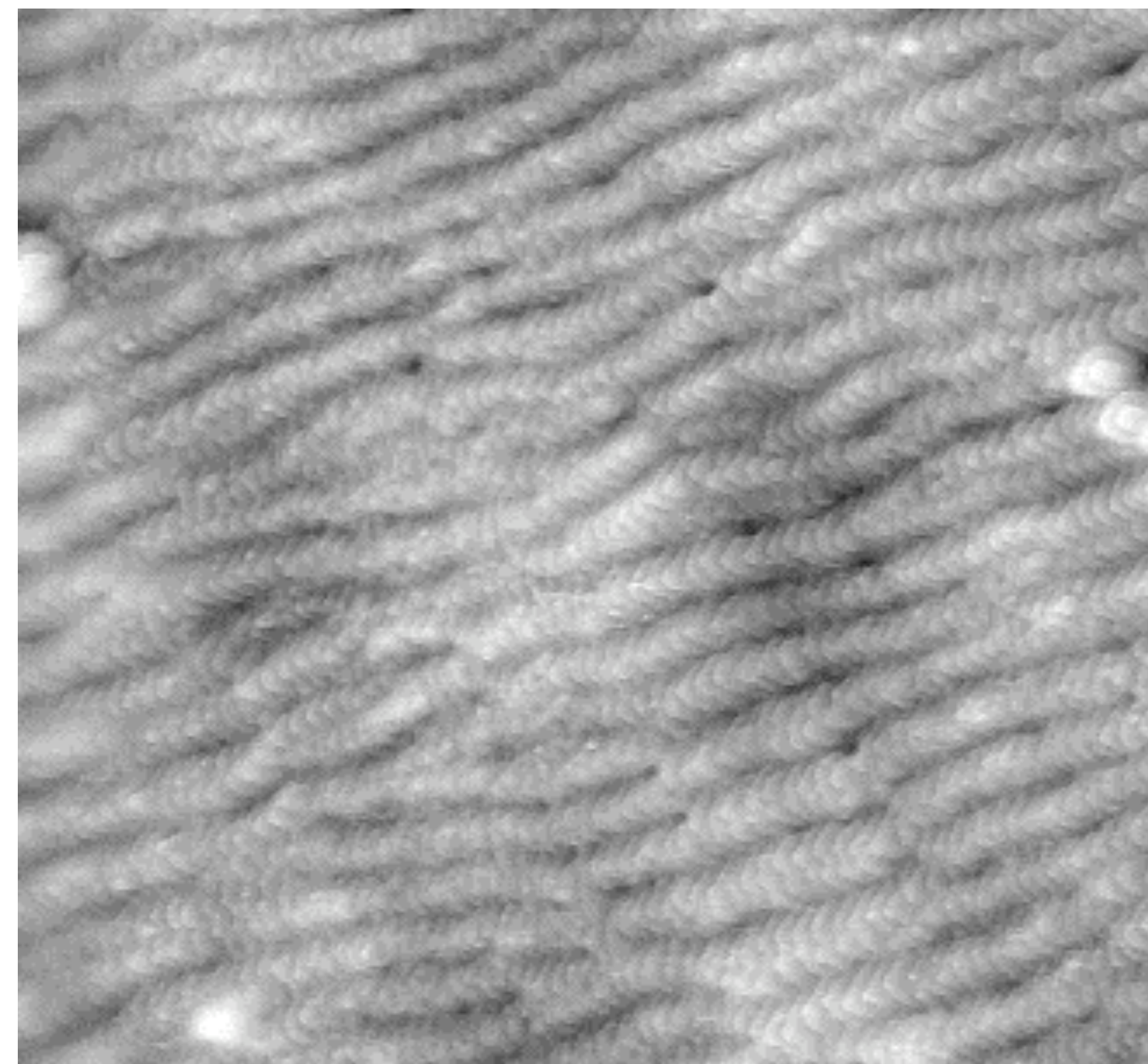


Cu (1,1,17)

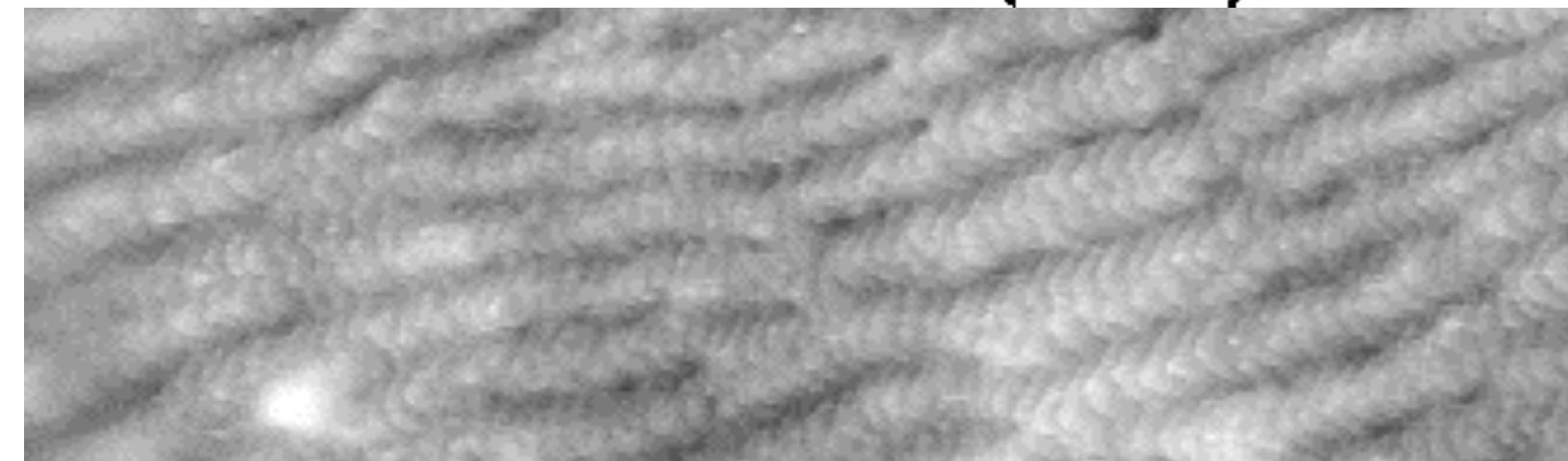
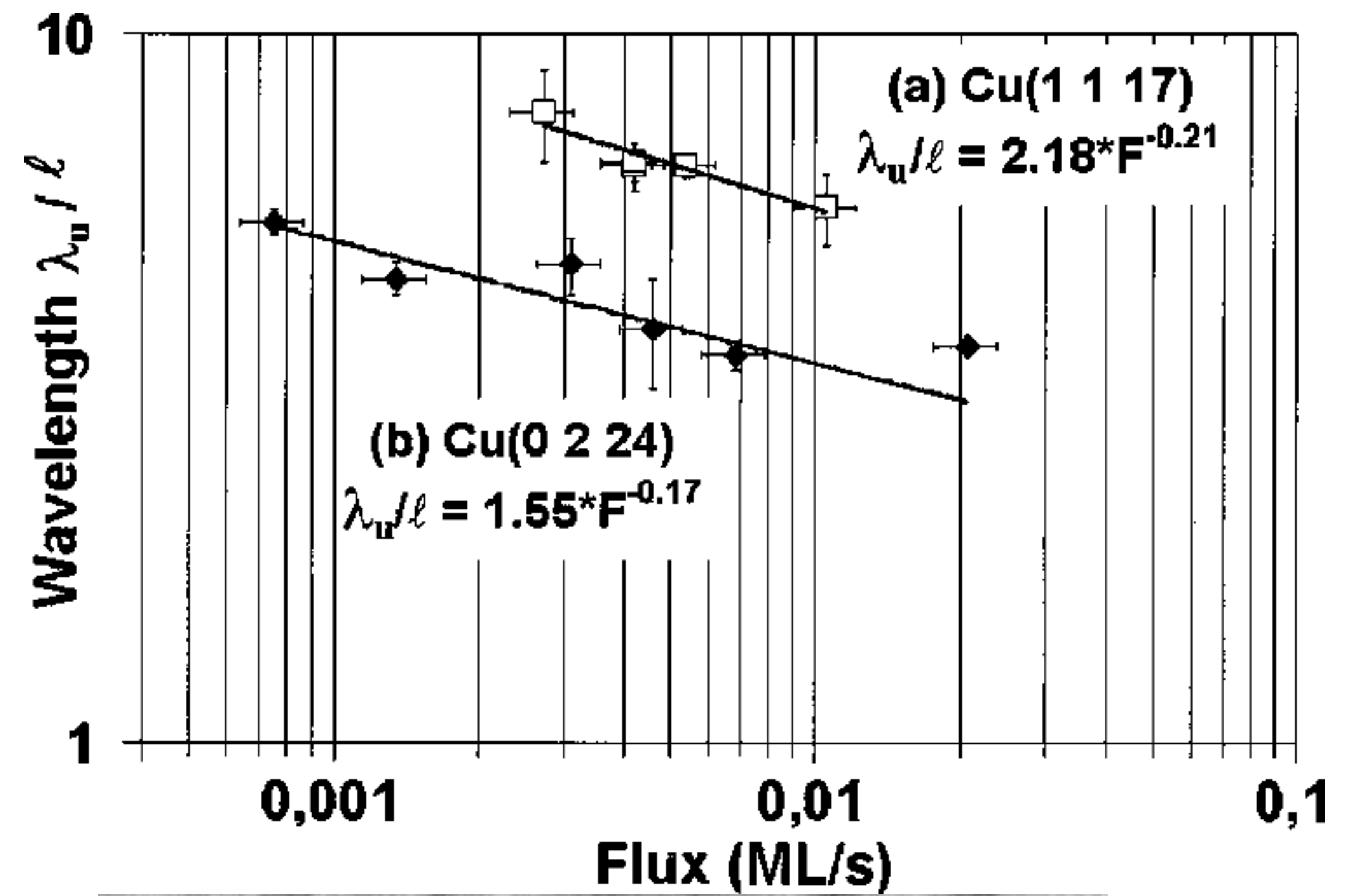
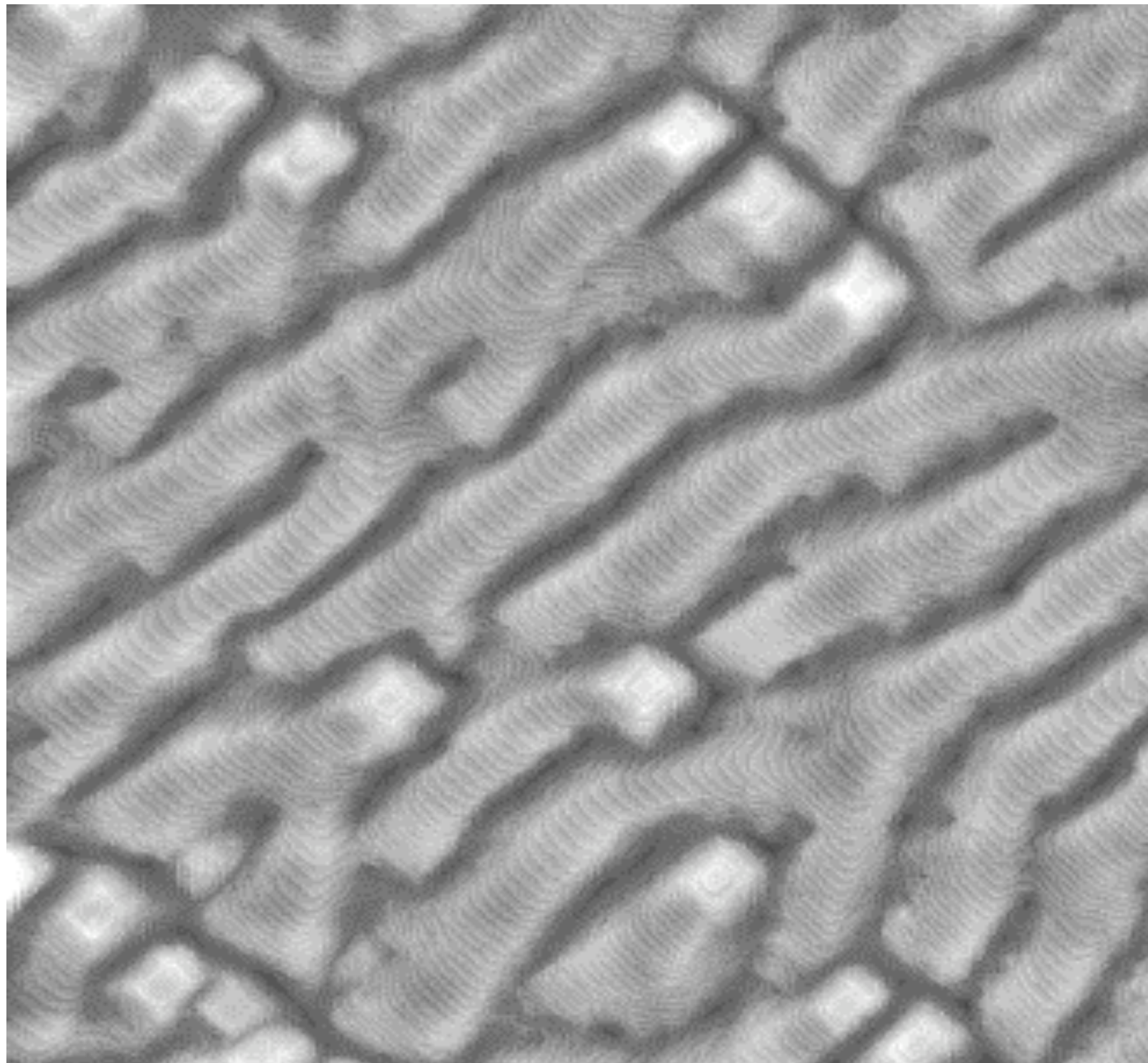




Cu (1,1,17)



Cu (0,2,24)

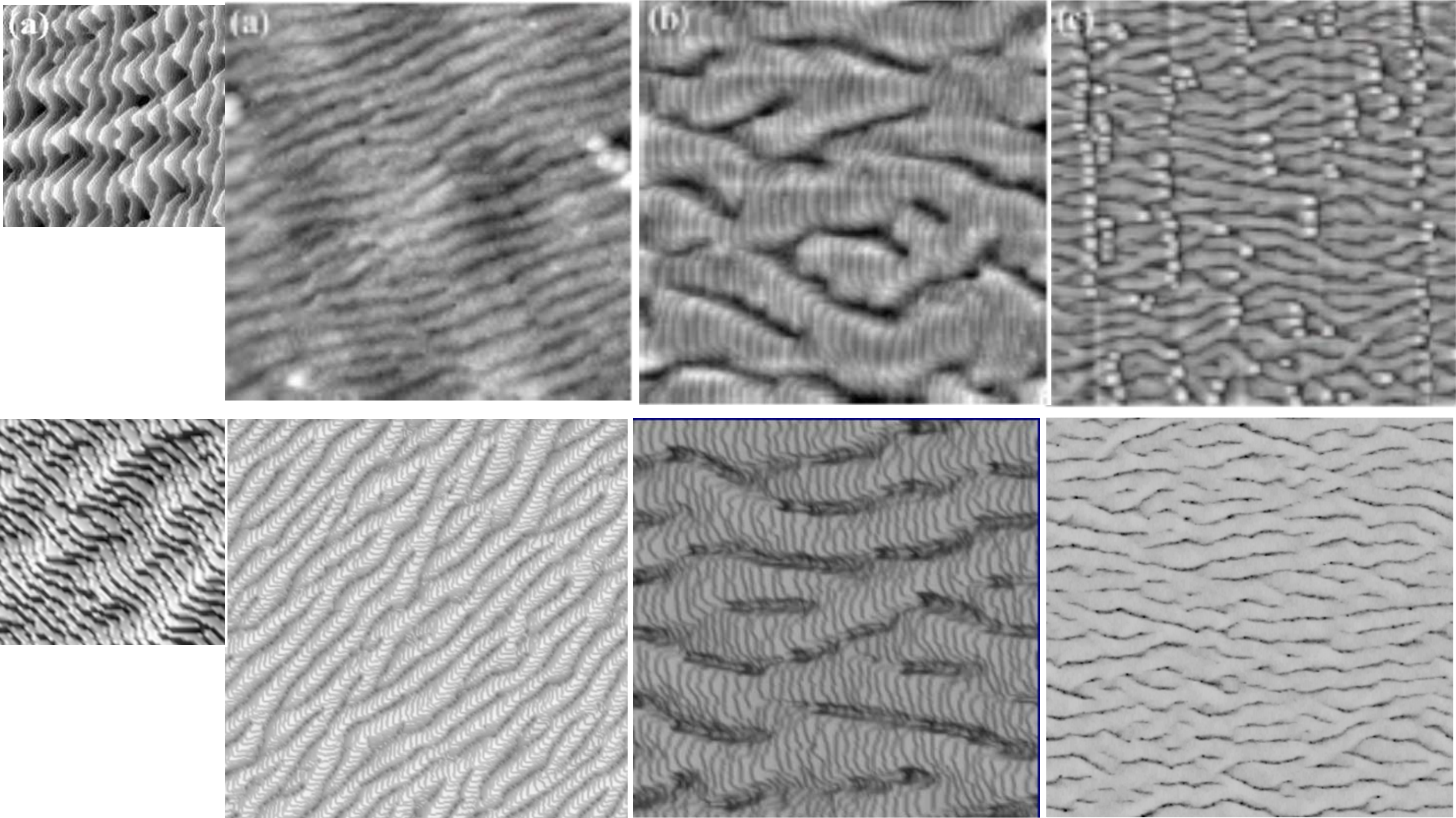


Cu vicinals — unstable growth:
meandering steps (?) + pyramids (?)

Cu (0 2 24)

Cu (1 1 17)

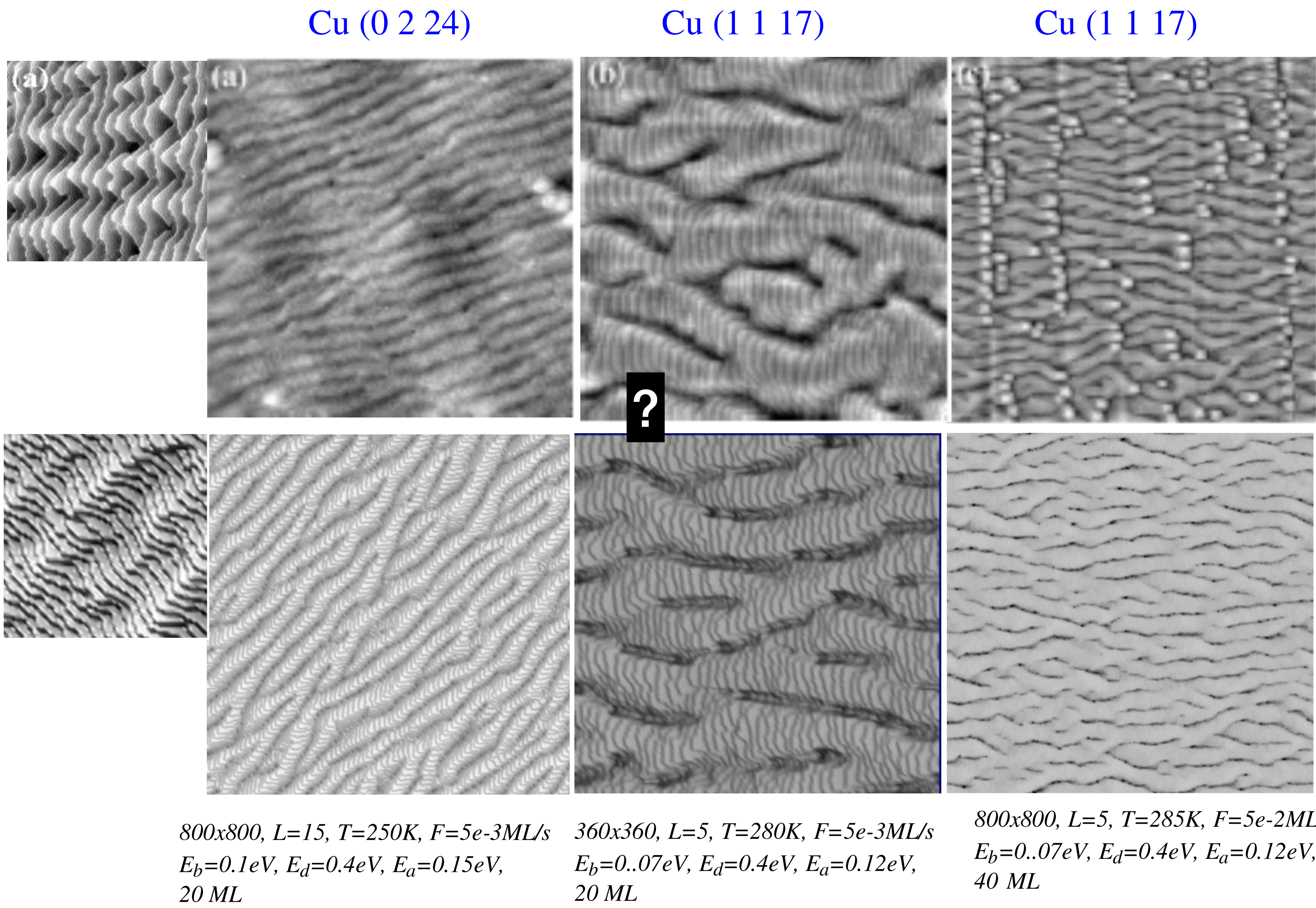
Cu (1 1 17)

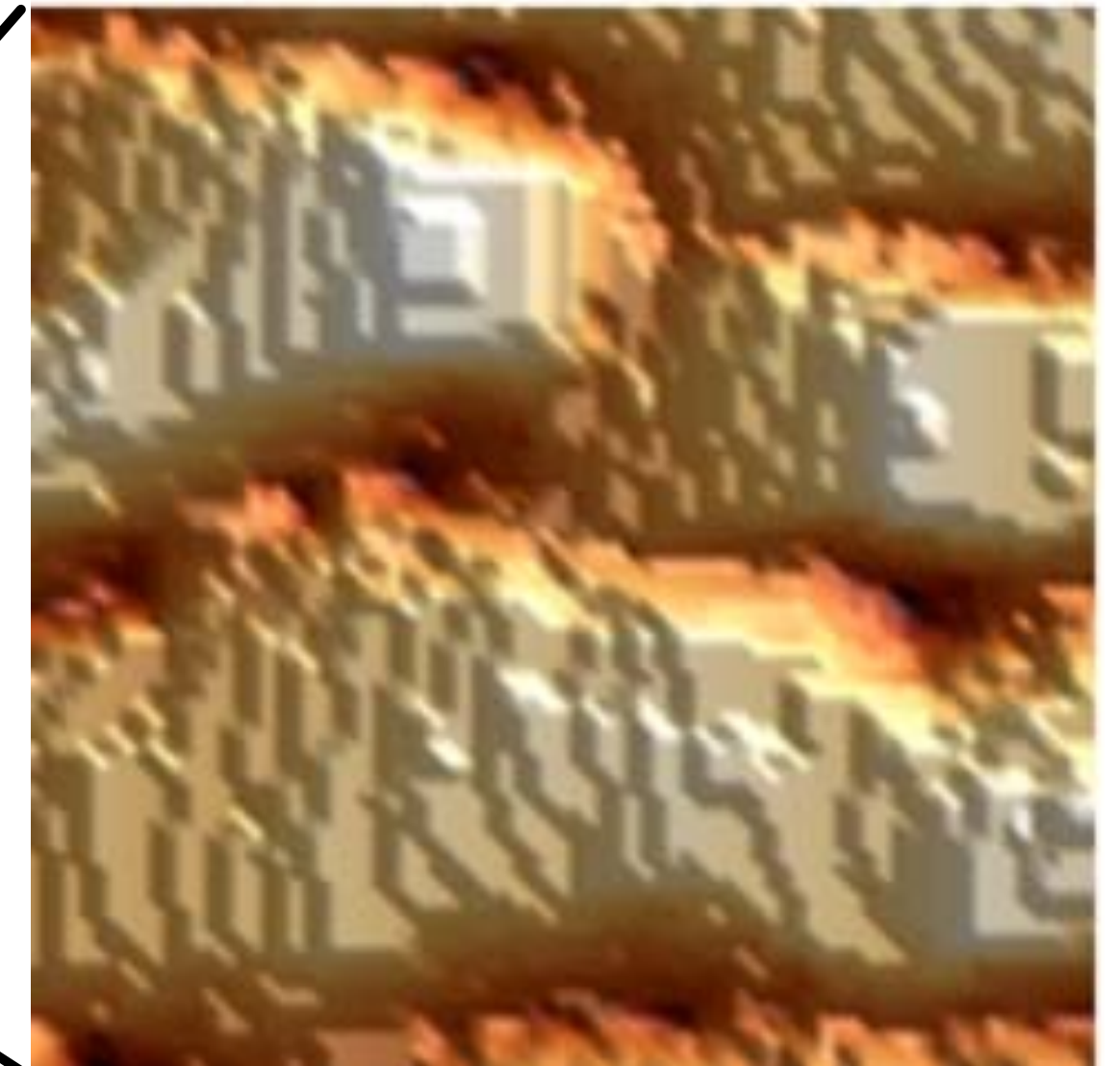
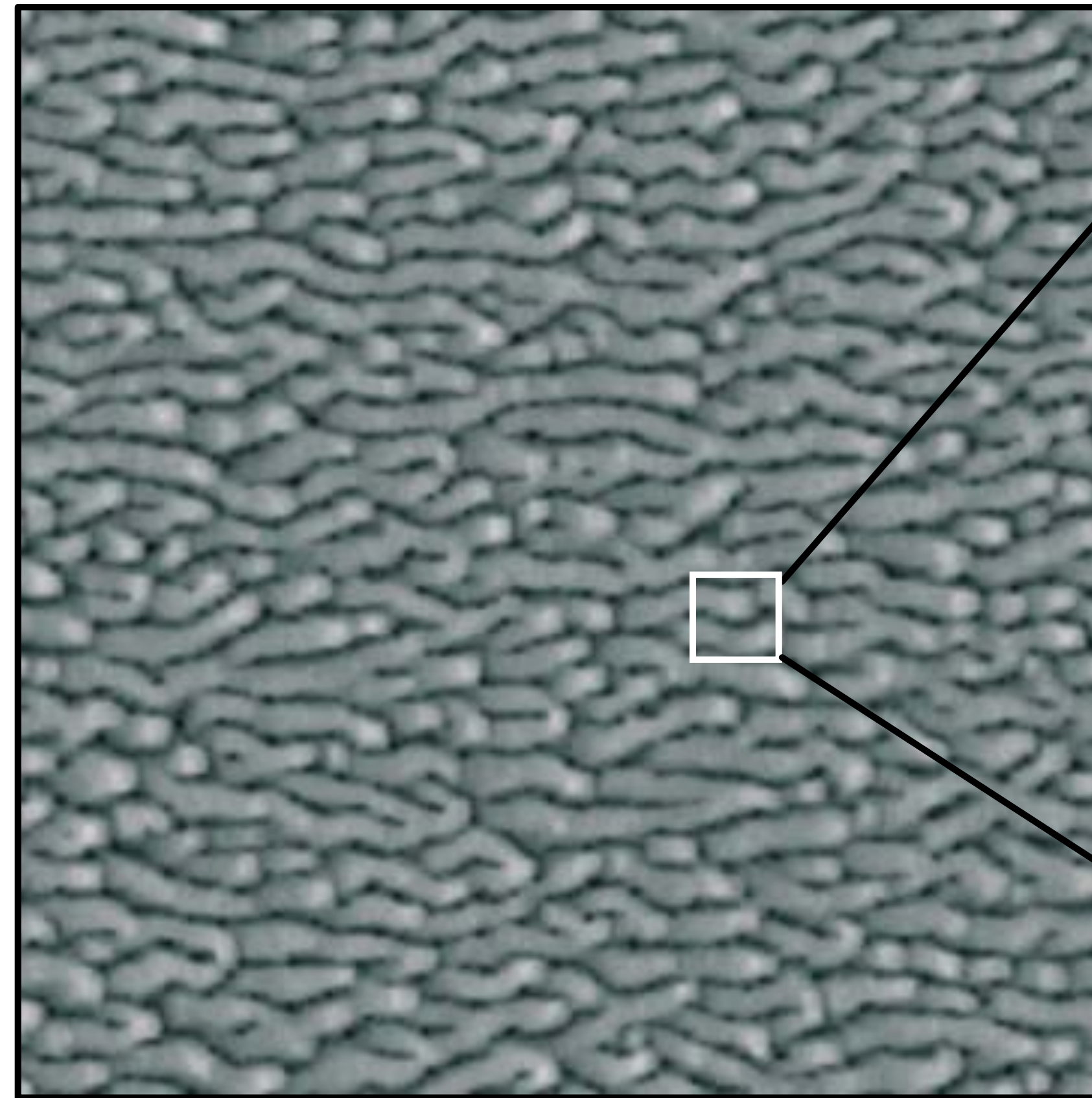
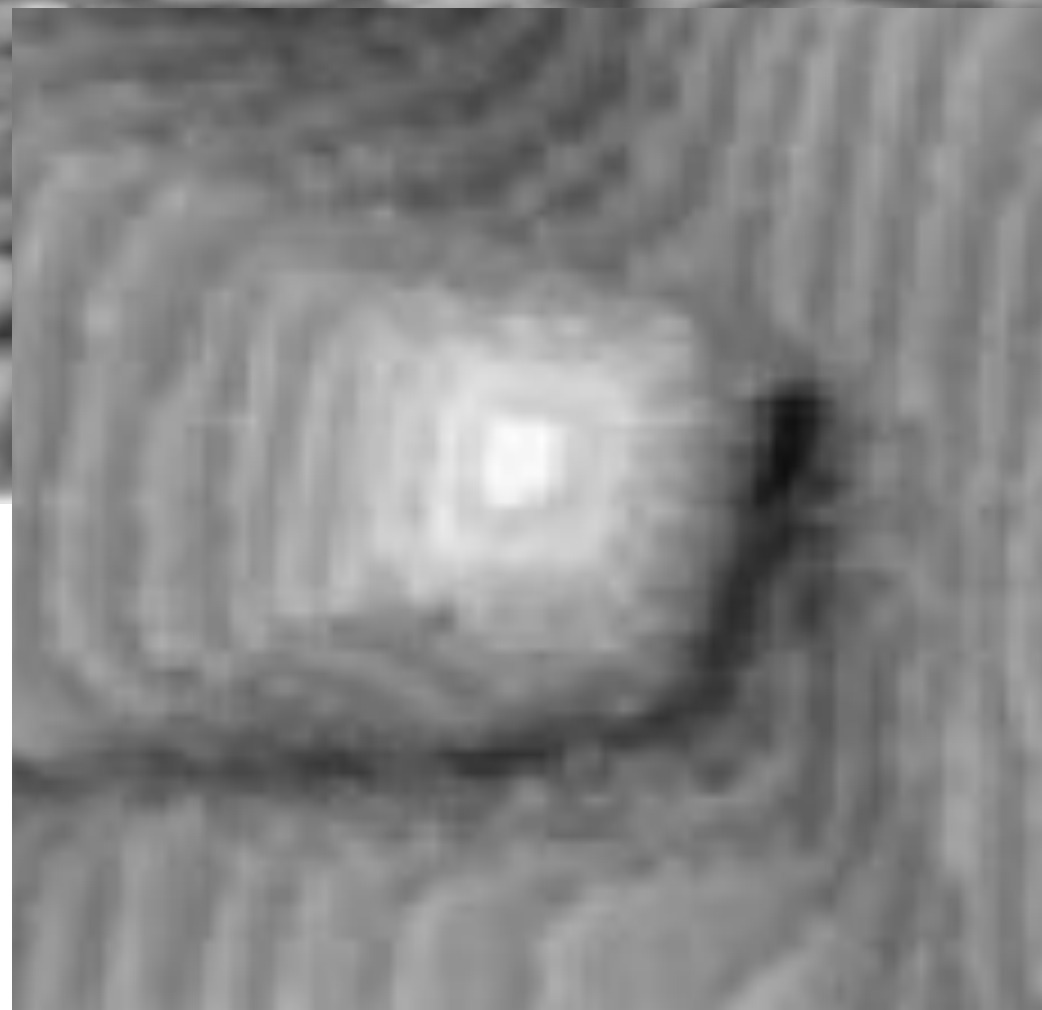
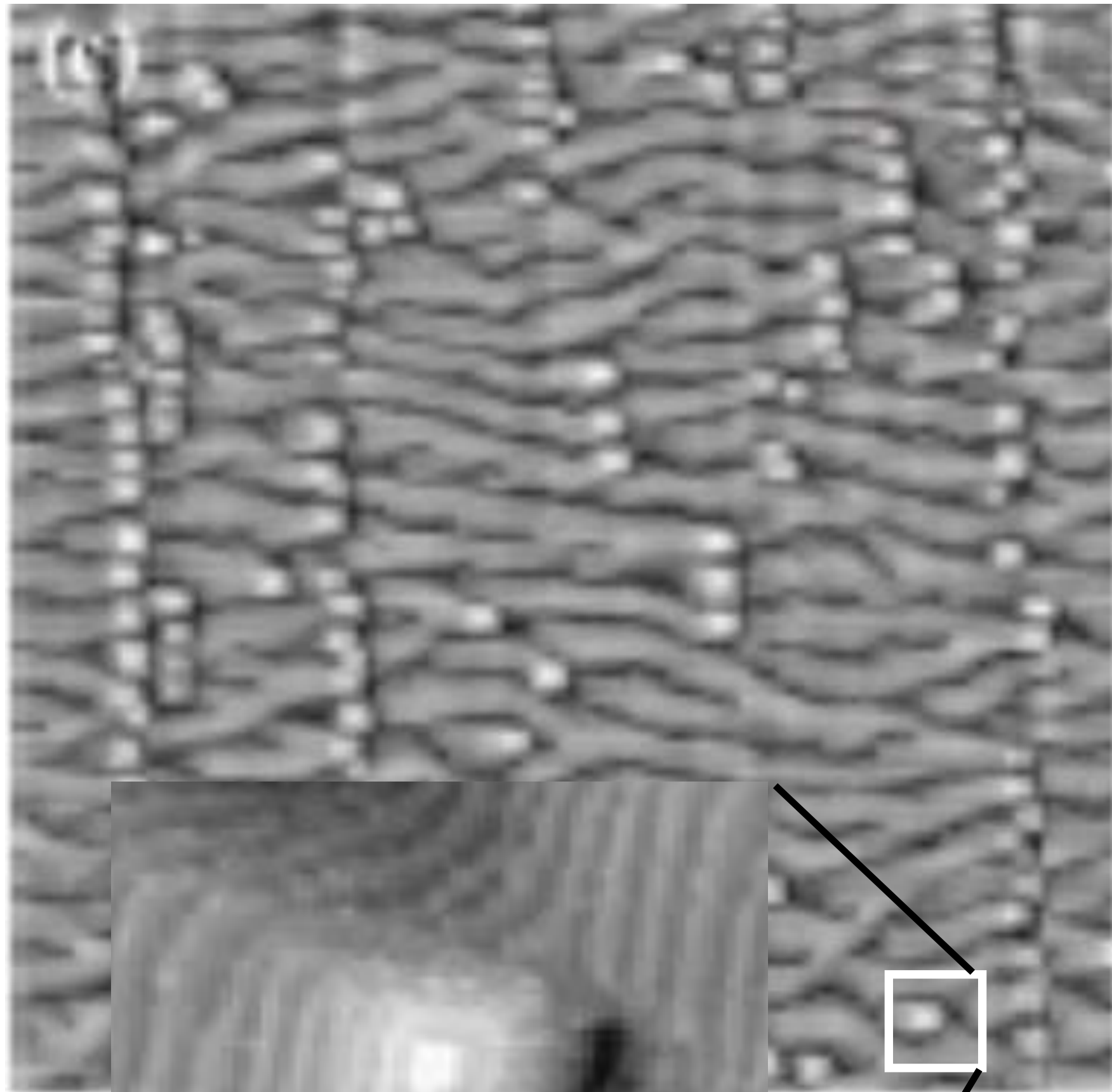


800x800, $L=15$, $T=250K$, $F=5e-3ML/s$
 $E_b=0.1eV$, $E_d=0.4eV$, $E_a=0.15eV$,
20 ML

360x360, $L=5$, $T=280K$, $F=5e-3ML/s$
 $E_b=0.07eV$, $E_d=0.4eV$, $E_a=0.12eV$,
20 ML

800x800, $L=5$, $T=285K$, $F=5e-2ML/s$,
 $E_b=0.07eV$, $E_d=0.4eV$, $E_a=0.12eV$,
40 ML





Add an impurity

DFT points to W

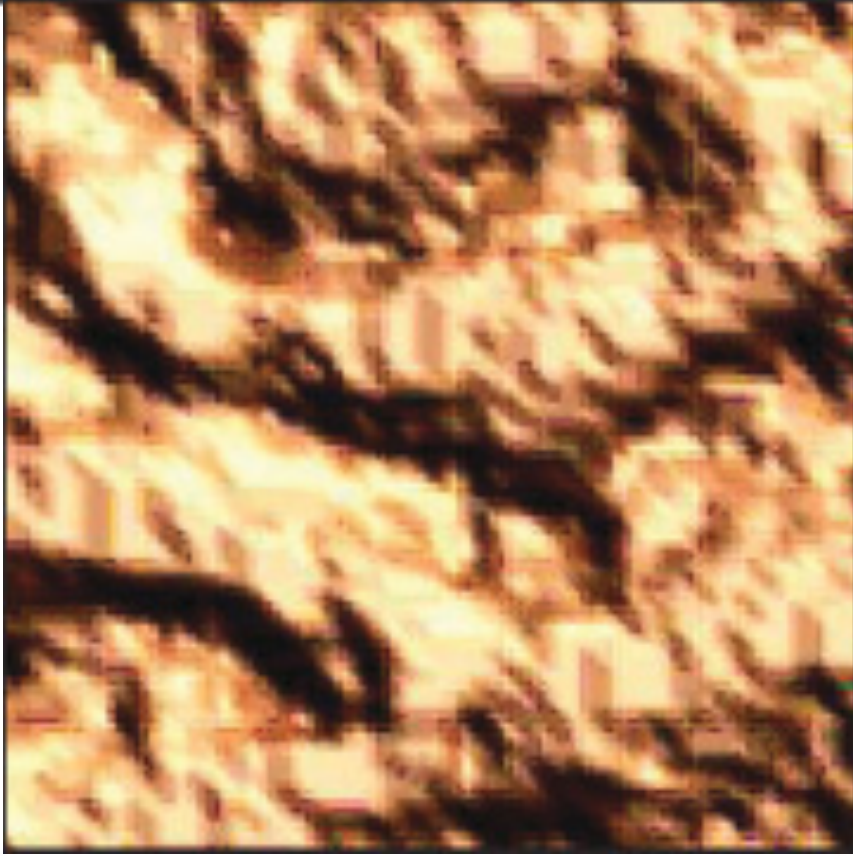
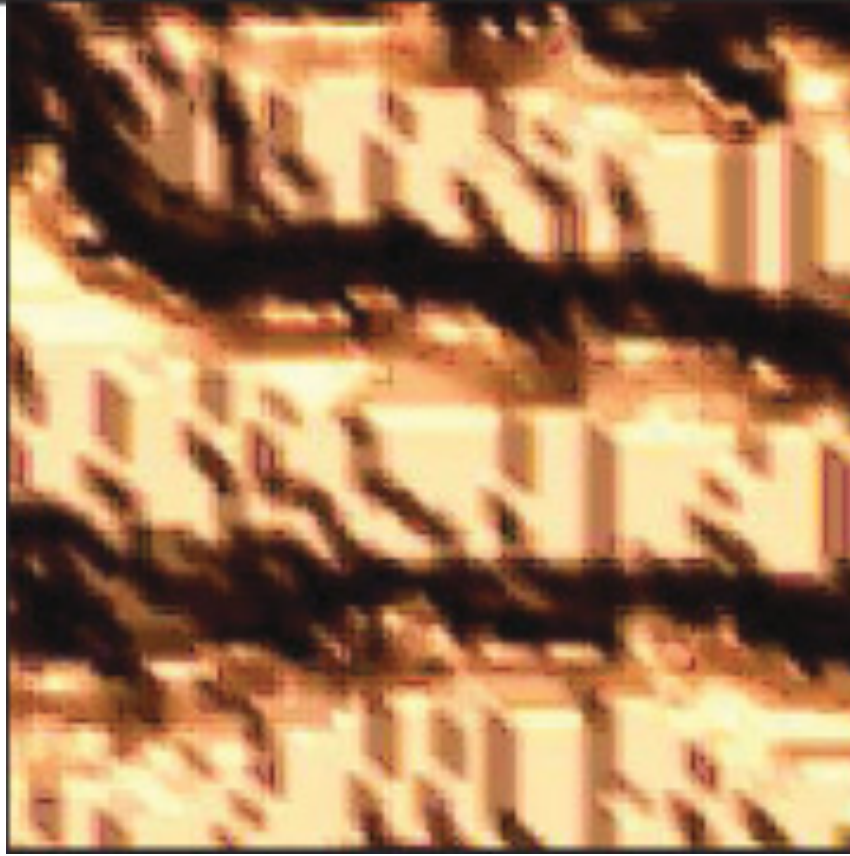
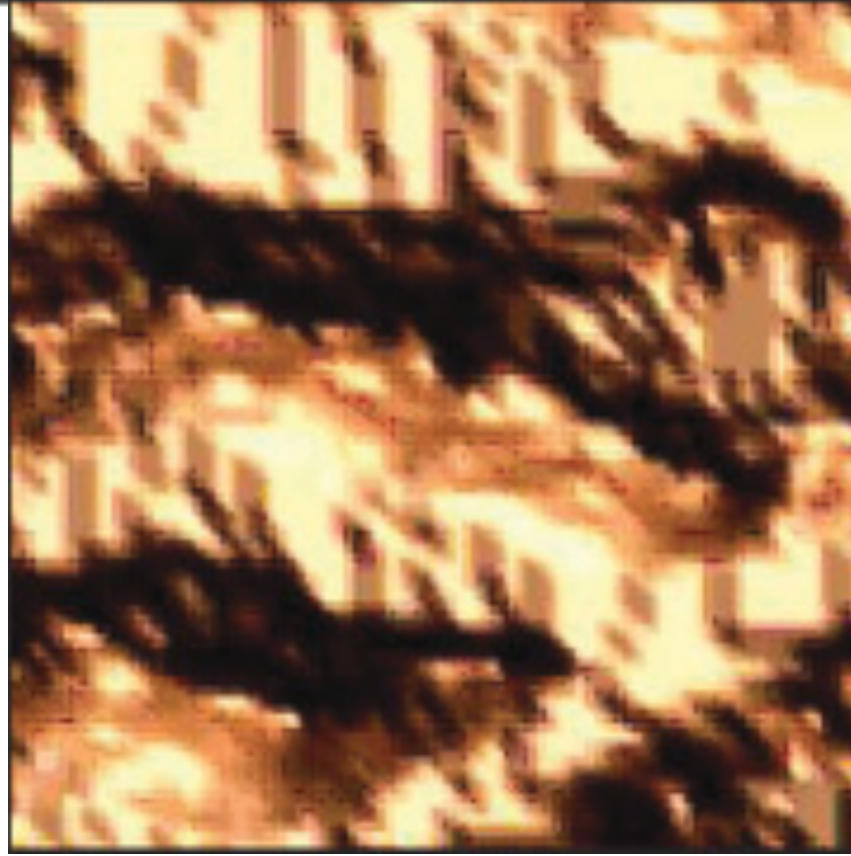
Imp.	Cu	C (Set 1)	Al (Set 2)	Ni (Set 3)	W (Set 4)
γ	0.45 ± 0.05	0.35 ± 0.06	0.50 ± 0.06	0.41 ± 0.05	0.17 ± 0.02
Zoom View					

FIG. 2. (Color online) Rough estimate of the exponent γ ($\lambda_m \sim F^{-\gamma}$) and of the possibility of pyramid formation for pure Cu and for Cu codeposited with a member of each of the four sets of impurities, the impurity being 2% of the flux. Simulations are done for five values of F : 0.005, 0.01, 0.05, 0.1, 0.5 ML/s.⁴¹ For pure Cu (Cu-Cu) γ is consistent with the Bales-Zangwill⁷ value ($\gamma = 0.5$) while for Cu-W ($\gamma = 0.17$) it is in the range of the experiment.³ The zoomed views (color online) are 7% of images as in Fig. 1: after 40 ML are deposited with $F = 0.05$ ML/s at $T = 425$ K.

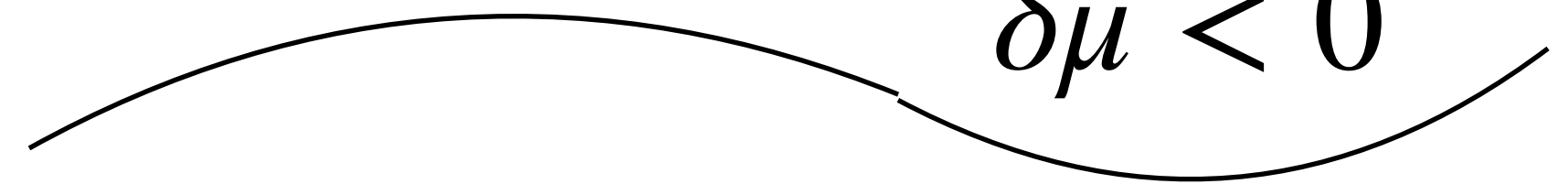
Growth Instabilities at Long Length Scales

- At scales $\ell \gg a$, where a is the lattice spacing, a crystal surface can be treated as a continuous object $z = z(x, y)$
- Assuming as reference state a flat surface, $z(x, y) = z_0$, an undulating surface is characterized by an excess chemical potential, proportional to the local

curvature $\nabla^2 z$:
$$\delta\mu = -\tilde{\sigma} \left(\frac{\partial^2 z}{\partial x^2} + \frac{\partial^2 z}{\partial y^2} \right)$$

$$\delta\mu > 0$$

$$\delta\mu < 0$$

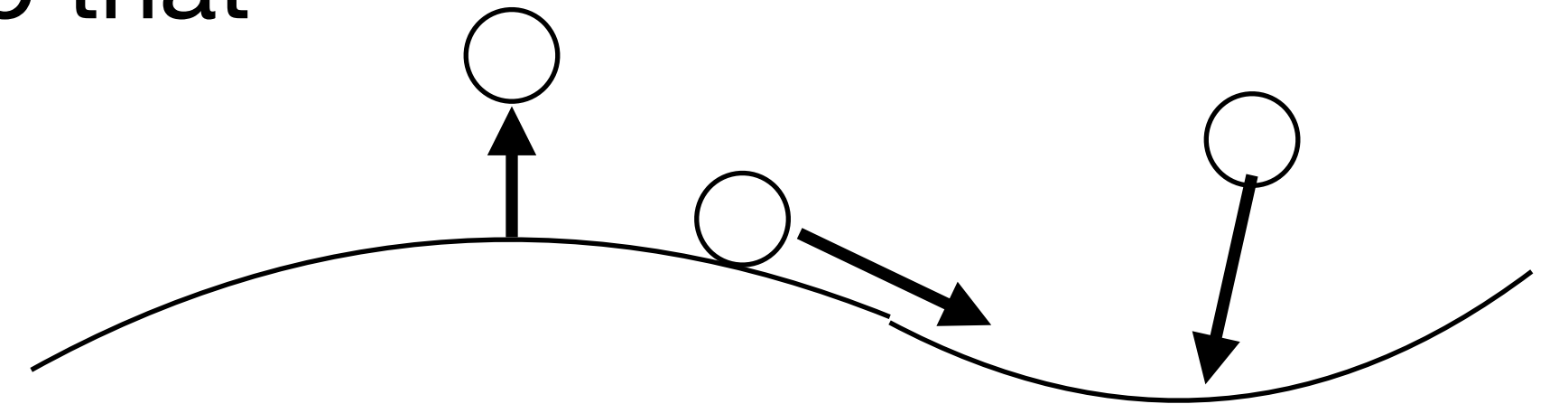


Growth Instabilities at Long Length Scales

- Linear thermodynamics yields two types of dynamics, corresponding to two different transport modes:

- Non conserved: $\dot{z} = -K\delta\mu = K\tilde{\sigma}\nabla^2 z$

- Conserved: $\dot{z} = -\text{div}\vec{J}$; $\vec{J} = -D_s\vec{\nabla}\delta\mu$, so that
 $\dot{z} = -\tilde{\sigma}D_s\nabla^4 z$

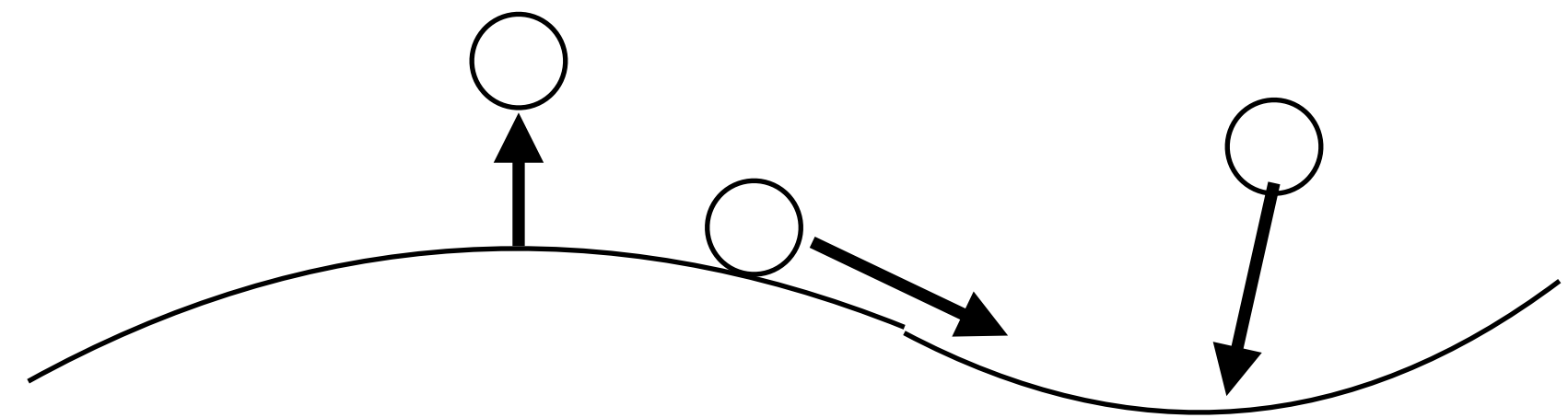


- These equations can describe the growth, dissolution or equilibration kinetics of a rough surface at long length scales

Growth Instabilities at Long Length Scales

- In particular, if $\tau(L)$ is the characteristic time scale of the surface evolution at length scale L :
 - Non conserved: $\dot{z} = K\tilde{\sigma}\nabla^2 z$ implies $\tau(L) \sim L^2$
 - Conserved: $\dot{z} = -\tilde{\sigma}D_s\nabla^n z$ implies $\tau(L) \sim L^n$
 - Ex. Kardar-Parisi-Zhang Eq.

$$\frac{\partial h(\vec{x}, t)}{\partial t} = \nu \nabla^2 h + \frac{\lambda}{2} (\nabla h)^2 + \eta(\vec{x}, t)$$



Growth Instabilities at Long Length Scales

The surface morphology can be quantified by measuring its height-height correlation function

$$G(r, t) = \overline{\langle [h(x + r, t) - h(x, t)]^2 \rangle}$$

Such function is observed to possess the scaling form

$$G(r, t) = t^{2\beta} g(r/\xi)$$

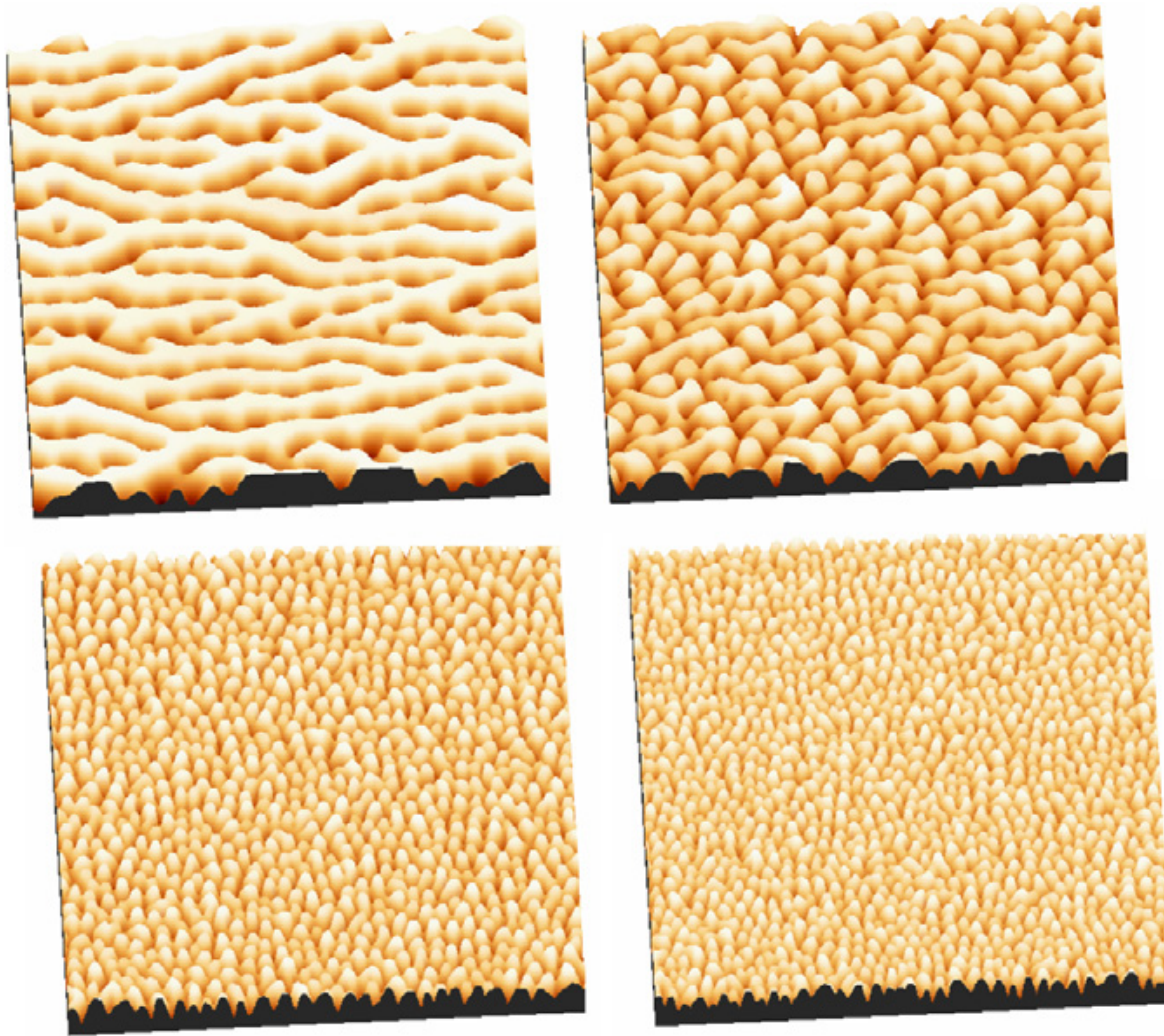
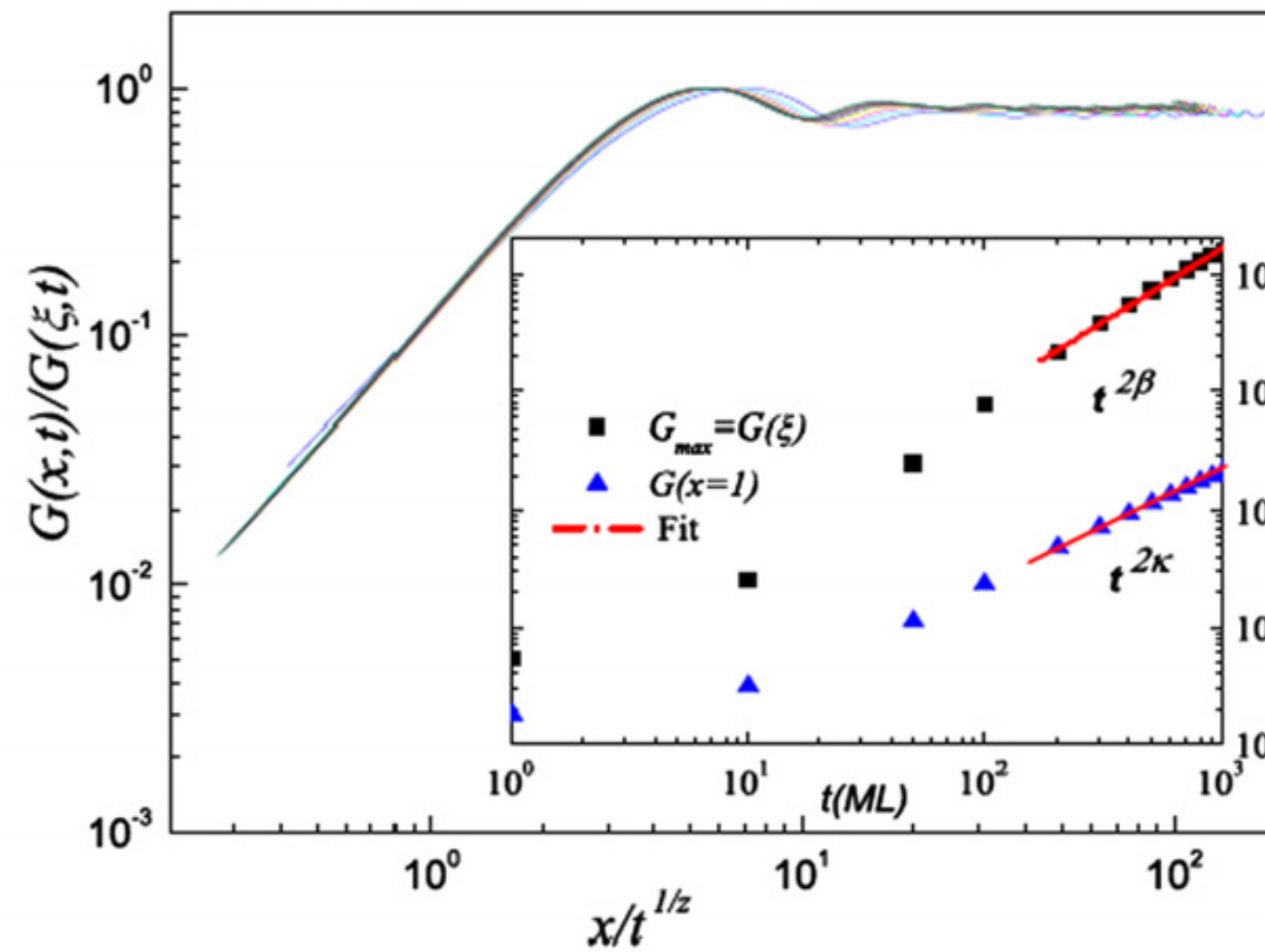
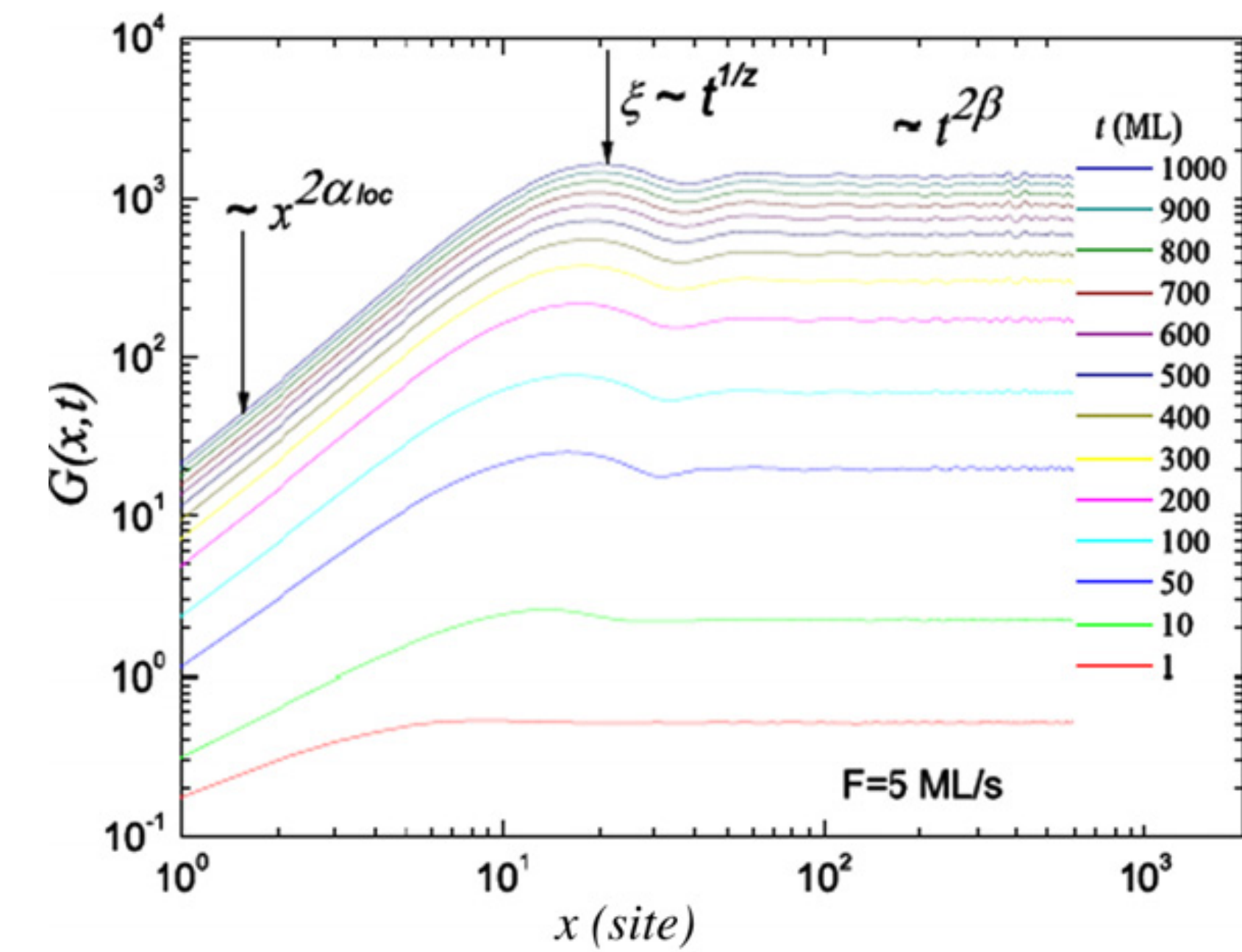
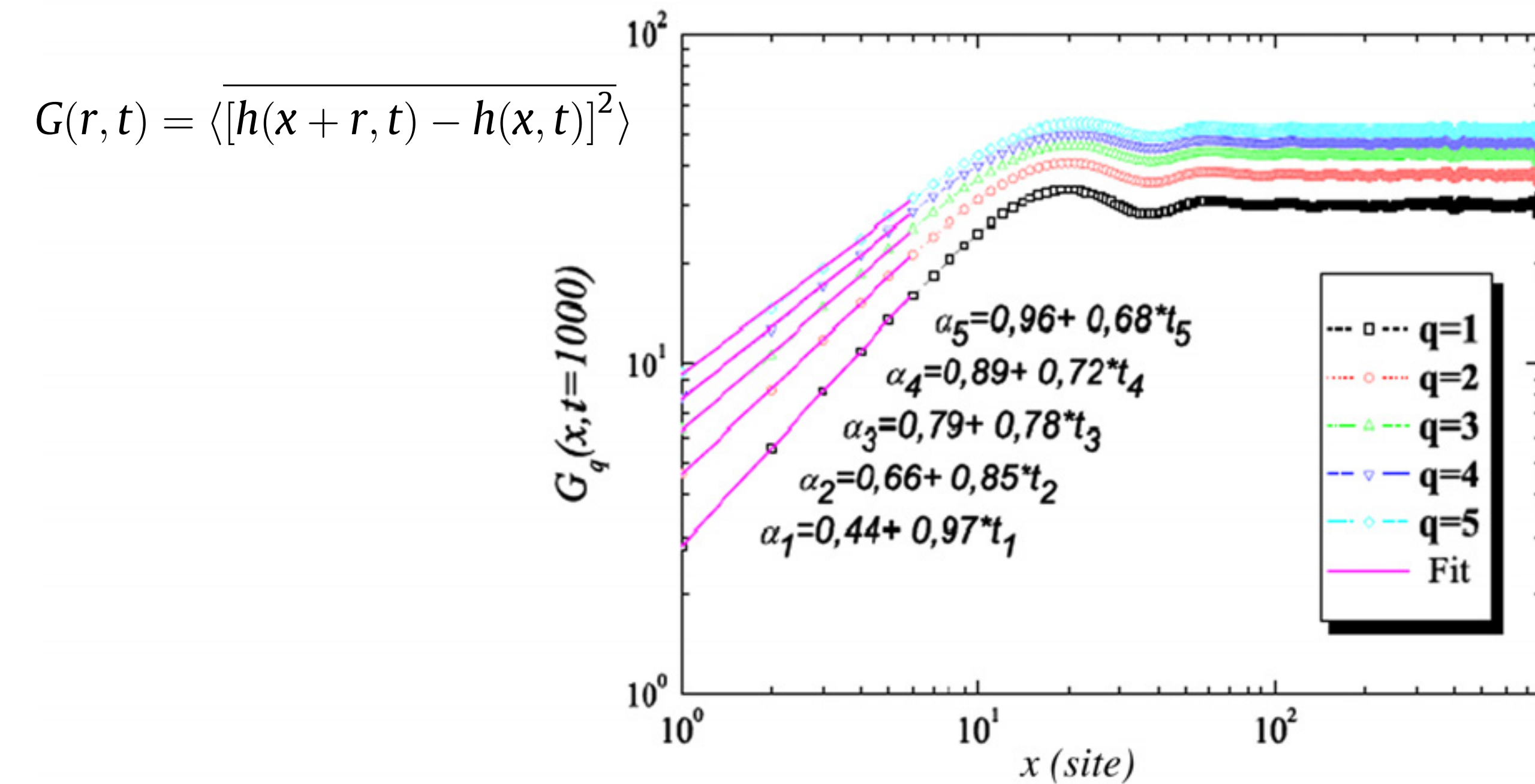


Fig. 5. Monte Carlo simulation images of the surface morphology evolution with increasing deposition flux: $F = 1, 5, 20, 40$ ML/s. The initial surface is a vicinal (800×800) sites (in lattice unit) with terrace width $L = 5$ sites.



$$G(r, t) = t^{2\beta} g(r/\xi)$$



$$\frac{\partial}{\partial t} h(x, y, t) = v_x \frac{\partial^6}{\partial x^6} h(x, y, t) + v_y \frac{\partial^6}{\partial y^6} h(x, y, t) + \eta(x, y, t),$$

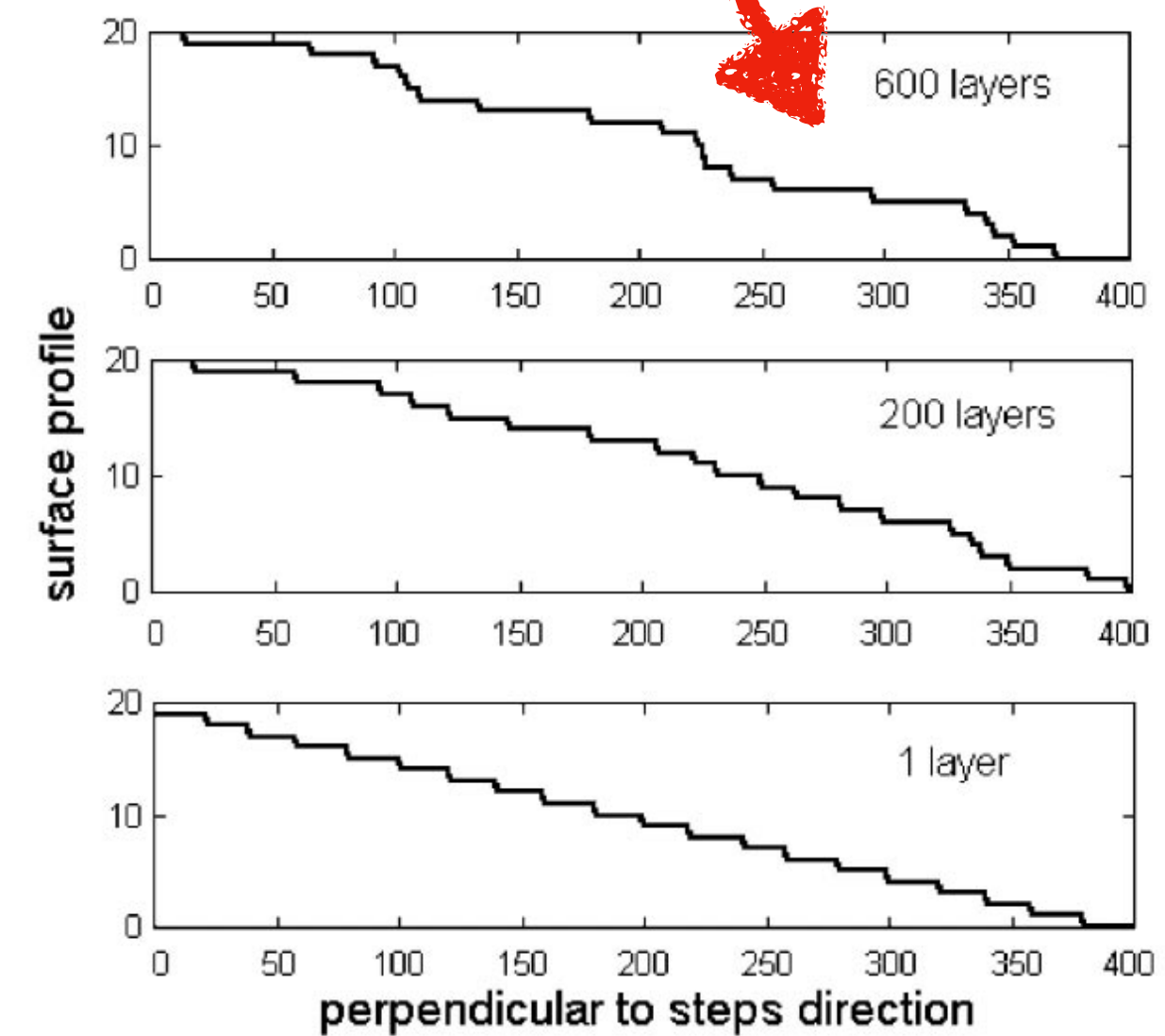
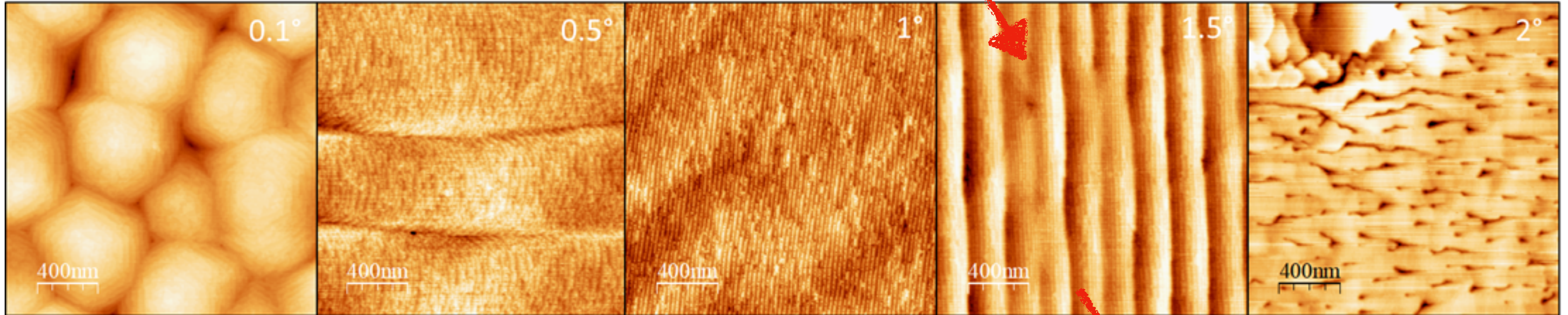
$$\text{with } \langle \eta(x, y, t) \rangle = 0,$$

and

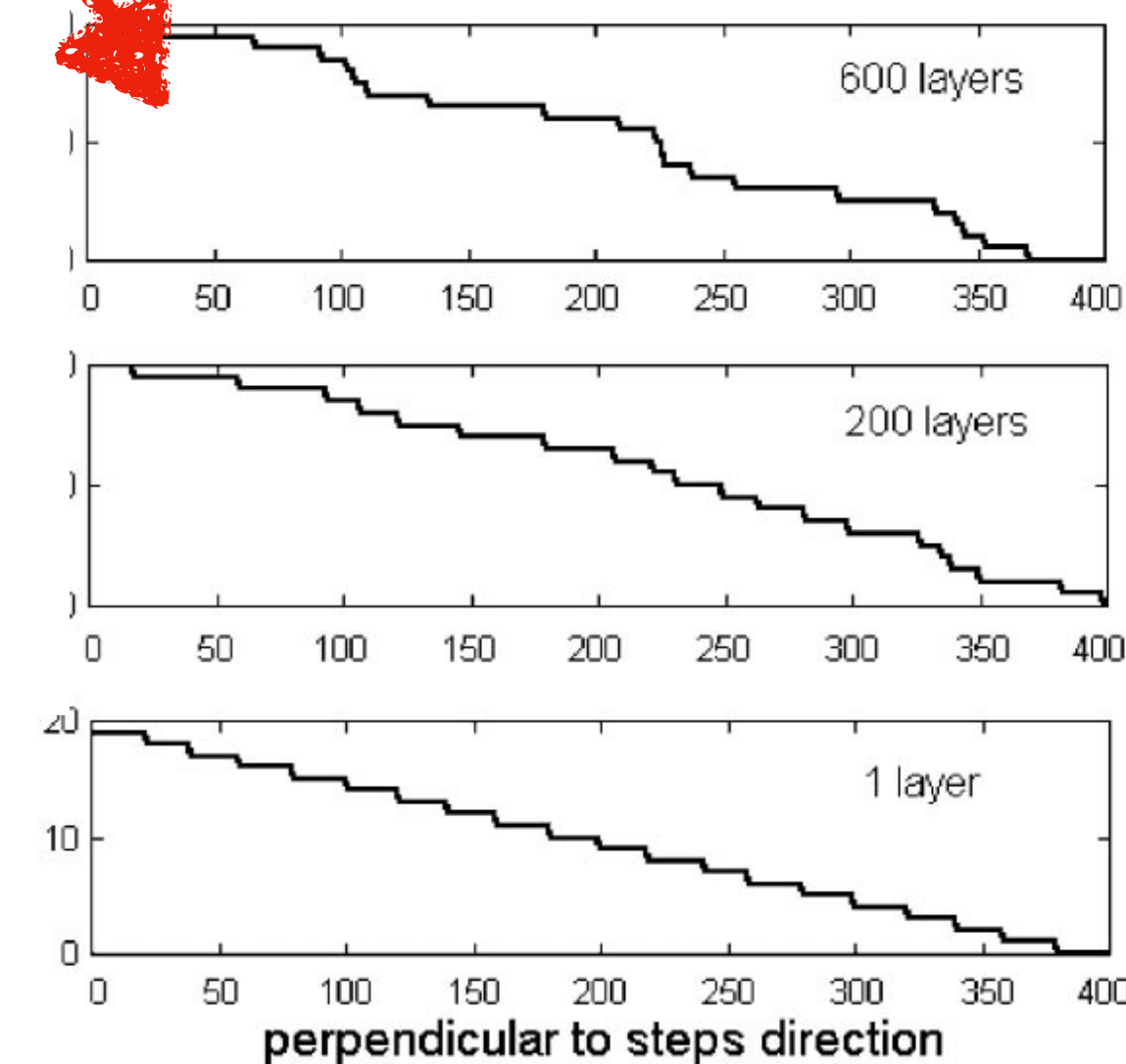
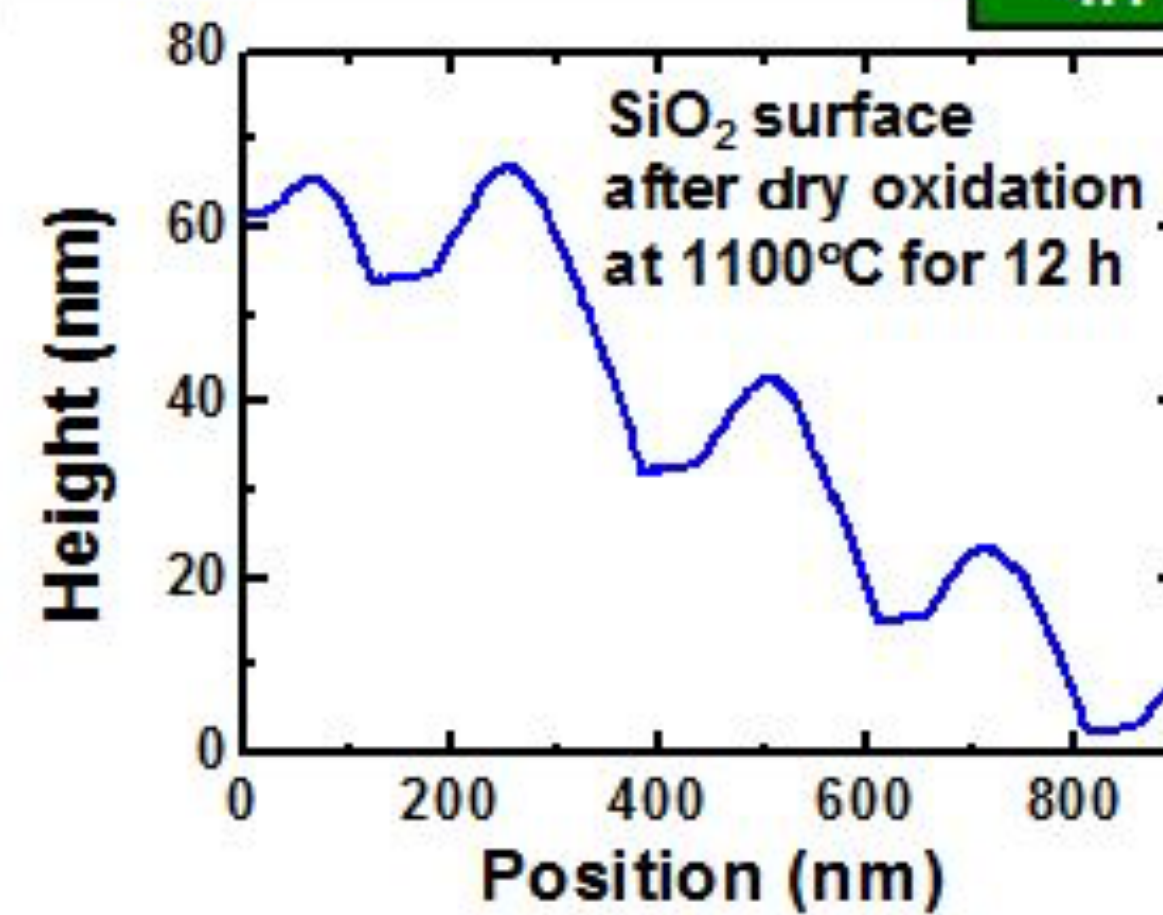
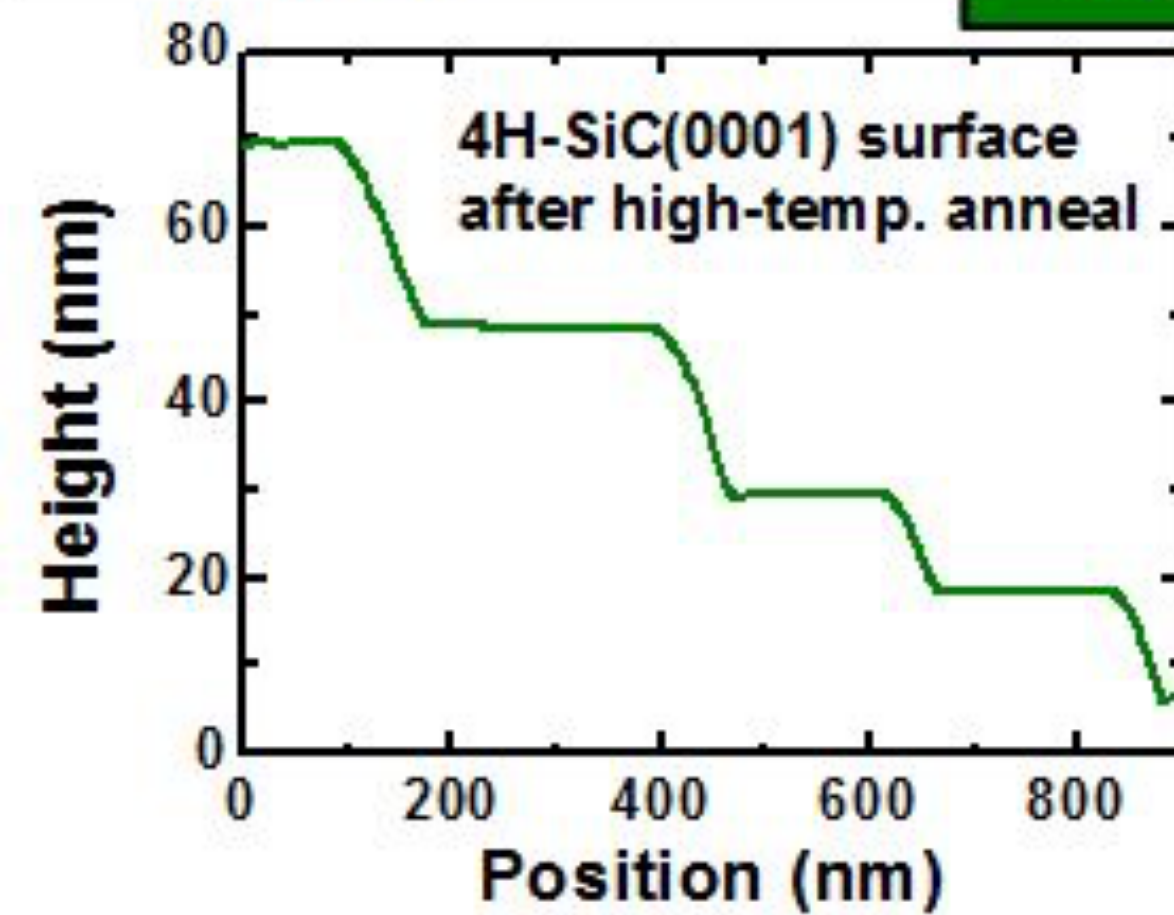
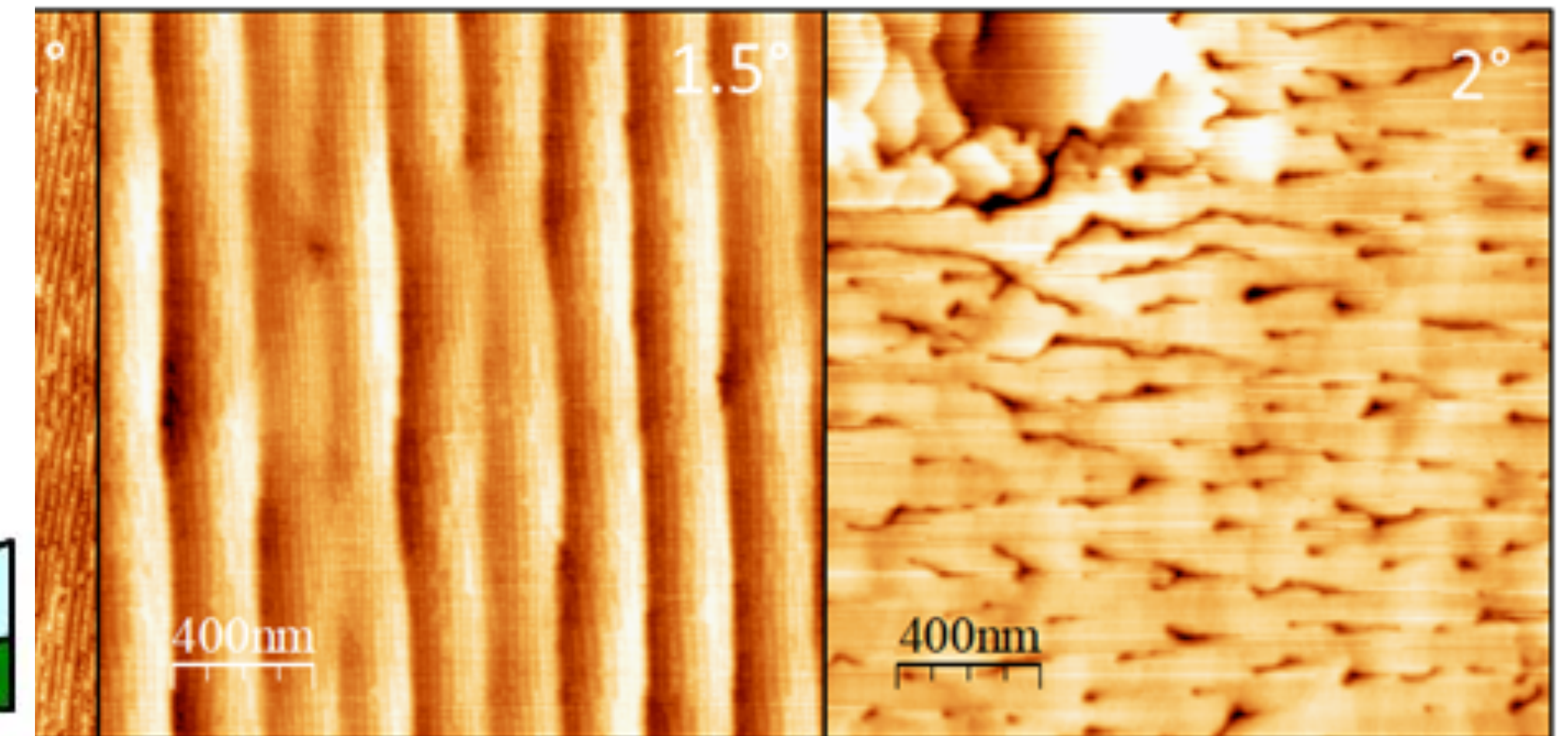
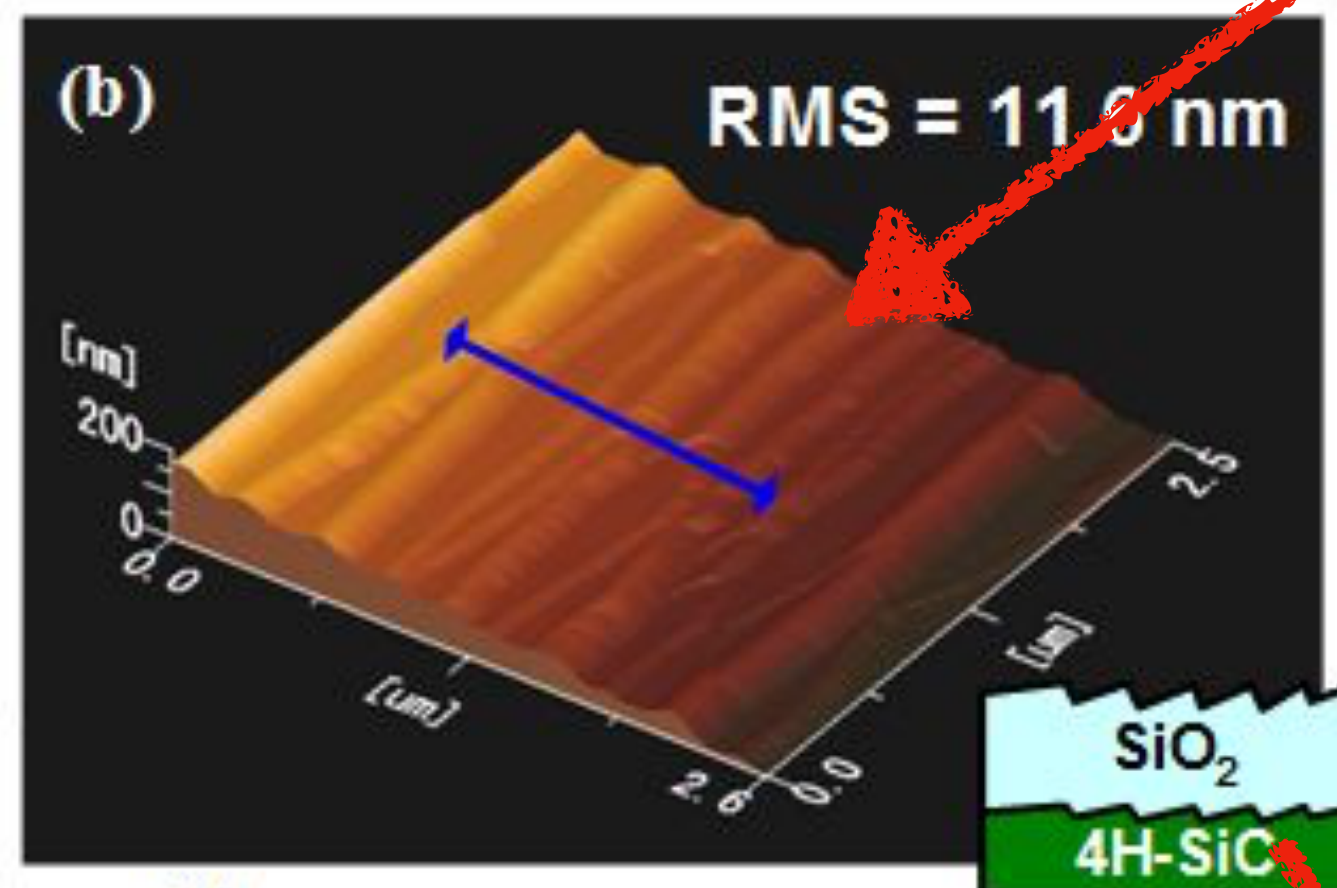
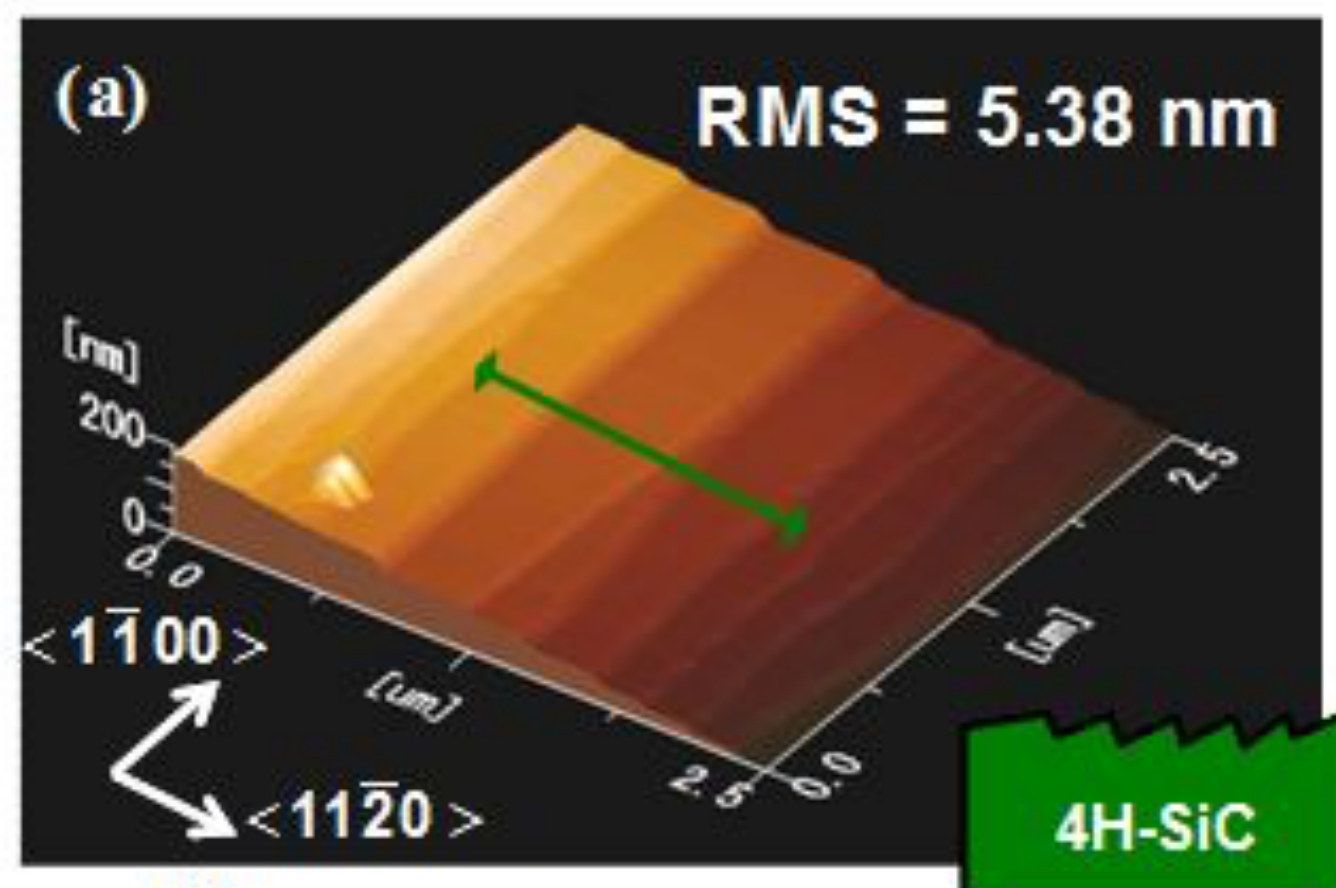
$$\langle \eta(x, y, t) \eta(x', y', t') \rangle = D((x - x')^2 + (y - y')^2)^{(2\rho-1)/2} |t - t'|^{2\theta-1}.$$

Multiscaling!

Other Instabilities: Step Bunching



Other Instabilities: Step Bunching



Inverse Schwoebel effect (iES)

R.L. Schwoebel, J. Appl. Phys. 40 (1968) 614.

Bunching: Net current “down”

$$J_{\text{down}} > J_{\text{up}}$$

Inverse Schwoebel effect (iES)

$$v_n = F(x_n - x_{n-1})$$

v_n

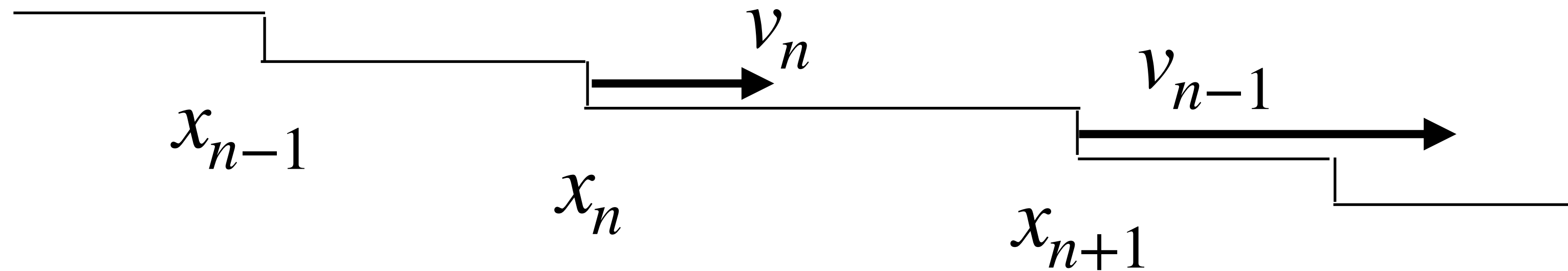
Extreme case: infinite iES effect

Net current “down”

$$J_{\text{down}} > J_{\text{up}}$$

Inverse Schwoebel effect (iES)

$$v_n = F(x_n - x_{n-1})$$



Atoms can only “step down”

$$J_{\text{up}} = 0$$

Inverse Schwoebel effect (iES)

$$v_n = F(x_n - x_{n-1})$$

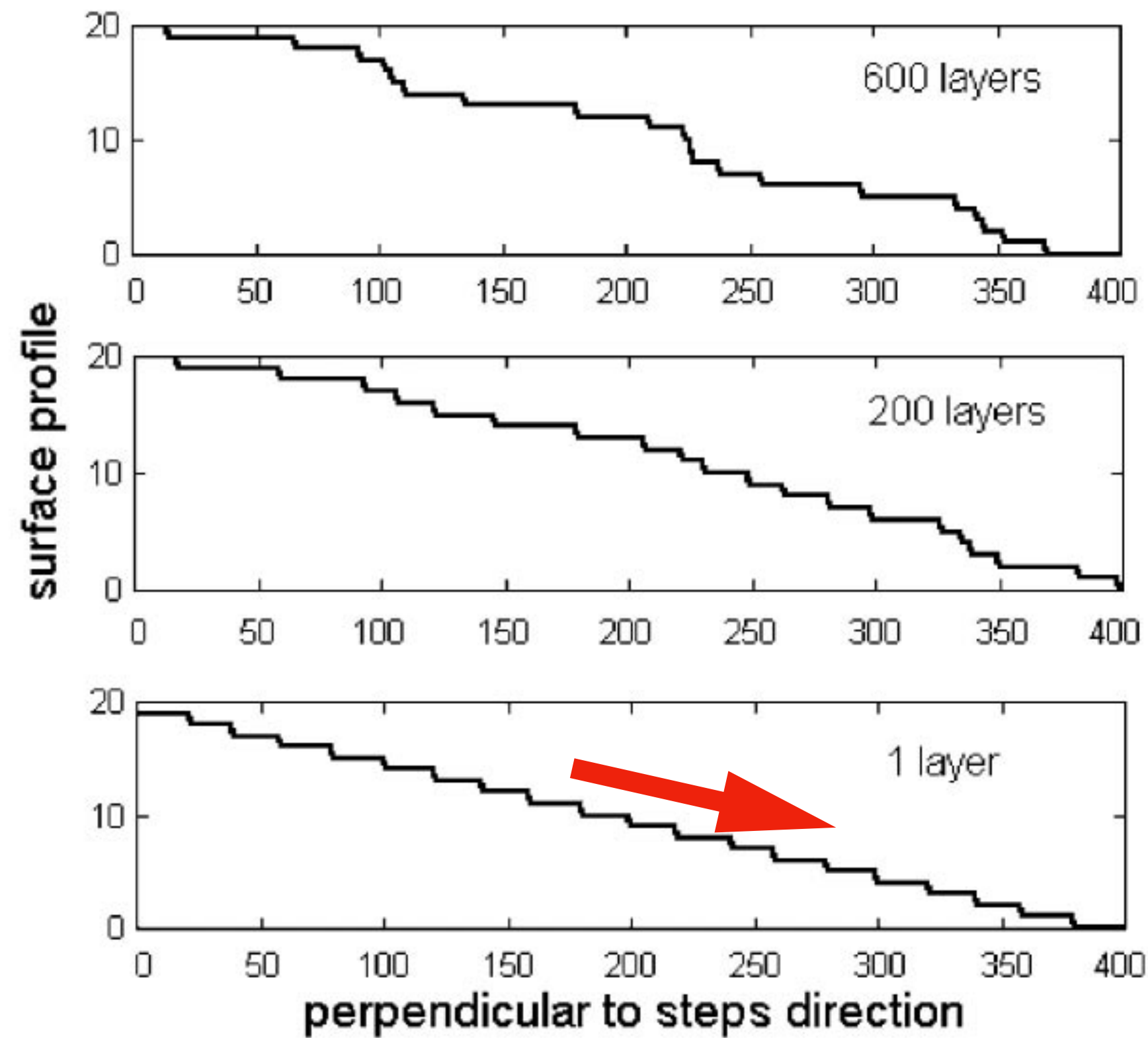
v_n

Attractive kinetic interaction

Net current “down”

$$J_{\text{down}} > J_{\text{up}}$$

Inverse Schwoebel effect



$$J_{\text{down}} > J_{\text{up}}$$

Net current “down”

Steps “attract”

Leads to step bunching

Common wisdom is:

- **Surface current “up”: step meandering;**
- **Surface current “down”: step bunching**

Bunching and meandering cannot coexist

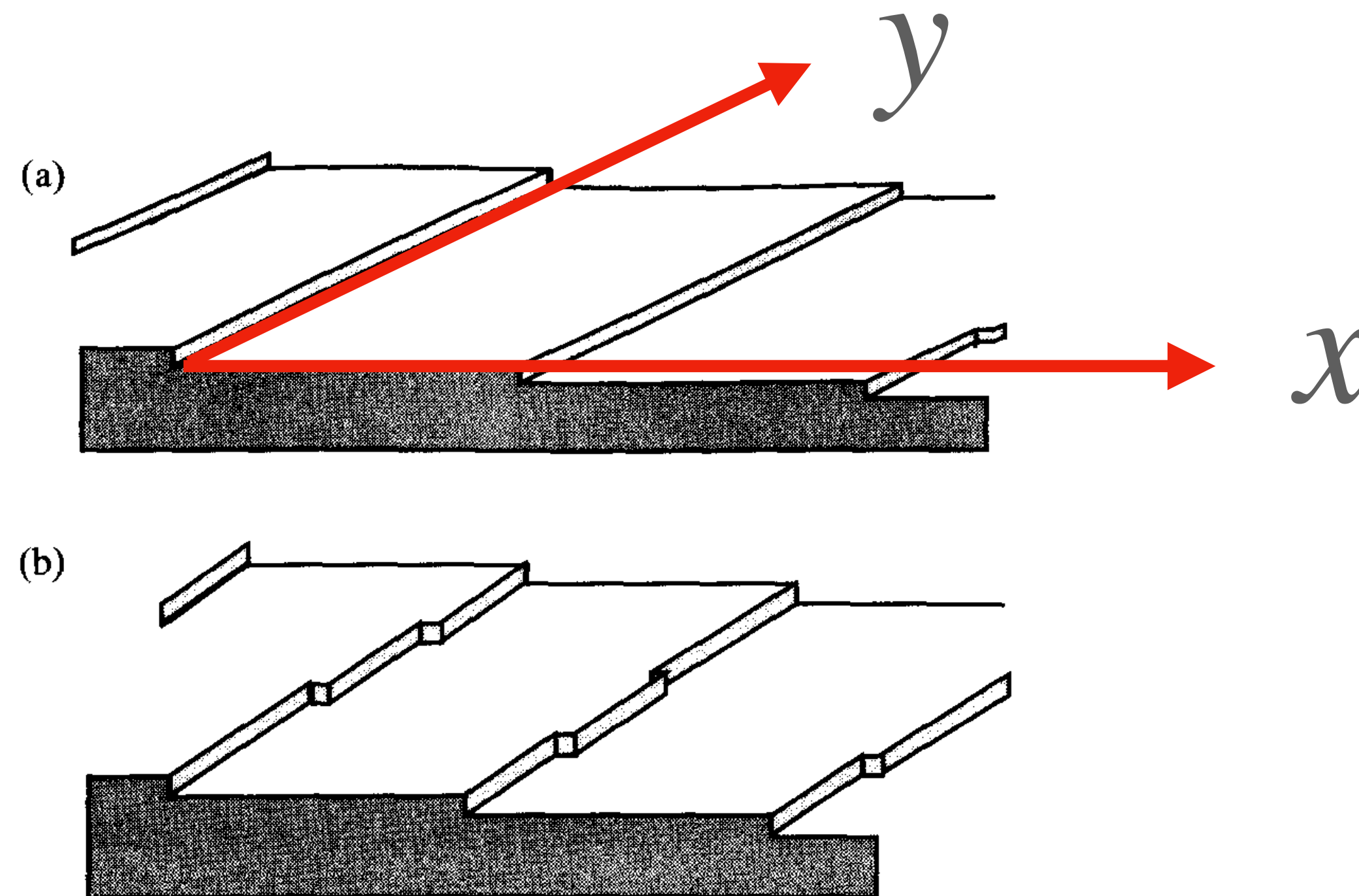
Common wisdom is:

- **Surface current “up”: step meandering;**
- **Surface current “down”: step bunching**

Linear stability analysis

Common wisdom is:

- Surface current “up”: step meandering;
- Surface current “down”: step bunching



Common wisdom is:

- **Surface current “up”: step meandering;**
- **Surface current “down”: step bunching**

$$\frac{\partial h}{\partial t} = \nu_{\parallel} \frac{\partial^2 h}{\partial y^2} + \nu_{\perp} \frac{\partial^2 h}{\partial x^2}$$

Common wisdom is:

- Surface current “up”: step meandering;
- Surface current “down”: step bunching

Common wisdom is wrong

$$\nu_{\parallel} = \ell J(\ell)$$

step meandering if $\nu_{\parallel} < 0$

$$\nu_{\perp} = -\ell^2 \frac{dJ(\ell)}{d\ell}$$

step bunching if $\nu_{\perp} < 0$

Common wisdom is:

- Surface current “up”: step meandering;
- Surface current “down”: step bunching

Common wisdom is wrong

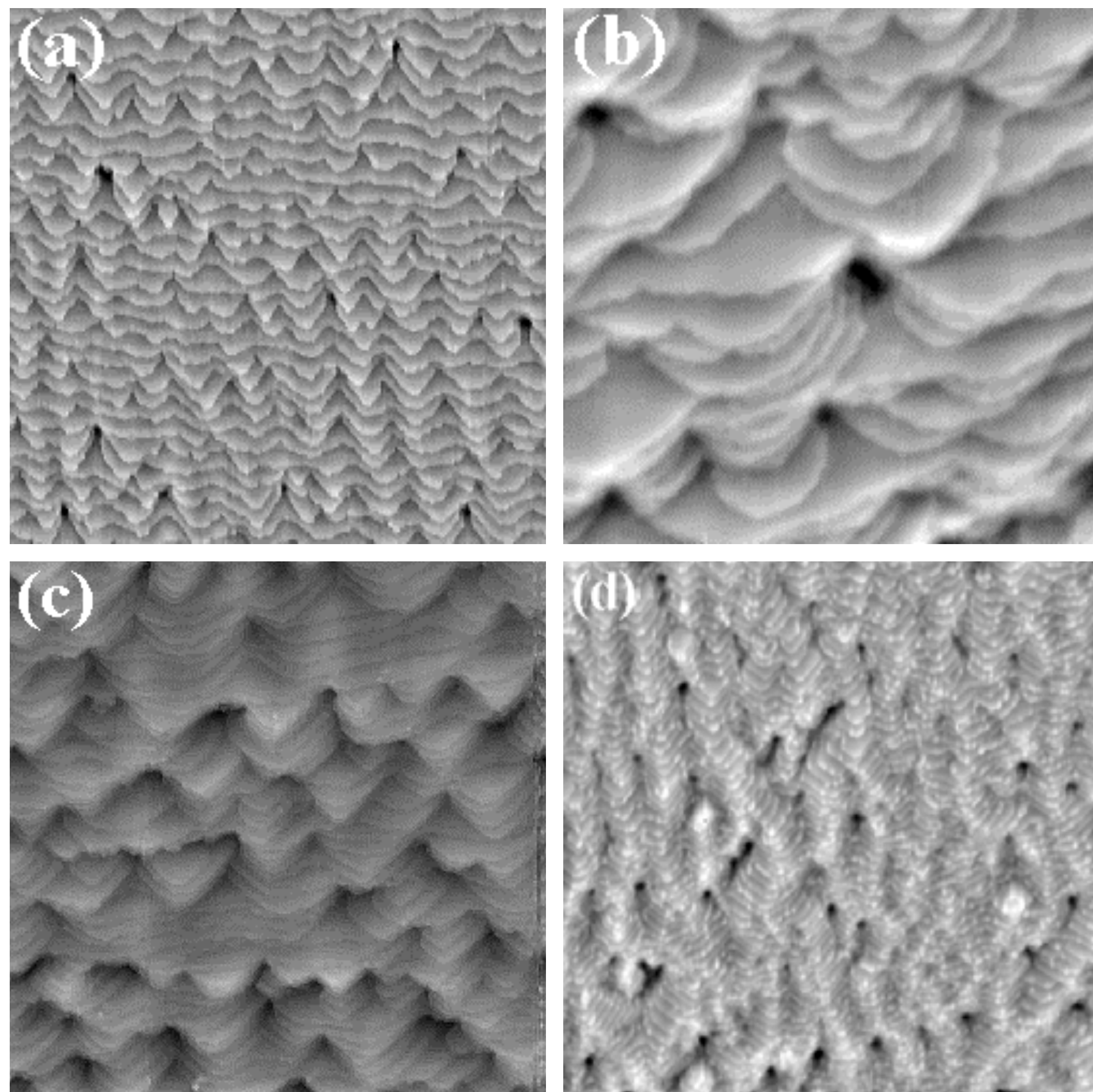
$$\nu_{\parallel} = \ell J(\ell)$$

step meandering if $\nu_{\parallel} < 0$

$$\nu_{\perp} = -\ell^2 \frac{dJ(\ell)}{d\ell}$$

step bunching if $\nu_{\perp} < 0$

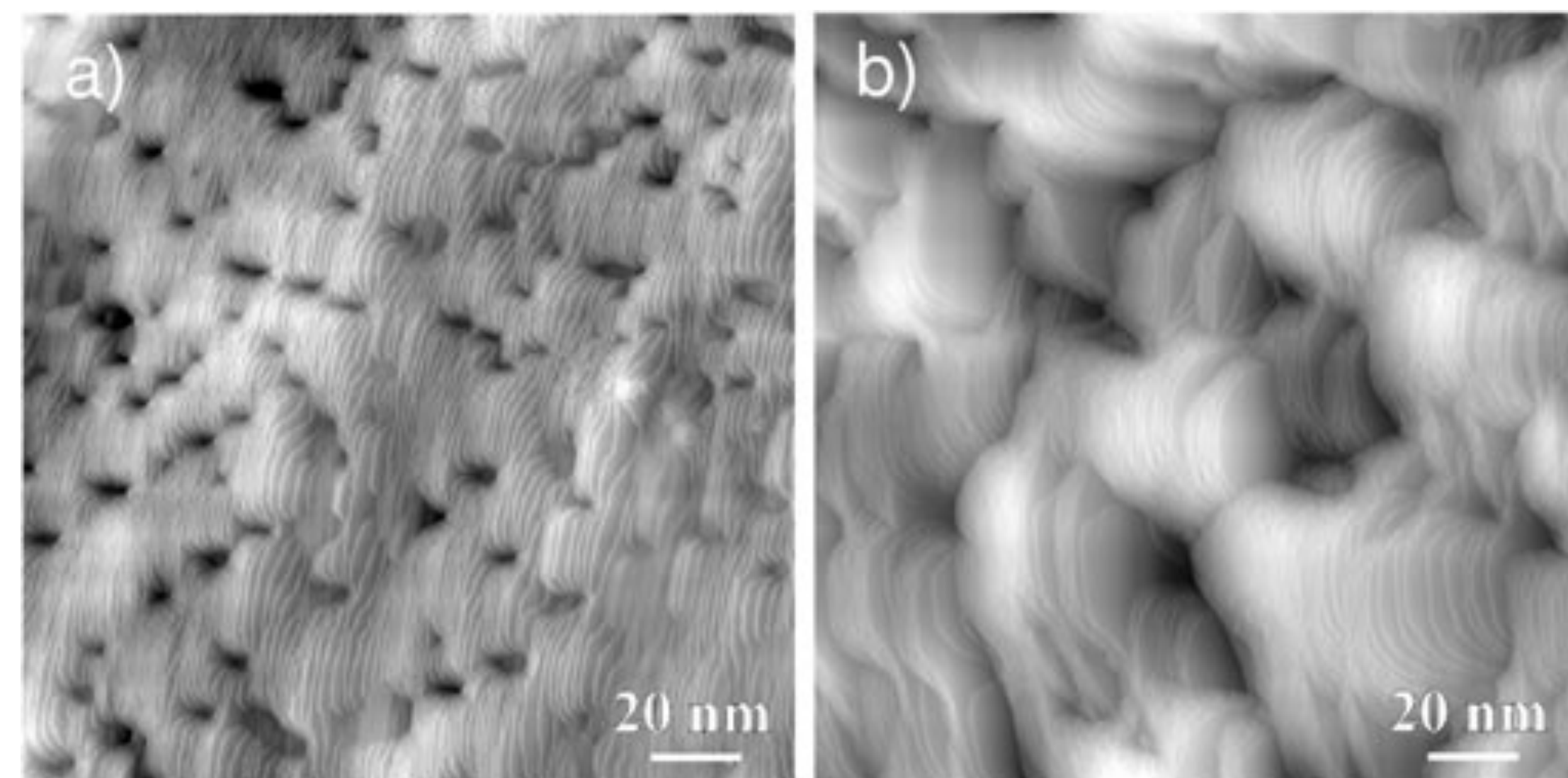
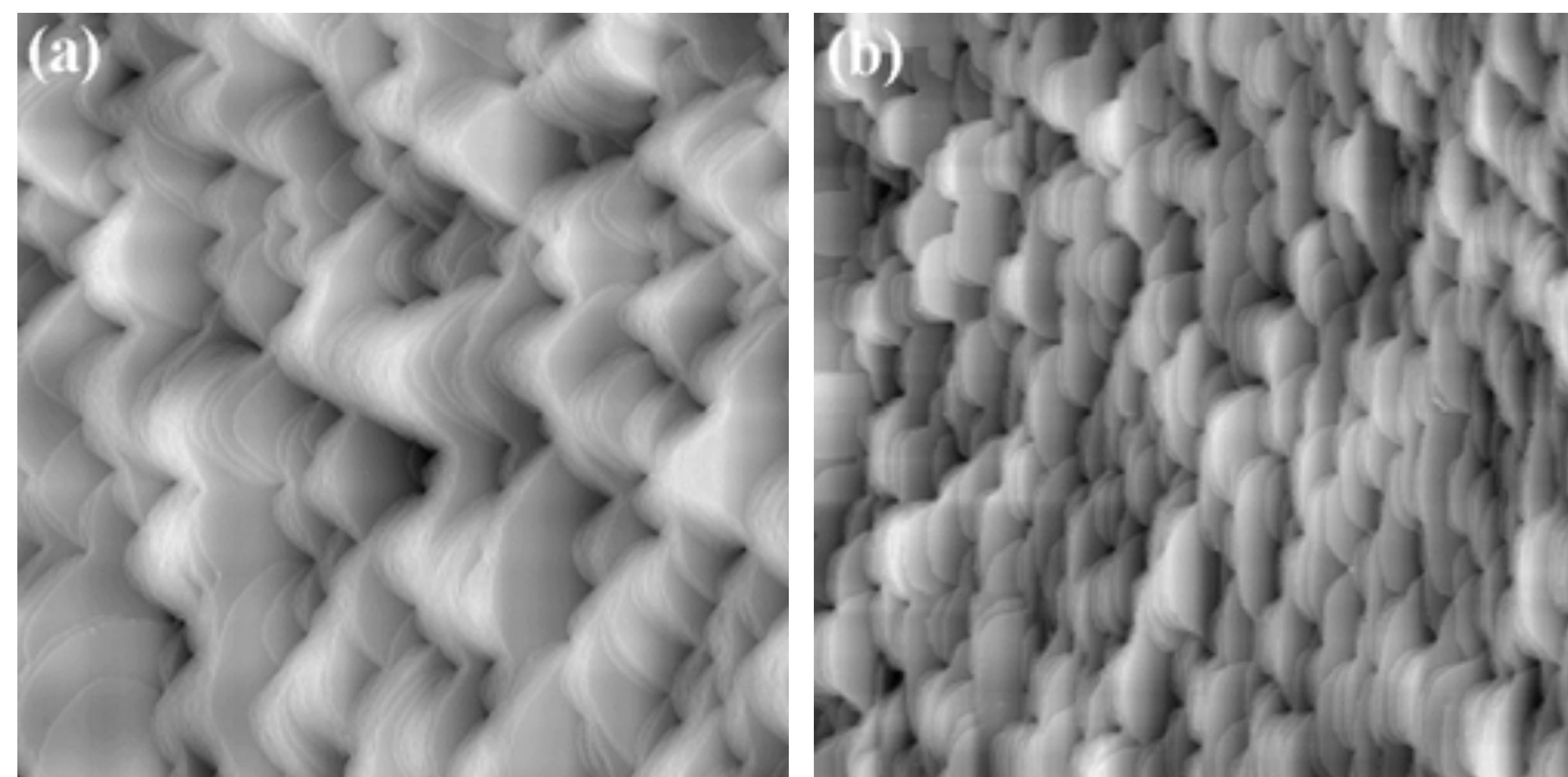
Bunching and meandering can coexist



Spontaneous structural pattern formation at the nanometre scale in kinetically restricted homoepitaxy on vicinal surfaces

N Néel, T Maroutian, L Douillard and H-J Ernst

CEA Saclay, DSM/Drecom/Spesi, 91191 Gif Sur Yvette, France



Complex growth: multiple components



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Available online at www.sciencedirect.com



Journal of Crystal Growth 258 (2003) 1–13



www.elsevier.com/locate/jcrysgr

A new model of morphological instabilities during epitaxial growth: from step bunching to mounds formation

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Received 21 July 2000; accepted 26 September 2000

Two-particle surface diffusion-reaction models of vapour-phase epitaxial growth on vicinal surfaces

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Journal of Crystal Growth 258 (2003) 14–25

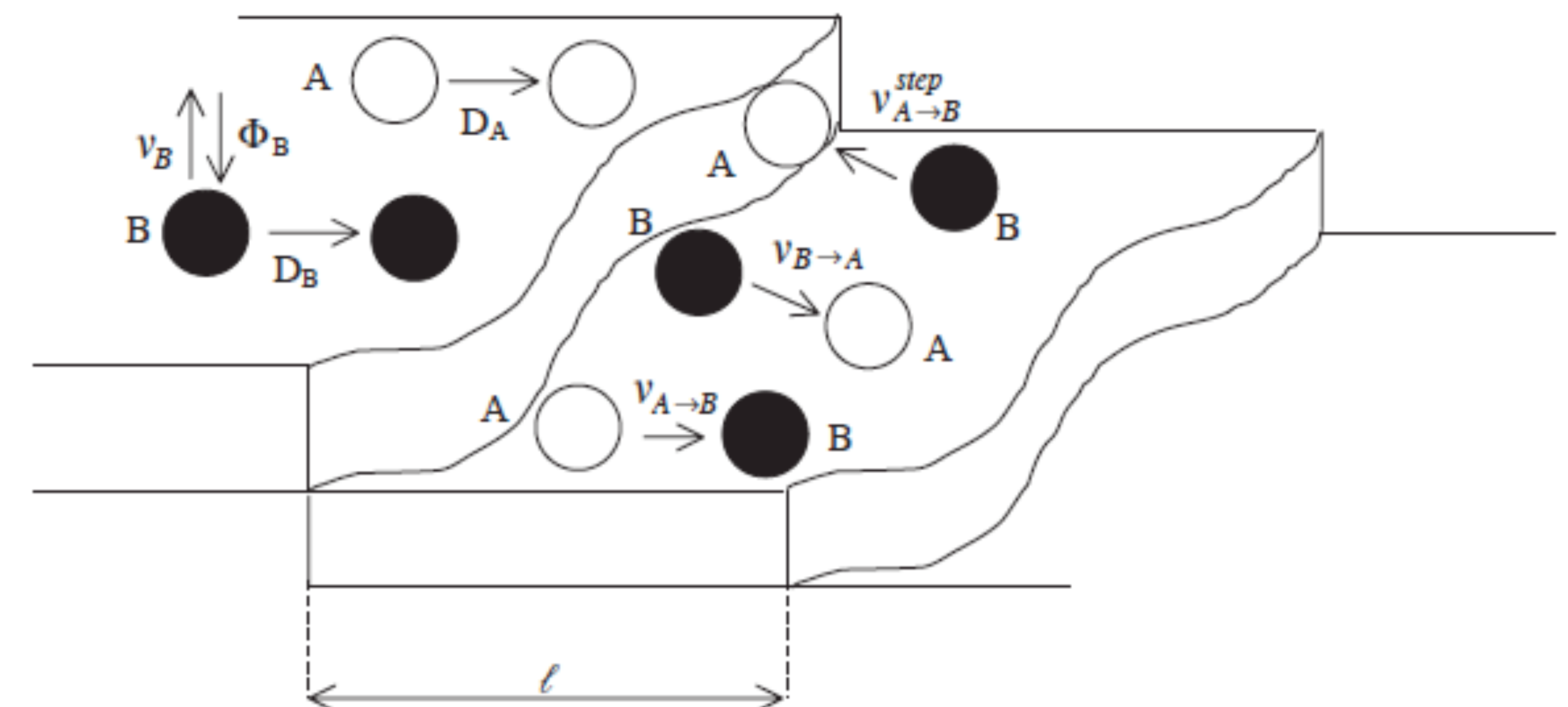


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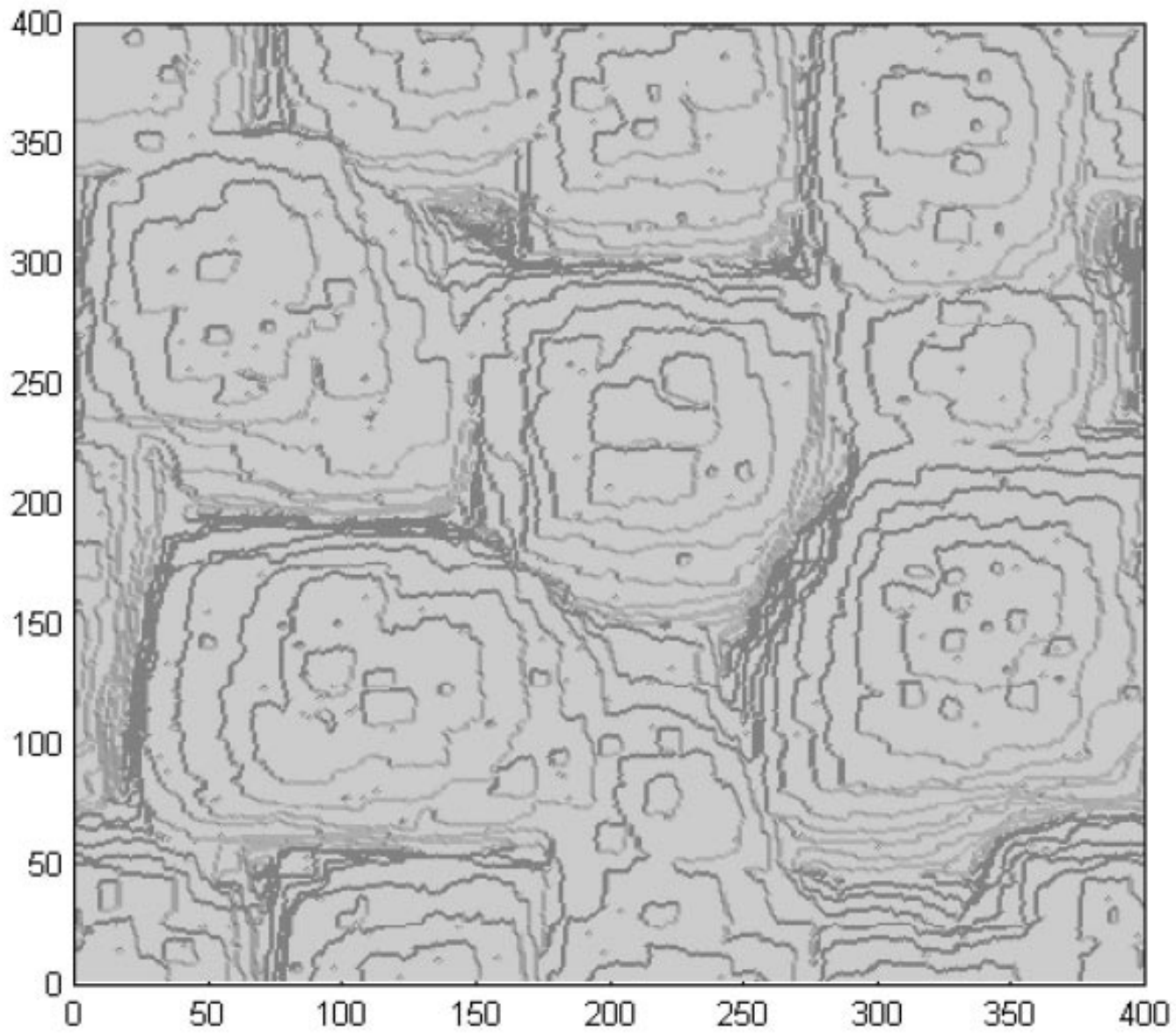
Direct condensation modelling for a two-particle growth system: application to GaAs grown by hydride vapour phase epitaxy

E. Gil-Lafon*, J. Napierala, A. Pimpinelli, R. Cadoret, A. Trassoudaine, D. Castelluci

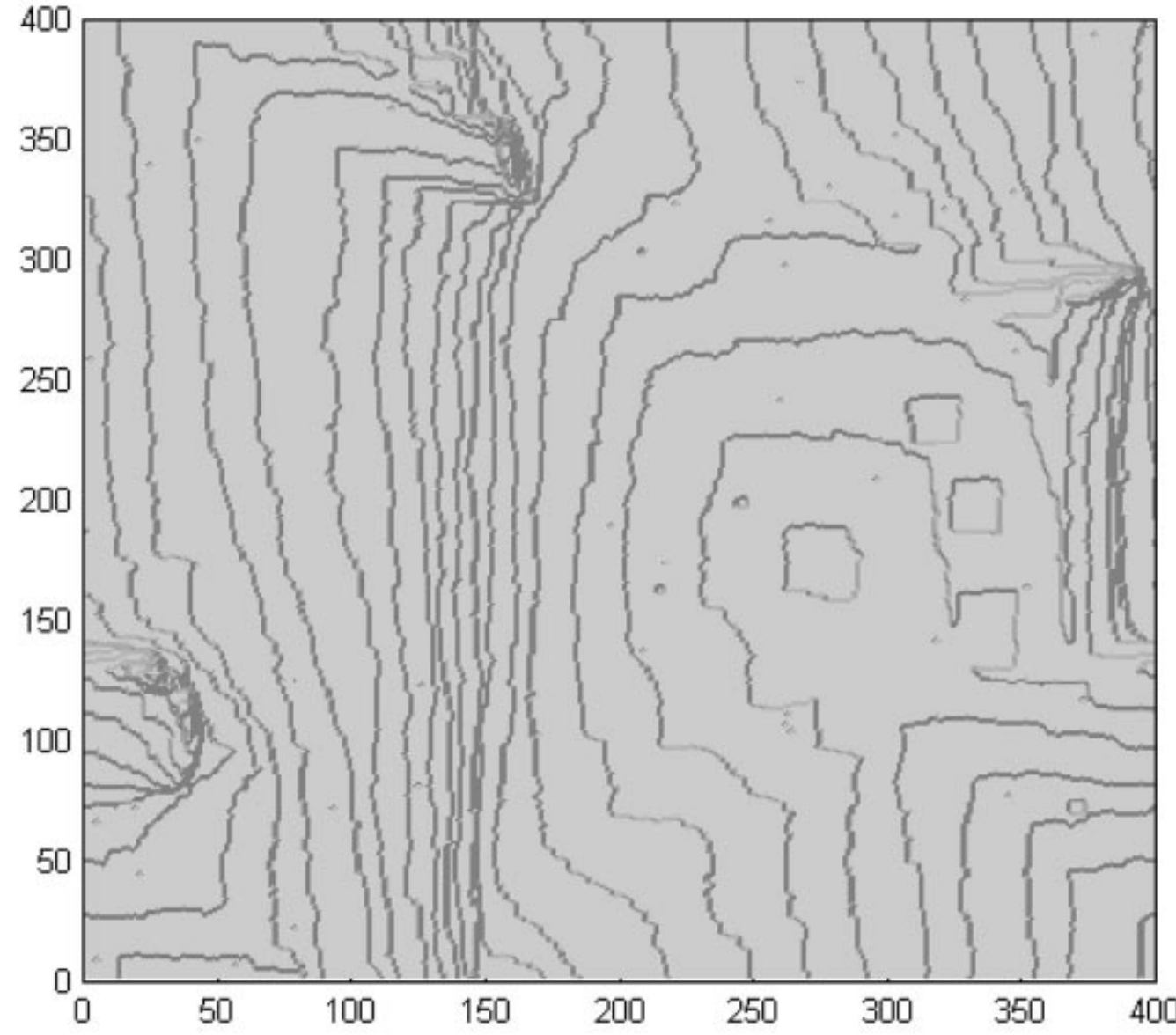
LASMEA, UMR CNRS 6602, Université Blaise Pascal, Clermont-Ferrand II, Complexe Scientifique des Cèzeaux, Les Cèzeaux, F-63177 Aubière Cedex, France



Mounds

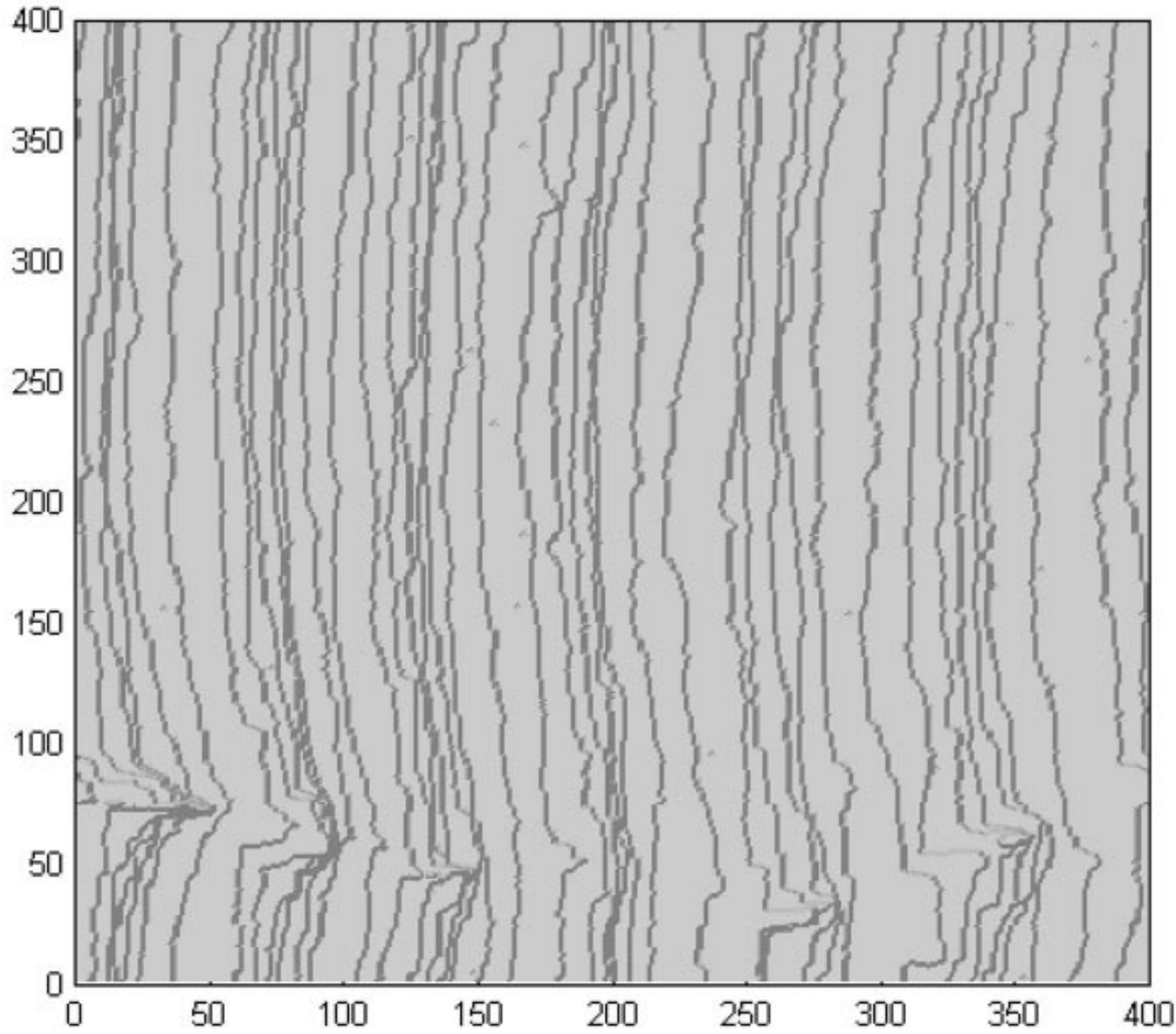


(a)

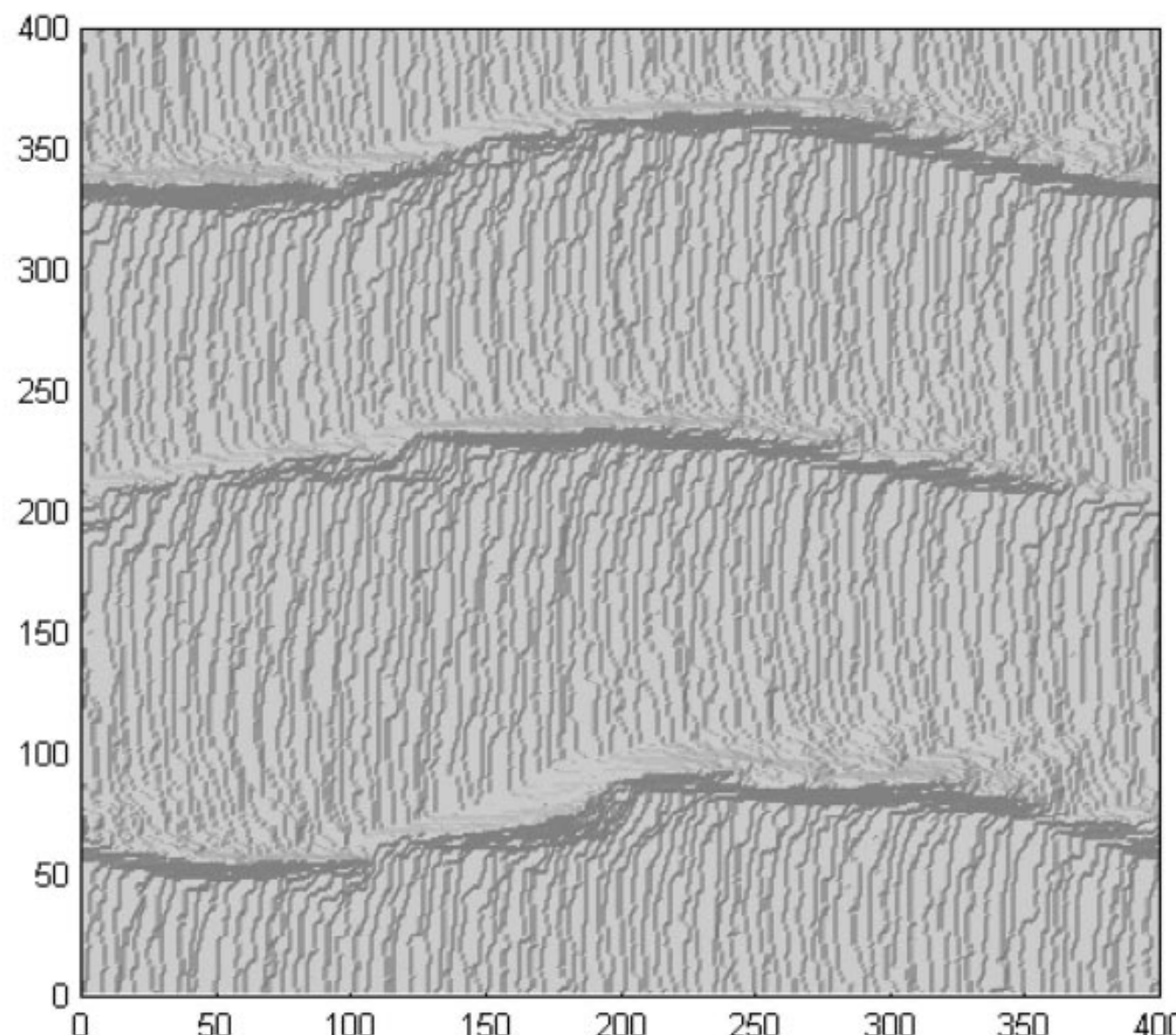


(b)

Bunches



(c)



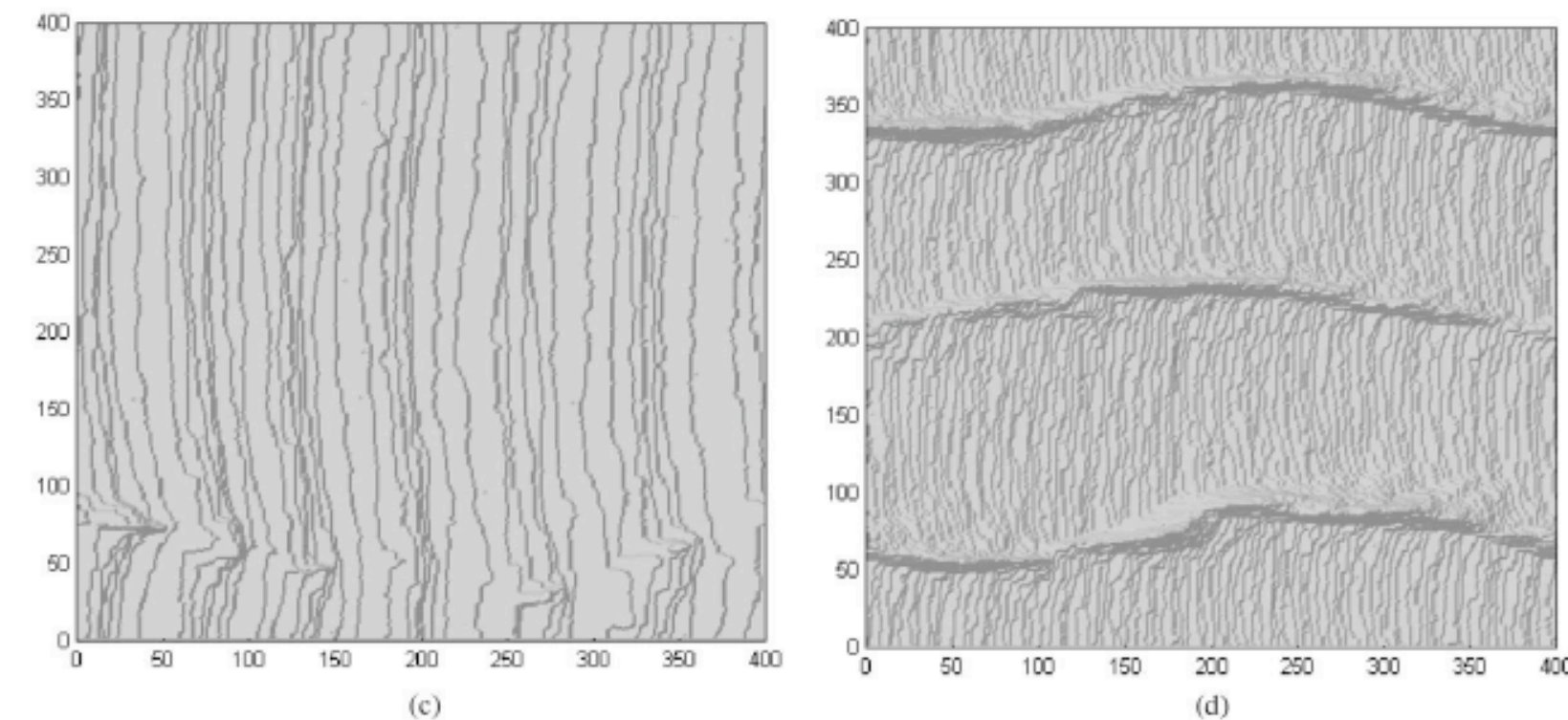
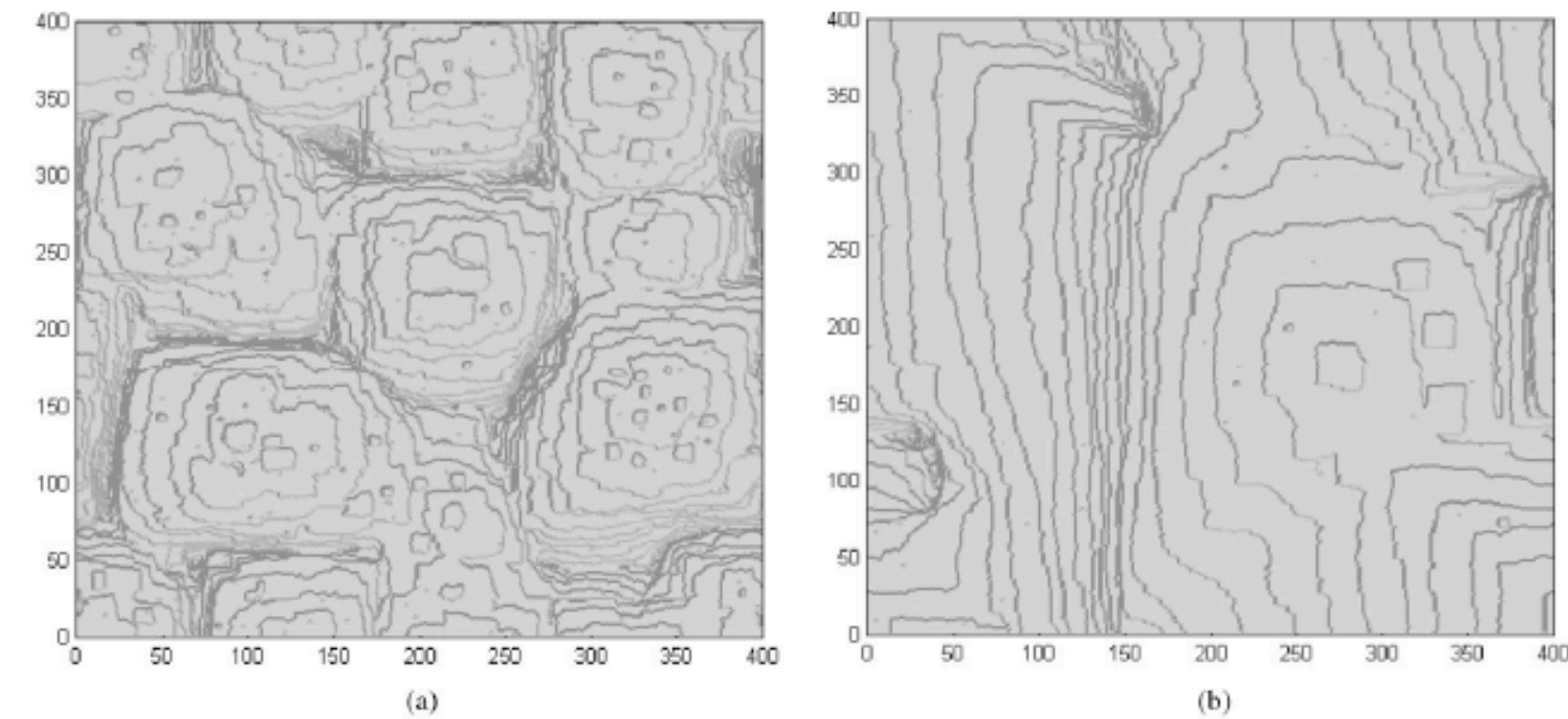
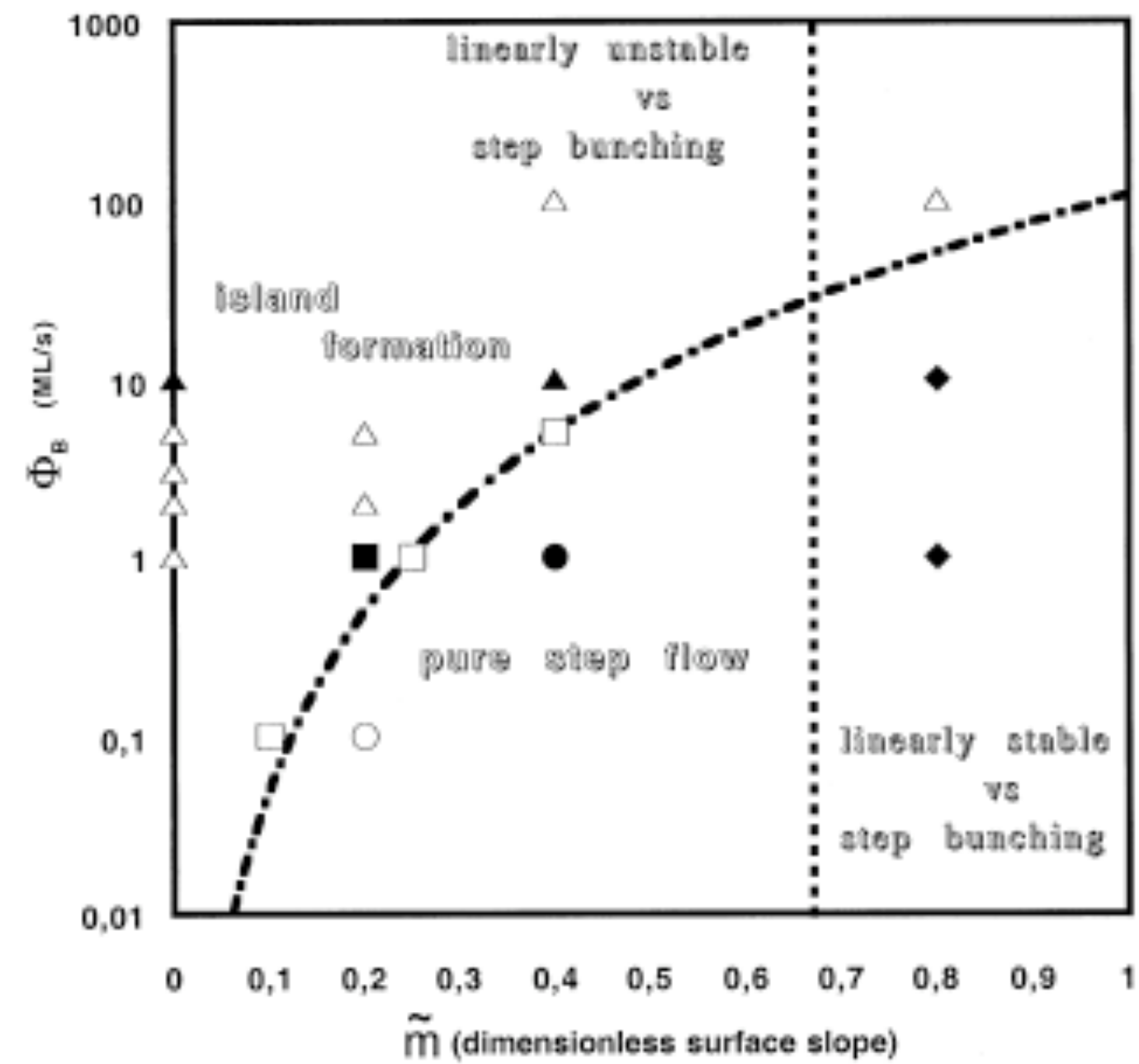
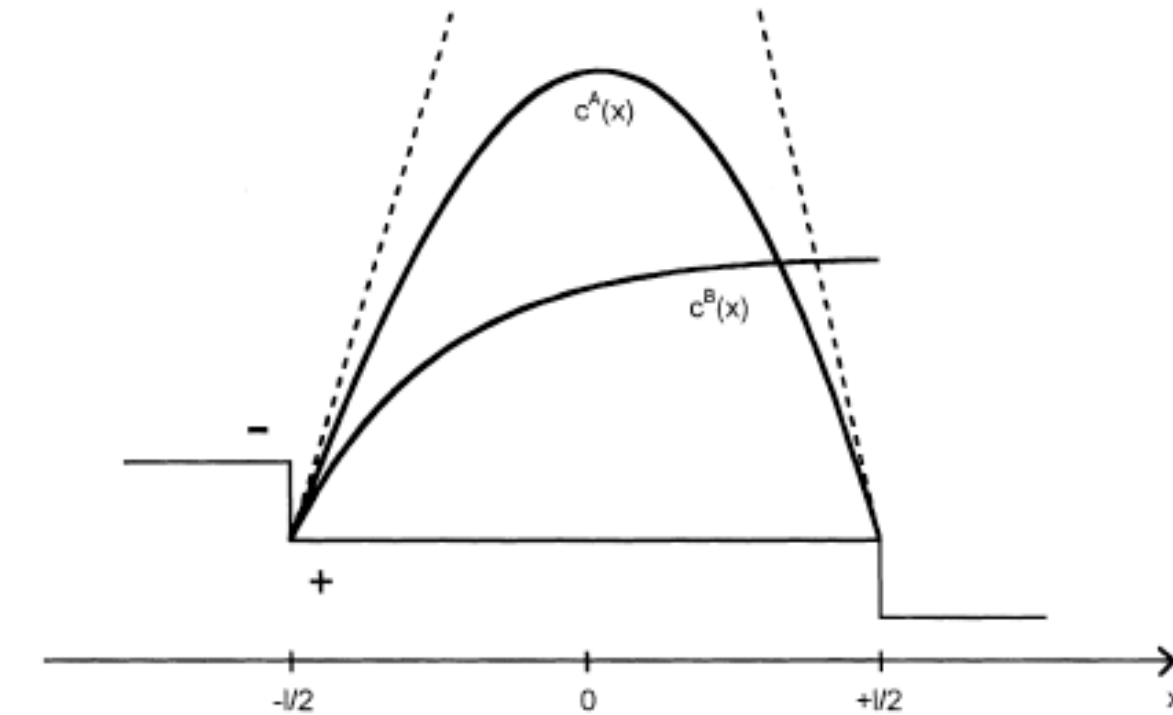
(d)

Meanders

ES barrier at step edges implies coexistence of mounding, bunching *and* meandering...

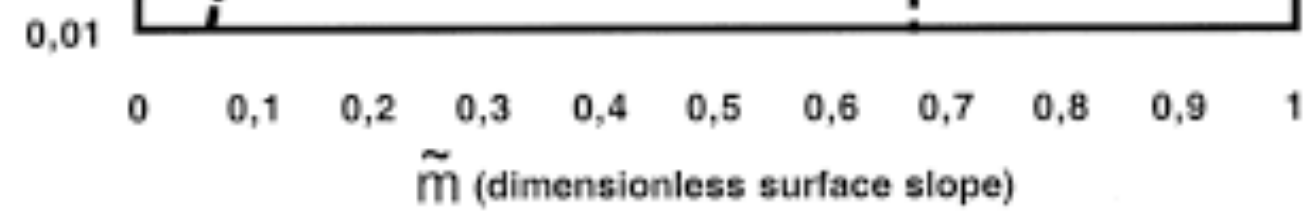
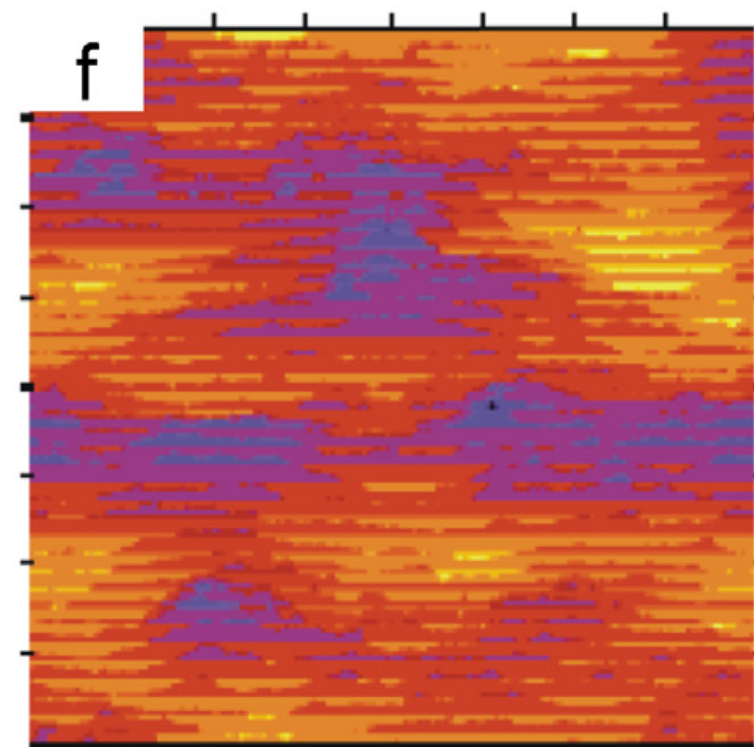
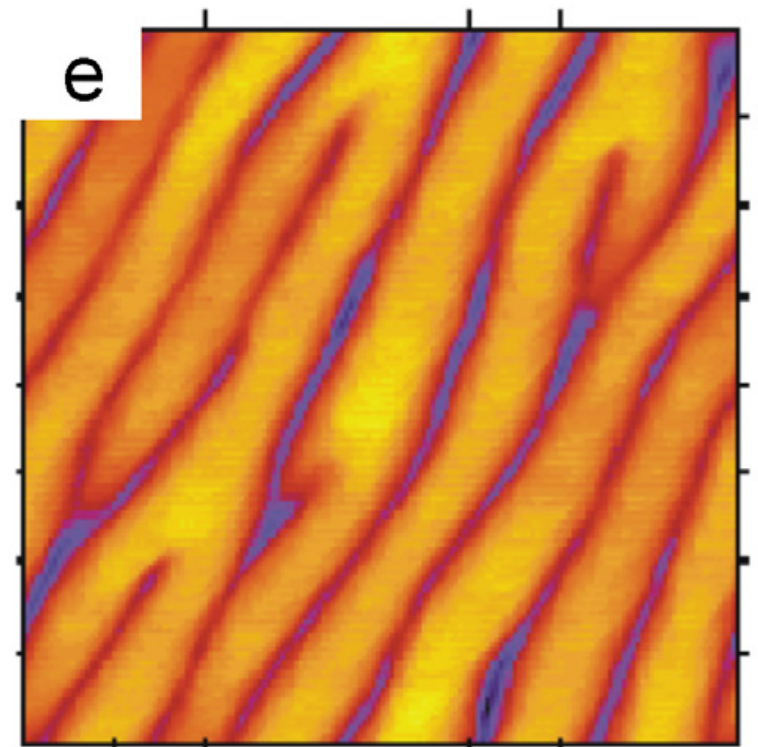
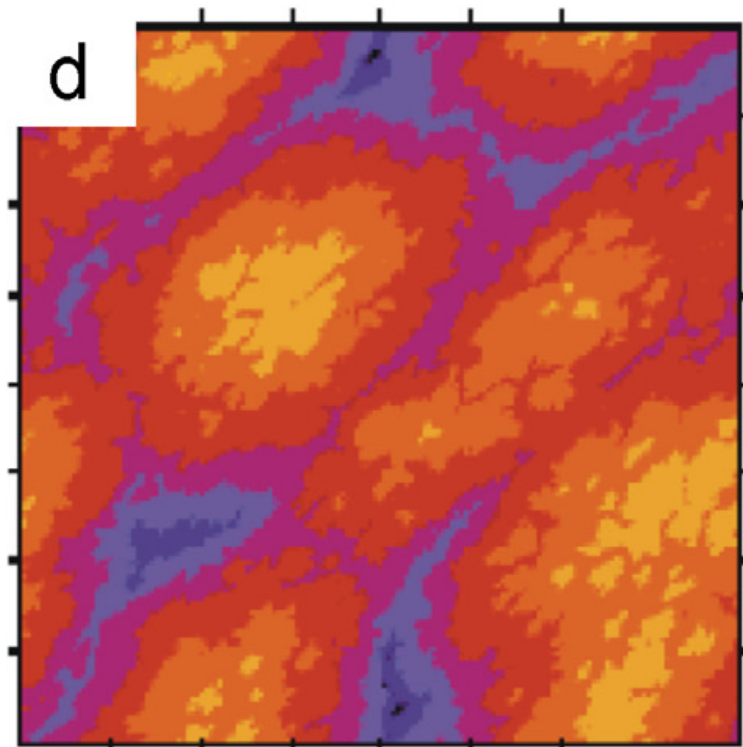
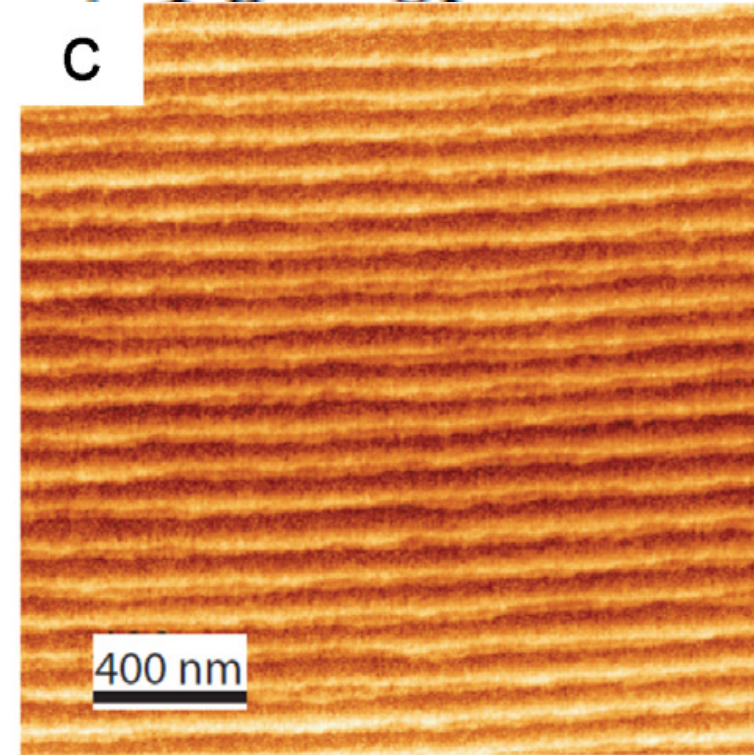
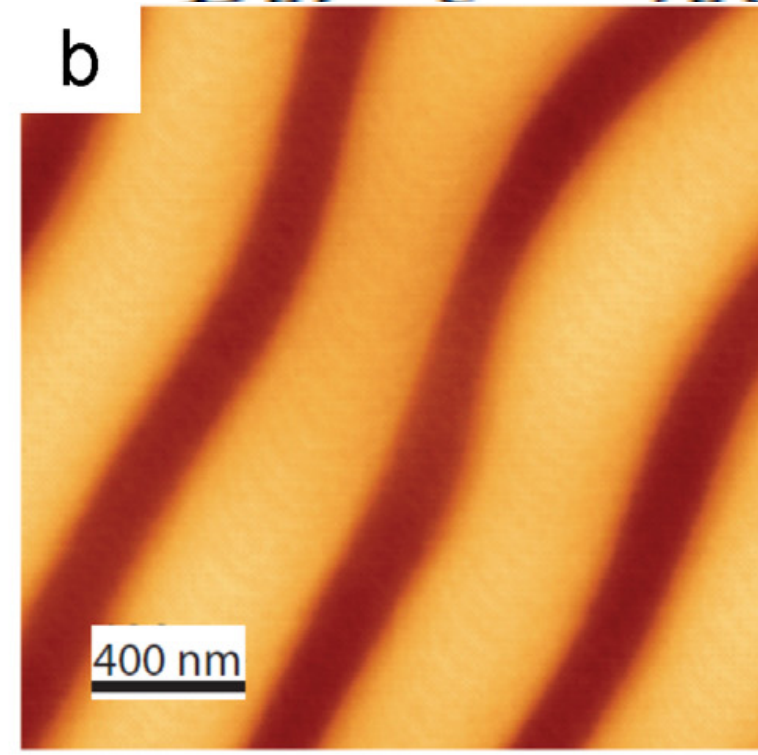
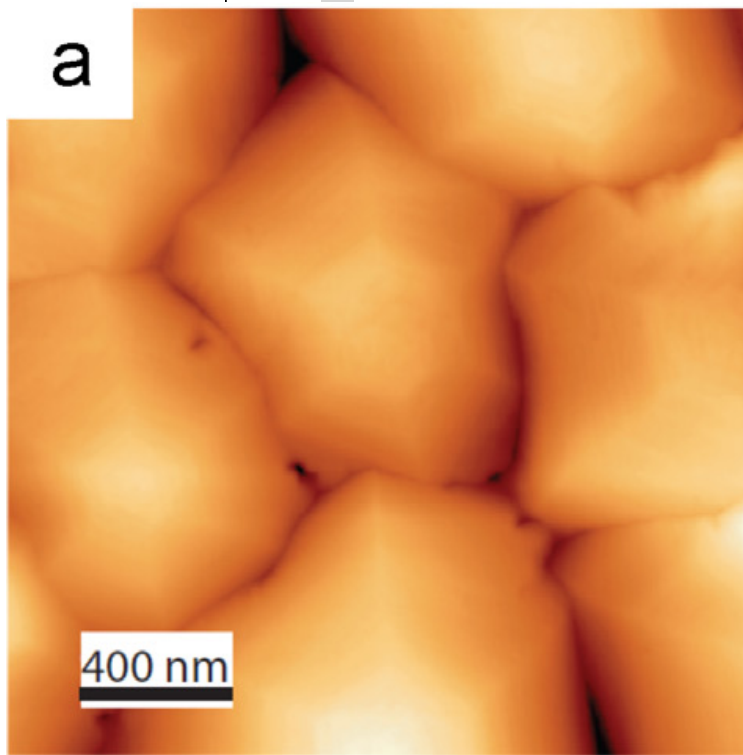
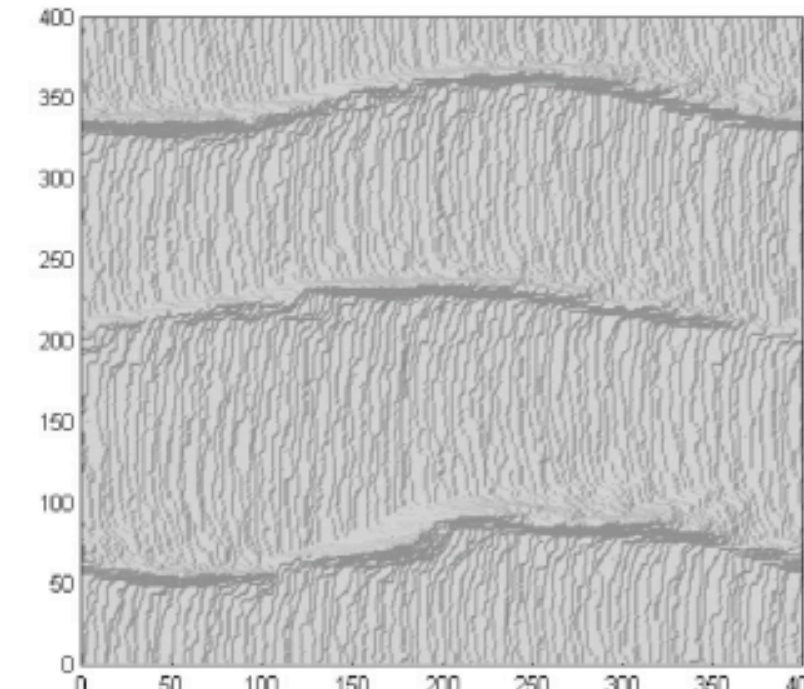
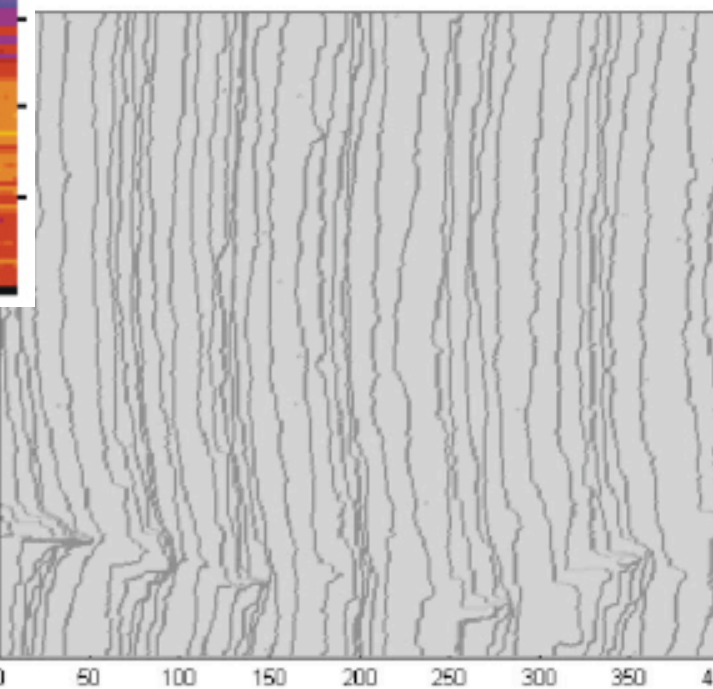
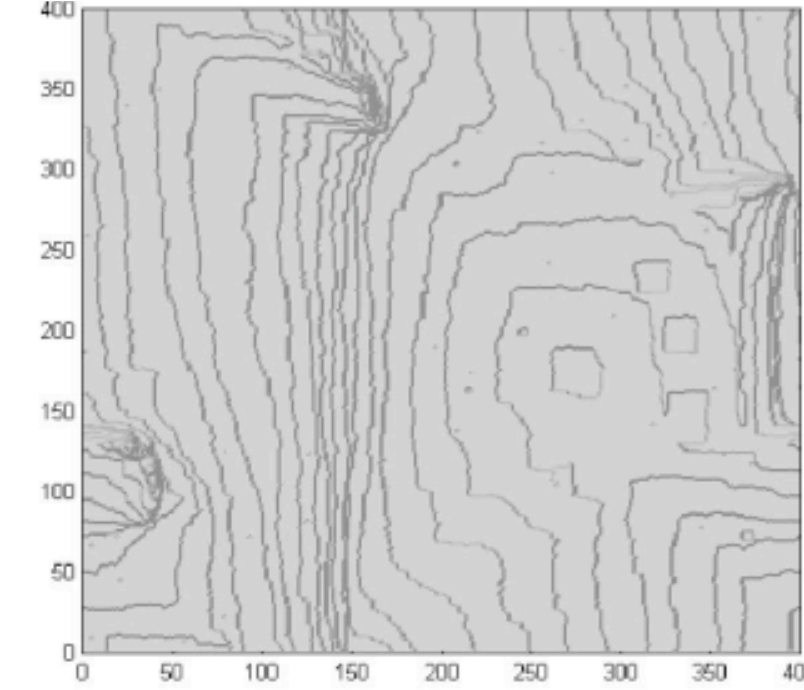
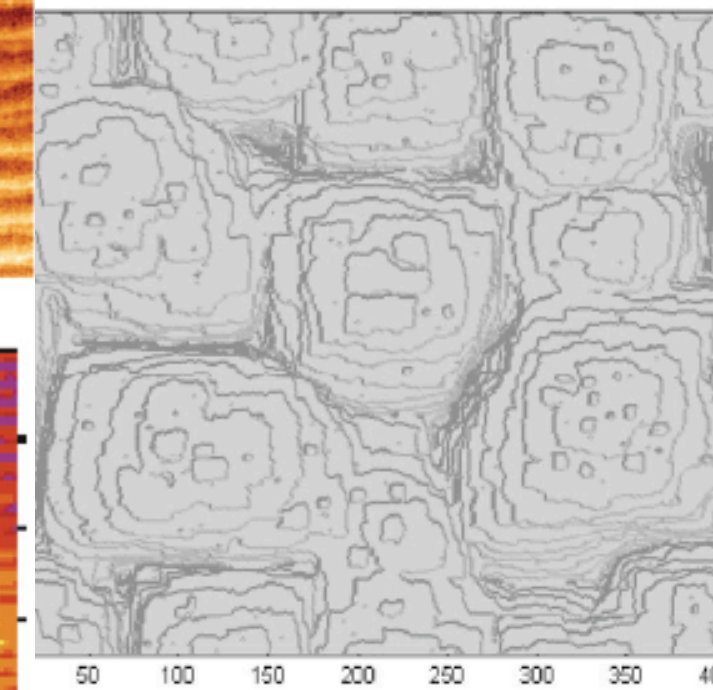
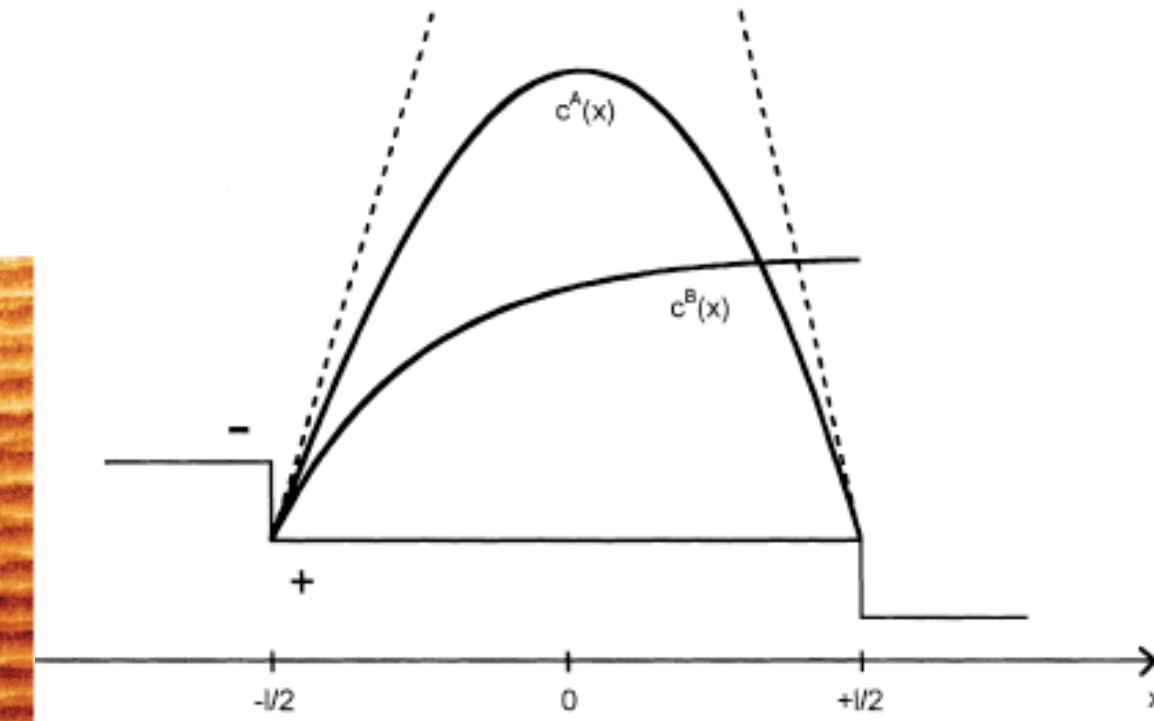
$$D_B \nabla^2 c^B - v_B c^B + \Phi_B = 0,$$

$$D_A \nabla^2 c^A + v_B c^B = 0,$$

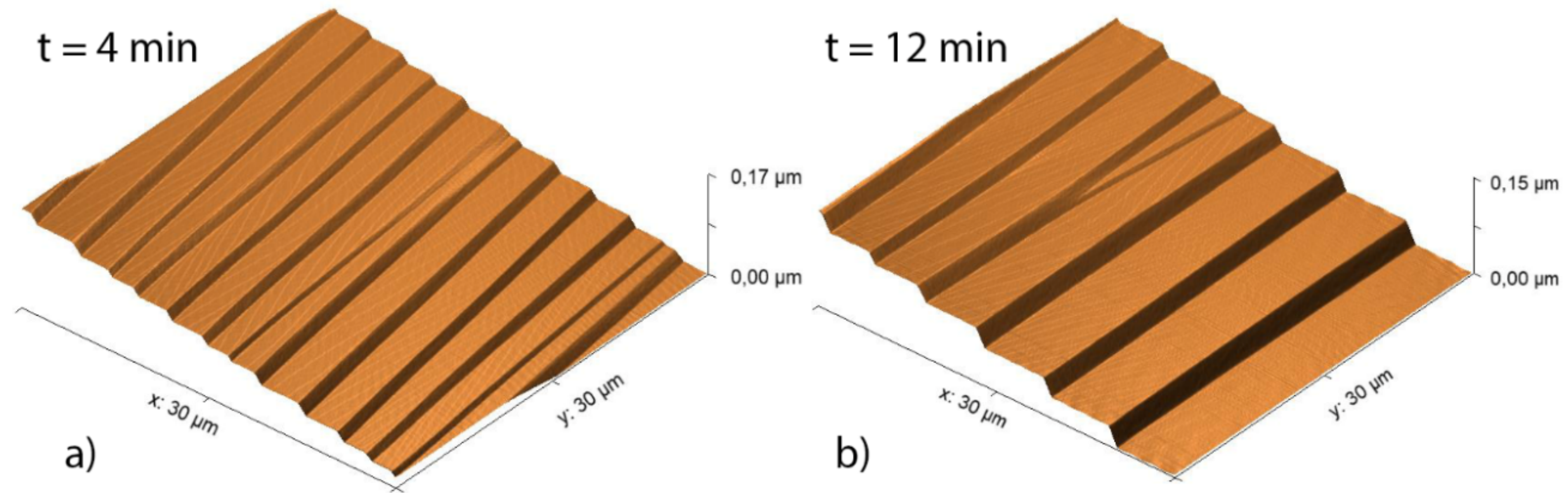


ES barrier at step edges implies coexistence of mounding, bunching *and* meandering...

$$D_B \nabla^2 c^B - v_B c^B + \Phi_B = 0.$$

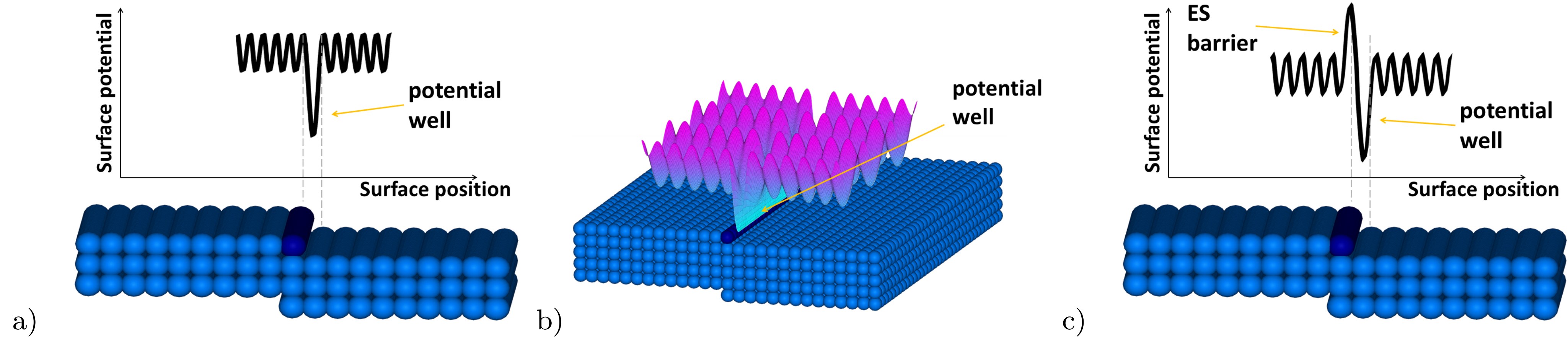


More Exotic Instabilities: Electromigration



AFM images of sublimating Si(111) vicinal surface,
miscut in the $[11-2]$ direction by 0.3° ,
heated by step-down direct current for 4 min (a) and 12 min (b).

More Exotic Models: Cellular Automata



Step meandering: The balance between the potential well and the Ehrlich – Schwoebel barrier

Marta A. Chabowska^{1,*}, Hristina Popova², and Magdalena A. Załuska-Kotur¹

More Exotic Models: Cellular Automata

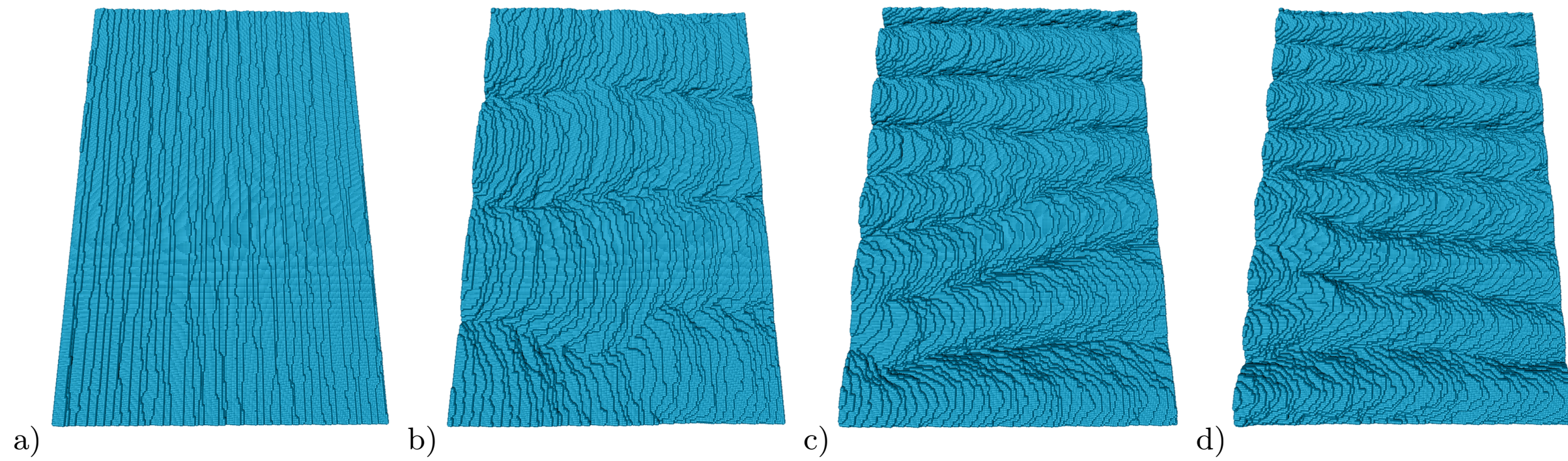


FIG. 2. Meanders obtained for $c_0 = 0.003$, $l_0 = 5$, $n_{DS} = 5$ and a) $\beta E_V = 0.0$, b) $\beta E_V = 2.0$, c) $\beta E_V = 3.5$, d) $\beta E_V = 6.0$, simulation time $t = 2 \cdot 10^6$. System size 200 x 300.

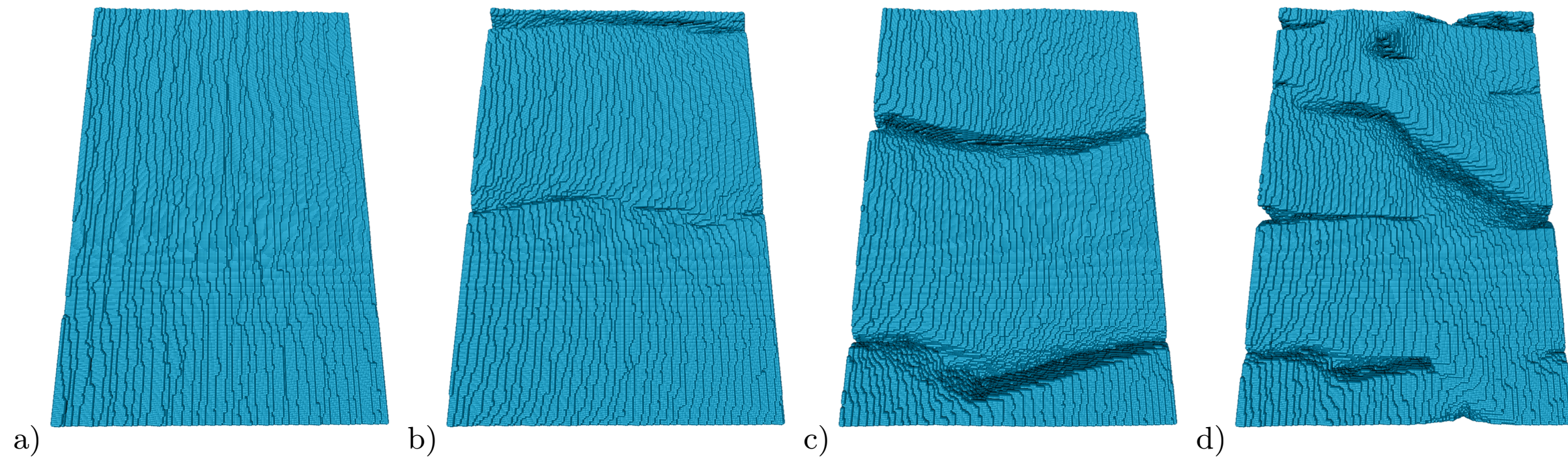


FIG. 3. Meanders obtained for $c_0 = 0.005$, $l_0 = 5$, $n_{DS} = 5$, $\beta E_V = 0.0$ and a) $\beta E_{ES} = 2.0$, b) $\beta E_{ES} = 4.0$, c) $\beta E_{ES} = 6.0$, d) $\beta E_{ES} = 8.0$, simulation time $t = 2 \cdot 10^6$. System size 200 x 300.

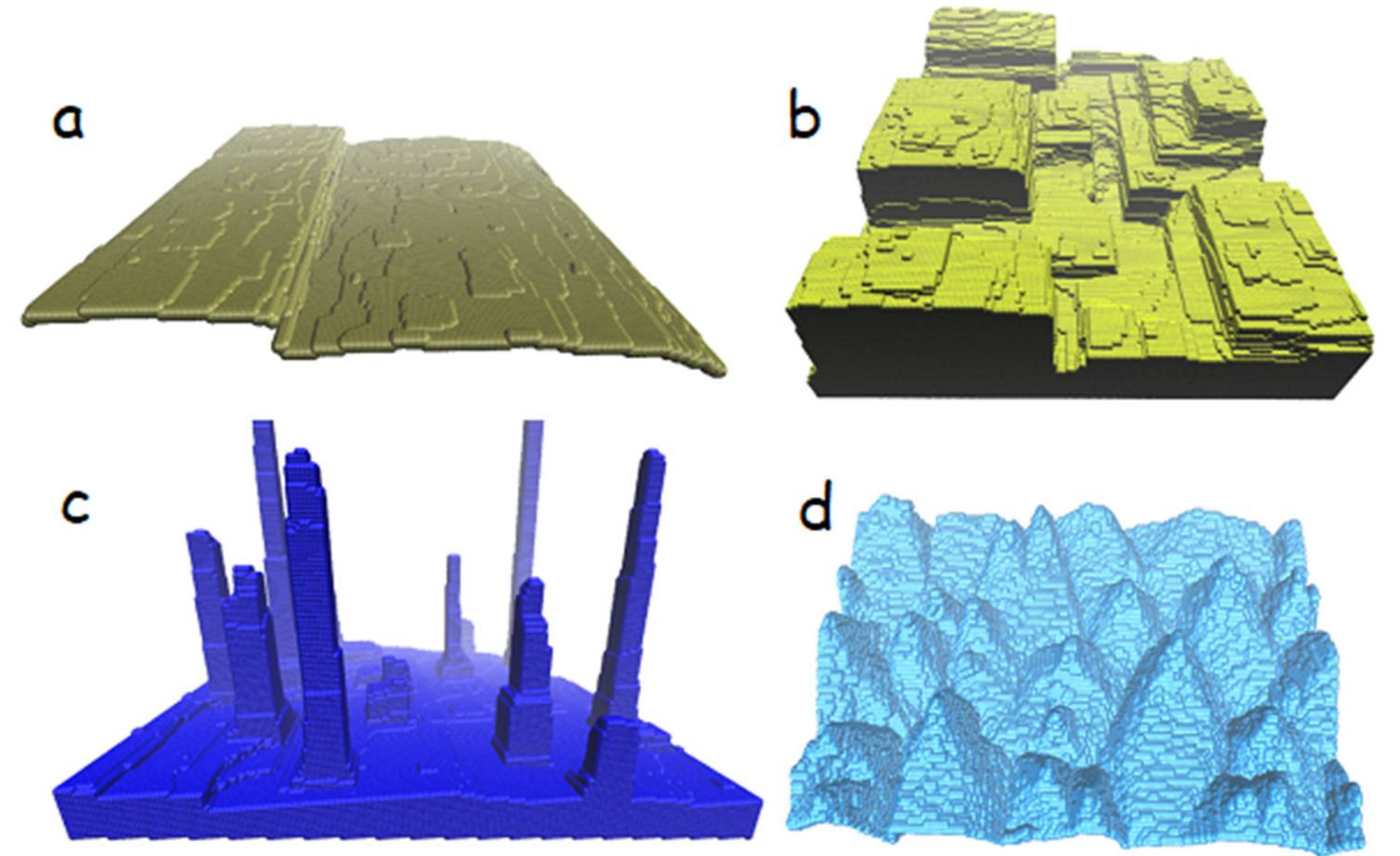
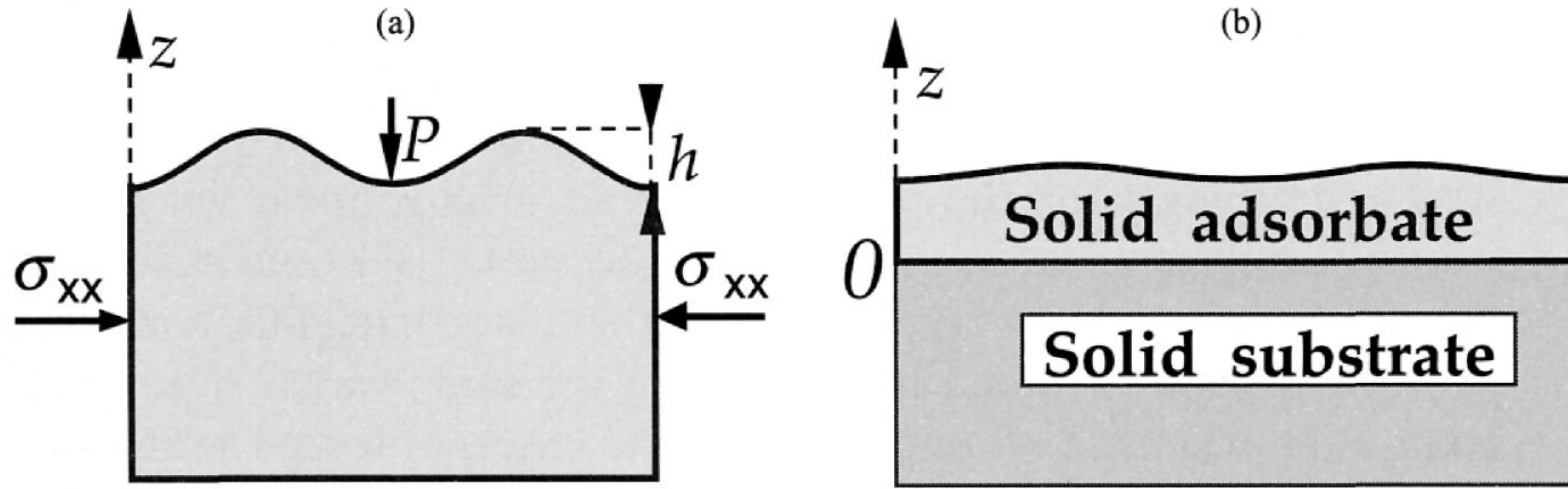


Figure 5. (a) Antibunches, $c_0 = 0.02$, $P_{iES} = 0.2$, $P_{dES} = 0.6$, $p_w = 1.66$, $l_0 = 10$, $n_{DS} = 10$, time steps 3×10^6 . (b) Nanopillars, $c_0 = 0.02$, $P_{iES} = 0.2$, $P_{dES} = 0.4$, $p_w = 2.5$, $l_0 = 10$, $n_{DS} = 10$, time steps 2×10^5 . (c) Nanowires, $c_0 = 0.02$, $P_{iES} = 0.1$, $P_{dES} = 0.2$, $p_w = 5$, $l_0 = 10$, $n_{DS} = 10$, time steps 3×10^5 . (d) Pyramids, $c_0 = 0.02$, $P_{iES} = 1.0$, $P_{dES} = 0.0$, $p_w = 1.0$, $l_0 = 5$, $n_{DS} = 10$, time steps 4×10^5 .

Elastic instabilities: Grinfeld



Misfit-induced strain: $\epsilon_{xx} = \epsilon_{yy} = \epsilon_0 = \delta a / a$

$$p_0 = -\frac{E}{1 - \zeta} \epsilon_0 = \frac{E}{1 - \zeta} \frac{\delta a}{a}$$

Contributions to the free energy of the undulating film $\delta Z(x, y) = Z(x, y) - \bar{Z}$.

$$\delta Z(x) = h \cos(qx)$$

$$1) \quad d\mathcal{F}_{\text{cap}}/d\mathcal{A} = \sigma h^2 q^2 / 2$$

Capillarity (surface free energy)

$$\text{curvature} \sim \partial^2 z / \partial x^2 \rightarrow z_q q^2 \sim h q^2$$

$$dV = h d\mathcal{A}$$

$$2) \quad d\mathcal{F}_{\text{relax}}/d\mathcal{A} \approx -h p_0 \epsilon$$

Relaxation, proportional to elastic energy and to volume change

$$dV = h d\mathcal{A}$$

$$3) \quad d\mathcal{F}_{\text{el}}/d\mathcal{A} \approx C \epsilon^2 / (2q)$$

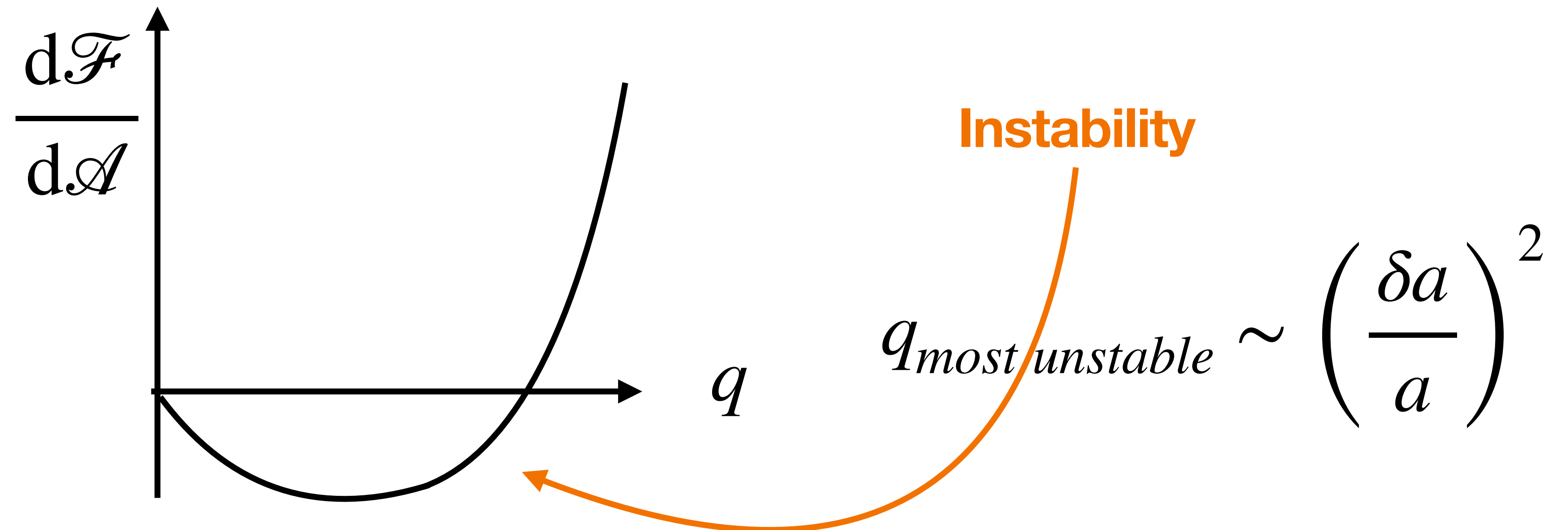
Elastic energy cost proportional to penetration length of the strain $\lambda \approx 1/q$

Minimizing 2 and 3 with respect to ϵ yields:

$$\epsilon \approx hqp_0/C$$

The total free energy in the undulating film is the sum of 1), 2) and 3):

$$d\mathcal{F}/d\mathcal{A} = \sigma h^2 q^2/2 - h^2 p_0^2 q/(2C)$$



Questions?

